Riemannian Geometry IV

Solutions, set 8.

Exercise 19. (a) Note that T(X, Y) = -T(Y, X), so we only have to prove linearity in the first argument. Moreover, we obviously have $T(X_1+X_2, Y) = T(X_1, Y) + T(X_2, Y)$. We are left with showing that

$$T(fX,Y) = fT(X,Y).$$

The calculation for this goes as follows:

$$T(fX,Y) = [fX,Y] - (\nabla_{fX}Y - \nabla_Y fX)$$

= $f[X,Y] - (Yf)X - (f\nabla_X Y - (Yf)X - f\nabla_Y X)$
= $f([X,Y] - (\nabla_X Y - \nabla_Y X)) - (Yf)X + (Yf)X = fT(X,Y).$

(b) It is, again, straightforward to check that

$$\nabla A(X_1, \dots, X_i + \tilde{X}_i, \dots, X_r, X_{r+1})$$

= $\nabla A(X_1, \dots, X_i, \dots, X_r, X_{r+1}) + \nabla A(X_1, \dots, \tilde{X}_i, \dots, X_r, X_{r+1}),$

for i = 1, 2, ..., r + 1. So it remains to show that

$$\nabla A(X_1,\ldots,fX_i,\ldots,X_r,X_{r+1}) = f\nabla A(X_1,\ldots,X_i,\ldots,X_r,X_{r+1}),$$

for i = 1, 2, ..., r + 1. Let i = 1, 2, ..., r. Then

$$\nabla A(X_1, \dots, fX_i, \dots, X_r, Y)$$

= $Y(fA(X_1, \dots, X_r)) - f \sum_{j=1}^n A(X_1, \dots, \nabla_Y X_j, \dots, X_r) - (Yf)A(X_1, \dots, X_r)$
= $fY(A_1(X_1, \dots, X_r)) - f \sum_{j=1}^n A(X_1, \dots, \nabla_Y X_j, \dots, X_r)$
= $f \nabla A(X_1, \dots, X_r, Y).$

Finally, we obtain

$$\nabla A(X_1, \dots, X_r, fY)$$

= $fY(A(X_1, \dots, X_2)) - \sum A(X_1, \dots, f\nabla_Y X_j, \dots, X_r)$
= $f\nabla A(X_1, \dots, X_r, Y).$

(c) Using (b), we obtain

$$\nabla G(X, Y, Z) = Z(\langle X, Y \rangle) - \langle \nabla_Z X, Y \rangle - \langle X, \nabla_Z Y \rangle.$$

Then $\nabla G \equiv 0$ means precisely that the affine connection ∇ has the "Riemannian property".

Exercise 20. (a) We have

$$\frac{\partial}{\partial x_1}\Big|_{\varphi^{-1}(x_1,x_2)} = (f'(x_1)\cos x_2, f'(x_1)\sin x_2, g'(x_1)),\\ \frac{\partial}{\partial x_2}\Big|_{\varphi^{-1}(x_1,x_2)} = (-f(x_1)\sin x_2, f(x_1)\cos x_2, 0).$$

This implies that

$$(g_{ij}) = \begin{pmatrix} (f'(x_1))^2 + (g'(x_1))^2 & 0\\ 0 & f^2(x_1) \end{pmatrix}$$

and

$$(g^{ij}) = \begin{pmatrix} \frac{1}{(f'(x_1))^2 + (g'(x_1))^2} & 0\\ 0 & \frac{1}{f^2(x_1)} \end{pmatrix}.$$

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Consequently, the Christoffel symbols are calculated as

$$\begin{split} \Gamma_{11}^{1} &= \frac{1}{2}g^{11}\left(g_{11,1} + g_{11,1} - g_{11,1}\right) = \frac{f'(x_1)f''(x_1) + g'(x_1)g''(x_1)}{(f'(x_1))^2 + (g'(x_1))^2} \\ \Gamma_{11}^{2} &= \frac{1}{2}g^{22}(g_{12,1} + g_{12,1} - g_{11,2}) = 0, \\ \Gamma_{12}^{1} &= \frac{1}{2}g^{11}(g_{11,2} + g_{12,1} - g_{12,1}) = 0 = \Gamma_{21}^{1}, \\ \Gamma_{12}^{2} &= \frac{1}{2}g^{22}(g_{12,2} + g_{22,1} - g_{12,2}) = \frac{f'(x_1)}{f(x_1)} = \Gamma_{21}^{2}, \\ \Gamma_{22}^{1} &= \frac{1}{2}g^{11}(g_{12,2} + g_{12,2} - g_{22,1}) = \frac{-f(x_1)f'(x_1)}{(f'(x_1))^2 + (g'(x_1))^2}, \\ \Gamma_{22}^{2} &= \frac{1}{2}g^{22}(g_{22,2} + g_{22,2} - g_{22,2}) = 0. \end{split}$$

