

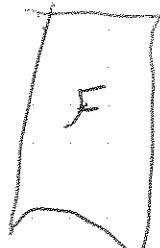
PLAN

A THETA LIFT IN  $SL(2, \mathbb{Z})$  AND  
LOCALLY HARMONIC MAASS FORMS

- HISTORY/MOTIVATION
- DEFINITIONS
  - LATTICE
  - GRASSMANIAN
  - SIEGEL THETA FUNCTIONS
  - HARMONIC WEAK MAASS FORMS
- MY SETTING
- THE LIFT
- LOCALLY HARMONIC MAASS FORMS
- ~~RELATIONSHIP~~ RELATIONSHIP WITH SHIMURA LIFT

$$T = \begin{pmatrix} 1 & \\ & 1 \end{pmatrix}$$

$$S = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$$



KEY DEFINITION

First 5 minutes reminder of final slides from last talk

$M_k$  MODULAR FORMS

key definition needed before proceeding

$\tau = u + iv$

-  $k \in \mathbb{Z}$ , even, non-negative. A function  $f: \mathbb{H} \rightarrow \mathbb{C}$  of weight " $k$ " or subgroup

- 1)  $f$  is holomorphic on  $\mathbb{H}$  - i.e. "nice", complex differentiable
- 2)  $f$  is weakly modular of weight  $k$ . -  $\gamma \tau = \frac{a\tau + b}{c\tau + d} \implies f(\gamma \tau) = (c\tau + d)^{-k} f(\tau)$  -  $\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{Z})$
- 3)  $f$  is holomorphic at  $\infty$ . - No singularities

PROPERTIES

-  $f(\tau + 1) = f(\tau) = \sum_{n=0}^{\infty} a_n e^{2\pi i n \tau} = \sum_{n=0}^{\infty} a_n q^n$  if  $a_n(0) = 0$  "cusp form" vanishes at singularities

- finite dimensional

Next I give an overview, which kind of ruins the surprise but will hopefully keep people intrigued before lots of definitions.

# HISTORY/MOTIVATION

- Modular forms are useful - as briefly described a lot of interesting applications  
 First examples of  $\frac{1}{2}$ -weight  $M_k$ , last time, Serret, physics etc.  
 (WILL DEFINE THESE SHORTLY) if people can't remember

- Theta functions (classical)

$$\Theta(\tau) = \sum_{\lambda \in L} e^{i\pi \frac{(\lambda, \lambda)}{2} \tau} \quad \text{or } \lambda^2 \in \frac{1}{2}\mathbb{Z} \quad \text{L a lattice}$$

- Kissing problem or  $\sqrt{6}\pi$

- Representation numbers - number of ways of summing squares to get "n".

## Shimura List

- Kickstarted theory of  $\frac{1}{2}$  modular forms, 1970s

- Note can look at general  $\frac{1}{q}$  weight, but not interesting, hard to work with

- Desired a map

$$\begin{aligned} \frac{H_{k+\frac{1}{2}}}{\mathbb{Z}} \cdot S_{k+\frac{1}{2}} &\rightarrow M_{2k} \\ \theta(z) = \sum_{n \in \mathbb{Z}} b(n) e^{2\pi i n z} &\rightarrow \sum_{n \in \mathbb{Z}} d(n) e^{2\pi i n z} \end{aligned}$$

$\rightarrow \sum_{n=1}^{\infty} a(n) n^{-s} = (s-k+1) \sum_{n=1}^{\infty} b(n^2) n^{-s}$

- Actually shown later roughly

$$\theta \mapsto \int_{\Gamma} \theta(z) \Theta(\tau) \quad \text{"a theta test"}$$

- Tunnels, congruent number problem



## $\frac{1}{2}$ Forms

Heegner

in the pure maths colloquium

- Constructing  $\hat{\pi}$  points on elliptic curves, last week Vlad Dokchister, BSD rank 1  
 - Class numbers - generating function of imaginary quadratic fields etc.

~~What~~ Harmonic <sup>weak</sup> Maass forms, locally harmonic Maass forms

$$H_k$$

- $\neq 1$  no cosets,  $k$  even
- $\neq$  neg
- $\neq$  pickup
- $\neq$  singularities
- $\neq$  unimodular even

- Generalizes modular forms

- Romanians work theta functions, is holomorphic part of  $\theta$  one of these  $(\frac{1}{2}$  weights)  
 1920 letter to Hardy,  $1729 = 1^3 + 12^3 = 9^3 + 10^3$

- Ono, Bruinier give  $\hat{\pi}$  use very similar "theta list" to one I describe today, formula of partition function

$$\int_{\mathcal{C}} \gamma(t) \Theta(t, z) dt$$



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AIM

$$H_{3/2-k}$$

$$LH_{2-2k}$$

Surprise

signature

$$S_{k+1/2}$$

Skinner

$$M_{2k}$$

$k \in \mathbb{N}$

$k=1, 2, 3, \dots$

$$Q(\gamma x) = r^2 Q(x)$$

$$(x, y) = Q(x+y) - Q(x) - Q(y)$$

DIFFERENT "k"

$k$  here = 1, 2, 3, 4, 5, ...

DEFINITIONS

LATTICE

— need this to define theta's

We define a lattice:

$L$  of signature  $(b^+, b^-)$

$V$  a rational vector space, quadratic

"a free  $\mathbb{Z}$  module  $L \subset V$ "  $V = L \otimes \mathbb{Q}$

$$V(\mathbb{R}) = V \otimes \mathbb{R}$$

$L \cong \mathbb{Z}^m$  usually

with an attached quadratic form  $Q(x) = \frac{1}{2}(x, x)$  bilinear form

isometric

to  $\mathbb{R}^{b^+, b^-} = \mathbb{R}^{b^+ + b^-}$

$$Q(x) = x_1^2 + \dots + x_{b^+}^2 - x_{b^+ + 1}^2 - \dots$$

ISOMETRIC

~~isometric~~  
signatureless  
map  
 $\sigma: V \rightarrow V$   
 $Q(\sigma(x)) = Q(x)$

- Assume  $L$  even —  $Q(x) \in \mathbb{Z}$

- Assume  $L$  unimodular —  $\det(T) = 1, -1$

$$T = \begin{pmatrix} (b_1, b_1) & & \\ & (b_1, b_1) & \\ & & (b_2, b_2) & \\ & & & (b_2, b_2) \end{pmatrix}$$

big cheat here, simplifies exposition for today. i.e. scalar valued forms

GRASSMANNIAN

$$Gr(L) = \{ Z \subset V(\mathbb{R}) \mid \dim Z = b^- \text{ and } Q|_Z < 0 \}$$

"set of negative definite  $b^-$ -dimensional subspaces"

can think of this as a section to  $\mathbb{H}$  actually in my case

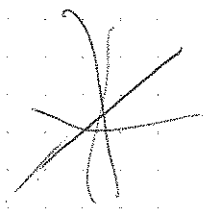
$$Q_Z(\lambda) = Q(\lambda_Z) - Q(\lambda_{\bar{Z}})$$

for a fixed  $Z = x + iy$

$$\lambda = \lambda_Z + \lambda_{\bar{Z}}$$

always positive

in case  $b^- = 1$



take line of length 1,  
 $Z$ -shell  
hyperboloid

like  $\mathbb{H}$

# SIEGEL THETA FUNCTION

the quadratic form is

"Generalises earlier version, which diverged to negative definite as  $V \rightarrow \infty$ "

$$\Theta_L(\tau, z, \rho) = v^{\frac{k^+ + m^-}{2}} \sum_{\lambda \in \Gamma} p(\lambda) e(Q(\lambda)\tau + Q_2(\lambda)z)$$

$$\lambda = \lambda_1 \mathbf{b}_1 + \lambda_2 \mathbf{b}_2 + \dots$$

basis of  $\mathbb{Z}^n$

"polynomial" degree  $m^+, m^-$  in  $\lambda_1, \lambda_2, \lambda_3, \dots$

- weight  $\frac{k^+ - k^-}{2} + m^+ - m^-$  modular form in  $\tau$ .