This implies that

$$\nabla_{\frac{\partial}{\partial x_1}} \frac{\partial}{\partial x_1} = \frac{f'(x_1)f''(x_1) + g'(x_1)g''(x_1)}{(f'(x_1))^2 + (g'(x_1))^2} \frac{\partial}{\partial x_1},$$

$$\nabla_{\frac{\partial}{\partial x_1}} \frac{\partial}{\partial x_2} = \frac{f'(x_1)}{f(x_1)} \frac{\partial}{\partial x_2},$$

$$\nabla_{\frac{\partial}{\partial x_2}} \frac{\partial}{\partial x_1} = \frac{f'(x_1)}{f(x_1)} \frac{\partial}{\partial x_2},$$

$$\nabla_{\frac{\partial}{\partial x_2}} \frac{\partial}{\partial x_2} = \frac{-f(x_1)f'(x_1)}{(f'(x_1))^2 + (g'(x_1))^2} \frac{\partial}{\partial x_1}$$

(b) Note that we have

$$\gamma_1'(t) = \frac{\partial}{\partial x_1}|_{\gamma_1(t)}.$$

This implies that

$$\begin{pmatrix} \frac{D}{dt}\gamma_1' \end{pmatrix}(t) = \nabla_{\gamma_1'(t)}\frac{\partial}{\partial x_1} = \left(\nabla_{\frac{\partial}{\partial x_1}}\frac{\partial}{\partial x_1}\right)(\gamma_1(t))$$
$$= \frac{f'(x_1+t)f''(x_1+t) + g'(x_1+t)g''(x_1+t)}{(f'(x_1+t))^2 + (g'(x_1+t))^2}\frac{\partial}{\partial x_1}|_{\gamma_1(t)} \in T_{\gamma_1(t)}M.$$

The condition $\frac{D}{dt}\gamma'_1 \equiv 0$ is equivalent to f'(t)f''(t) + g'(t)g''(t) = 0 for all $t \in (a, b)$. This is equivalent to $(f'(t))^2 + (g'(t))^2 = constant$. Since

$$||c'(t)||^2 = (f'(t))^2 + (g'(t))^2,$$

we conclude that $\frac{D}{dt}\gamma'_1$ vanishes identically if and only if c is parametrized propertional to arc-length. Since c and γ_1 are obtained from each other by an isometry of \mathbb{R}^3 , namely a rotation by the angle x_2 around the vertical Z-axis, c is parametrized proportional to arc-length if and only if γ_1 is parametrized proportional to arc-length. As explained in Example 25, the property $\frac{D}{dt}\gamma'_1 \equiv$ 0 is equivalent to the fact that γ_1 is a geodesic.

(c) We have

$$\gamma_2'(t) = \frac{\partial}{\partial x_2}|_{\gamma_2(t)}.$$

This implies that

$$\begin{pmatrix} \frac{D}{dt}\gamma_2' \end{pmatrix}(t) = \nabla_{\gamma_2'(t)}\frac{\partial}{\partial x_2} = \left(\nabla_{\frac{\partial}{\partial x_2}}\frac{\partial}{\partial x_2}\right)(\gamma_2(t)) \\ = \frac{-f(x_1)f'(x_1)}{(f'(x_1))^2 + (g'(x_1))^2}\frac{\partial}{\partial x_1}|_{\gamma_2(t)} \in T_{\gamma_2(t)}M.$$

Since f > 0, the condition $\frac{D}{dt}\gamma'_2 \equiv 0$ is equivalent to $f'(x_1) = 0$. Since γ_2 is a parallel of the surface of revolution M and $f_1(x_1)$ its radius, the condition $\frac{D}{dt}\gamma'_2 \equiv 0$ is satisfied, e.g., if f assumes a local maximum or minimum at x_1 . Again, as explained in Example 25, the property $\frac{D}{dt}\gamma'_2 \equiv 0$ is equivalent to the fact that γ_2 is a geodesic.

(d) Since c is assumed to be arc-length parametrized, we have

$$(f'(x_1))^2 + (g'(x_1))^2 = 1$$
 for all $x_1 \in (a, b)$.

This implies that

$$(g_{ij}) = \begin{pmatrix} 1 & 0\\ 0 & f^2(x_1) \end{pmatrix},$$

and therefore, $\sqrt{\det(g_{ij})} = f > 0$. Since φ is an almost global coordinate chart, we conclude that

$$\operatorname{vol}(M) = \int_0^{2\pi} \int_a^b f(x_1) \, dx_1 \, dx_2 = 2\pi \int_a^b f(x_1) \, dx_1.$$