Invariant under  $Q(L)$  (preserves quadratic), polynomial ~~relates~~ will create a modular form weight in  $\tau = x + iy$

- Surjective homo,  $\mathbb{Z} \oplus \text{Spin}(V) \rightarrow \text{SO}^+(V)$ , acting via conjugation

## NOTE ON $\frac{1}{2}$ WEIGHT FORMS

-  $g(\gamma\tau) = (c\tau + d)^{k/2} g(\tau)$  leads to contradictions  
choice of square root

$$g(\gamma\tau) = \left(\frac{c}{d}\right)^{2k+1} e^{2\pi i k} (c\tau + d)^{k/2} g(\tau) \leftarrow \text{can define half weight forms like this}$$

Junk

group of linear maps, isometries preserving  $\mathbb{Q}$

$$= \text{SL}_2(\mathbb{Q})$$

in our case

- actually fractional weight forms should be under "central extensions of  $\text{SL}_2(\mathbb{Z})$ , i.e. braid group  $B_3$ "

-  $M_{2,2}(\mathbb{Z})$  unique connected double cover of  $\text{SL}_2(\mathbb{Z})$

Subgroup of this

$$\overline{M}_{P_2}(\mathbb{R}) = (\gamma, \phi_\gamma) \quad \gamma \in \text{SL}_2(\mathbb{R}) \quad \gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

$$\phi_\gamma(\tau)^2 = (c\tau + d)$$

$$(\gamma, \phi_\gamma(\tau)) (\gamma', \phi_{\gamma'}(\tau)) = (\gamma\gamma', \phi_{\gamma\gamma'}(\tau))$$

- so junk is projected from reps of  $M_{P_2}(\mathbb{Z})$

- I work in case non-unimodular, so have "vector valued forms" ( )

- unique rep, "weil representation"

$\mathbb{Z}/2$  has cosets  $e_h$

$$P_L(\tau)(e_h) = e(Q(h))e_h$$

$$P_L(S)(e_h) = \frac{e((b-b^+/\delta))}{\sqrt{|\mathbb{Z}/2|}} \sum_{k \in \mathbb{Z}/2} e(-\langle h, k \rangle) e_k$$

NATURAL from Heisenberg group

automorphisms of Heisenberg six centre, are symplectic groups, gives action, using stone von Neumann theorem

$g(\tau) = \text{HARMONIC WEAK MAASS FORMS}$  - Generalize modular forms:

②  $\Delta_k(g) = 0$

where  $\Delta_k = -v^2 \left( \frac{\partial^2}{\partial u^2} + \frac{\partial^2}{\partial v^2} \right) + i k v \left( \frac{\partial}{\partial u} + i \frac{\partial}{\partial v} \right)$



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- not finite dimensional

-  $g = g^+ + g^-$ ,  $g^+ = \sum_{n \geq 0} c^+(n) e(n\tau)$ ,  $g^- = \sum_{n < 0} c^-(n) \Gamma(1-k, 4\pi n v) e(n\tau)$

- for  $k \geq 2$   $g^-$  vanishes \ holomorphic part | non-holomorphic

$g \in H_k$  OPERATORS

- we have lowering, raising operators

$R_k := z i \frac{\partial}{\partial \tau} + k v^{-1}$        $L_k := -2i v^2 \frac{\partial}{\partial \tau}$

$R_k(g)$  ~~weight~~ weight  $k+2$        $L_k(g)$  weight  $k-2$

$\hat{L}_k(g) = v^{k-2} L_k(g) = R_{-k} v^k g(\tau)$  basically the same

Theorem

$\hat{L}_k : H_k \rightarrow S_{2-k}$  surjectively Jens + Bruinier

- Background Done

he did this in

- My work based on german thesis, case  $k=1$ , MY SETTINGS

-  $V \cong \mathbb{Z} \oplus \mathbb{Z}$

-  $V = \text{Mat}_2(\mathbb{Q})$ , traceless

-  $Q(X) = -\det(X)$

-  $L = \begin{pmatrix} b & -a \\ c & -b \end{pmatrix}$   $a, b, c \in \mathbb{Z}$

-  $G(L) \cong \mathbb{Z} \times \mathbb{Z} \hookrightarrow \text{Hyperbolic space}$        $\mathbb{Z} = X + iy$

-  $\forall \lambda \in V \exists \lambda = \lambda_1(z) b_1(z) + \lambda_2(z) b_2(z) + \lambda_3(z) b_3(z)$  orthonormal basis

$\lambda_3(z) = \frac{1}{\sqrt{2}y} (c|z|^2 - bx + a)$

# MY THETA

$$\Theta_k(\tau, z) = v^{\frac{k-1}{2}} \sum_{\lambda \in L} \underbrace{\lambda_3(z) (\lambda_2(z) + i \lambda_1(z))}^{\#(\lambda)} e^{(Q_1(u) + Q_2(v)iv)} = \lambda_1(z)^2 + \lambda_2(z)^2 - \lambda_3(z)^2$$

Shimura theta

$$\lambda_1(z) - i \lambda_2(z)$$

- ask the audience

- weight  $\frac{1}{2} - 1 + k - 1 - 1 = k - \frac{3}{2}$  in  $\tau$ ,  $MP_2(\mathbb{R})$

- weight  $2 - 2k$  in  $z$  ( $\#_z(\delta, \lambda) = \#_z(\emptyset, \lambda) (cz+d)^{-2}$ )

- in  $\Gamma_0(N) \subset Spin(V) = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{Z})$   $c \equiv 0 \pmod{N}$

## THETA LIFT


$$\mathbb{D}_k(z, g) = \int_{\mathbb{R}^2}^{reg} g(\tau) \Theta(\tau, z) \frac{dudv}{v^2}$$

actually peterson scalar product

~~In our case ( $g \in H_0$ )~~

-  $g \in H_k$ , then  $\Theta$  in  $\tau \in M_k$  to make sense

- Also as ( $g \in H_k$ ) can grow exponential but  $v \rightarrow \infty$

- So take  $\lim_{v \rightarrow \infty} \int_{\mathbb{R}^2}$  

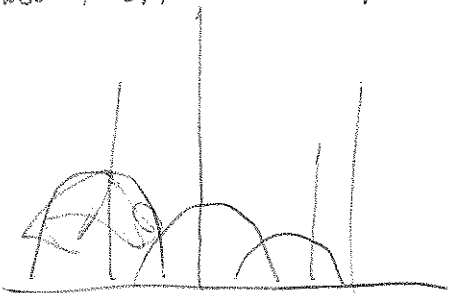
- Can prove (some work) this converges.

- In our case  $g \in H_{\frac{3}{2}-k}$

-  $\mathbb{D}_k(z, g)$  clearly of weight  $2 - 2k$  (but what is it?)

## EXCITING SINGULARITIES

- Smooth on  $\mathbb{H}$  except



- Step/Jump singularities

(ie. jump of  $\sum''$  or  $\lambda(z) + i\lambda_2(z) = (cz^2 - bz + a)^{k-1}$  polynomial jump



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- Why?

- When showing convergence, turns out only finite number of vectors  $\lambda$  contribute to potential singularities ~~at~~ and  $\lambda \perp z$

and  $\int_{v=1}^{\infty} e(\rho Q(u)) i v^{-1/2} dv$  ~~is~~ ~~singular~~

$$\begin{pmatrix} b/2 & -a \\ c & -b/2 \end{pmatrix} \begin{pmatrix} -x & x+y^2 \\ -1 & x \end{pmatrix}$$

$$c|z|^2 - bz + a = 0$$

$$= \Pi\left(\frac{1}{z}, -Q^T Q(u)\right)$$
 has singularities at  $O = Q(u)$

$\lambda \perp z$ ,  $\lambda$  fixed then  $z$  defines a geodesic

and  $C^+(n, h)$

$-\infty < h < \infty$   
so finitely many geodesics

LOCALLY HARMONIC MASS FORMS

- Defined ~~recently~~ recently

- Same as ~~the~~  $H_k$  but have singularities  $LH_k$

$$\Theta_k(z, g) \in LH_{2-2k}$$

- Proof (Very sketchy)

$$4\Delta_{k-3/2} \Theta_R(t, z) = \Delta_{2-2k} \Theta_R(\tau, z) + (6+4k) \Theta_R(\tau, z)$$

$\uparrow$  in  $t$                        $\uparrow$  in  $z$

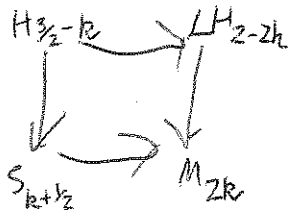
- ~~proof using~~ Casimir element, Lie Algebra, chain rule

- and 
$$\int_{F_R} g \overline{\Delta_R g} v^{k-2} dudv - \int_{F_R} \Delta_R g \overline{g} v^{k-2} dudv = \int_{\partial F_R} \overline{g} g$$
 STOKES      ~~g~~  $g$    
 ~~stoss~~

- So 
$$\Delta_{2-2k} \Theta_R(\tau, z) = \int g \cdot \Delta_{2-2k} \Theta = \int \Delta_{2-2k}(g) \Theta = 0$$

# EXPANSION

- Can check



about with  $\xi$  and  
 Shimura  $\Theta \leftrightarrow$  my  $\Theta$

~~Shimura  $\Theta \leftrightarrow$  my  $\Theta$~~

so  $\Phi_{\text{Shim}}(\xi_{3/2-k}(\theta)) = \xi_{2-2k}(\Phi_k(\theta))$

- Also would like explicit expansion

- lots of work, just wildaspoine series

- have to carefully re-

$$\Theta(t) = \sum_{a,b=1}^{\infty} (ct+ad)^k g\left(\frac{at+b}{ct+ad}\right)$$

then  $\int_{F_T}$  so  $\int_0^{1/2} \int_{-1/2}^{1/2} g(\tau) g(\bar{\tau}) \frac{dudv}{\sqrt{v}}$

Shimura  $\Phi(\xi_{3/2-k}(\theta))$   
 $= \int \xi(\theta) \theta^i = \int \theta \xi \theta^i$   
 $= \int (\int \theta \theta^i)$   
 $= \text{my way}$

$$- \sum_{m \geq 1} \sum_{0 \leq h \leq k} \binom{k-1}{h} \frac{1}{\sqrt{h+1}} (my)^{k-1-h} B_{1+h}(-mx) c^+\left(-\frac{m^2}{4}, \frac{rm}{2}\right)$$

$$+ \sum_{m \geq 1} \sum_{n \geq 1} \sum_{0 \leq i \leq 2k-2} (my)^k (\pi my)^{k-2-i} (\text{stuff}) e(-nmz) c^-\left(-\frac{m^2}{4}, \frac{rm}{2}\right)$$

+ constant (L-function at  $\theta$ )  $\Gamma(\dots)$

$$+ \frac{1}{y} \int \theta = \left( \sum_{m \geq 1} m^{1+k} c^+\left(-\frac{m^2}{4}, \frac{rm}{2}\right) \right)$$

$\theta$  - 1-dim

- Full Expansion of lift,  
 (Time permitting)

Shimura  $\theta \in M_{k+1/2}$   $g = \sum c(n,h) e(h\tau)$

$$= \text{constant} + \sum_{n \geq 1} \sum_{\substack{d \geq 1 \\ d|h}} d^{k-1} \left\langle c\left(\frac{h^2}{d^2 4}, \frac{hr}{2d}\right) e(nz) \right\rangle$$

Expansion of Shimura lift

after  $\xi$

$$\Phi(\xi_{3/2-k}(\theta)) = A_{\theta}(\xi_{3/2-k}(\theta)) - (\pi)^{2-k} \sum_{n \geq 1} \sum_{\substack{d \geq 1 \\ d|h}} \frac{n}{d^k} c^-\left(\frac{h^2}{d^2 4}, \frac{hr}{2d}\right) e(nz)$$