# Cartan's C-class equations: solving differential equations by differentiation

Boris Doubrov

Belarusian State University

Geometric and Algebraic Aspects of Integrability Durham, 28/07/2016

Collaborators: Andreas Čap, Dennis The

#### Outline

#### History

Cartan's C-class equations Two classical examples More recent examples

Geometry of the solution spaces Wilczynski invariants Structures on the solution spaces

#### C-class equations of arbitrary order

Cartan connections for ODE of arbitrary order Understanding the curvature tensor C-class equations in general case

#### C-class equations

E. Cartan, Les esapces généralisés et l'intégration de certaines classes d'équations différentielles (1938):Definition. A given class of differential equations of order n

$$\frac{d^{n}y}{dx^{n}} = F\left(x, y, \frac{dy}{dx}, \dots, \frac{d^{n-1}y}{dx^{n-1}}\right)$$

will be said to be a class (C) if there exists an infinite group (in the sense of Lie)  $\mathcal{G}$  transforming equations of the class into equations of the class and such that the differential invariants with respect to  $\mathcal{G}$  of an equation of the class be the first integrals of the equation.

(日) (同) (三) (三) (三) (○) (○)

#### Example 1: 2nd order ODEs

$$\frac{d^2 F_{y'}}{dx^2} - 4 \frac{d^2 F_{yy'}}{dx} - F_{y'} \frac{dF_{y'}}{dx} + 4F_{y'}F_{yy'} - 3F_yF_{y'} + 6F_y = 0.$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Here  $F_y = \partial_y F$ ,  $F_{y'} = \partial_{y'} F$  and  $\frac{d}{dx} = \partial_x + y' \partial_y + F \partial_{y'}$ .

#### Example 1: 2nd order ODEs

$$\frac{d^2 F_{y'}}{dx^2} - 4 \frac{d^2 F_{yy'}}{dx} - F_{y'} \frac{dF_{y'}}{dx} + 4F_{y'}F_{yy'} - 3F_yF_{y'} + 6F_y = 0.$$

Here  $F_y = \partial_y F$ ,  $F_{y'} = \partial_{y'} F$  and  $\frac{d}{dx} = \partial_x + y' \partial_y + F \partial_{y'}$ .

► The pseudogroup *G* consists of all point transformations:

$$(x,y)\mapsto (A(x,y),B(x,y)).$$

It preserves the class of 2nd order ODEs satisfying the above condition.

#### Example 1: 2nd order ODEs

$$\frac{d^2 F_{y'}}{dx^2} - 4 \frac{d^2 F_{yy'}}{dx} - F_{y'} \frac{dF_{y'}}{dx} + 4F_{y'}F_{yy'} - 3F_yF_{y'} + 6F_y = 0.$$

Here  $F_y = \partial_y F$ ,  $F_{y'} = \partial_{y'} F$  and  $\frac{d}{dx} = \partial_x + y' \partial_y + F \partial_{y'}$ .

The pseudogroup G consists of all point transformations:

$$(x,y)\mapsto (A(x,y),B(x,y)).$$

It preserves the class of 2nd order ODEs satisfying the above condition.

It can be shown that any equation from this class, except for thouse equivalent to y'' = 0, can be integrated by means the operation of differentiation and at most two quadratures.

Any affine connection {Γ<sup>i</sup><sub>jk</sub>} on the plane (x, y) defines an equation on *unparametrized* geodesics:

 $y'' = -\Gamma_{11}^2(y')^3 + (\Gamma_{22}^2 - 2\Gamma_{12}^1)(y')^2 + (2\Gamma_{12}^2 - \Gamma_{11}^1)y' + \Gamma_{11}^2.$ 

Any affine connection {Γ<sup>i</sup><sub>jk</sub>} on the plane (x, y) defines an equation on *unparametrized* geodesics:

 $y'' = -\Gamma_{11}^2(y')^3 + (\Gamma_{22}^2 - 2\Gamma_{12}^1)(y')^2 + (2\Gamma_{12}^2 - \Gamma_{11}^1)y' + \Gamma_{11}^2.$ 

Two affine connections are said to be projectively equivalent if their unparametrized geodesics coincide. A class of projectively equivalent affine connections is called *a projective* connection.

Any affine connection {Γ<sup>i</sup><sub>jk</sub>} on the plane (x, y) defines an equation on *unparametrized* geodesics:

 $y'' = -\Gamma_{11}^2(y')^3 + (\Gamma_{22}^2 - 2\Gamma_{12}^1)(y')^2 + (2\Gamma_{12}^2 - \Gamma_{11}^1)y' + \Gamma_{11}^2.$ 

- Two affine connections are said to be projectively equivalent if their unparametrized geodesics coincide. A class of projectively equivalent affine connections is called *a projective* connection.
- There is a 1-1 correspondence between projective connections and second order ODEs:

$$y'' = A_0(x, y) + A_1(x, y)y' + A_2(x, y)(y')^2 + A_3(x, y)(y')^3.$$

Any affine connection {Γ<sup>i</sup><sub>jk</sub>} on the plane (x, y) defines an equation on *unparametrized* geodesics:

 $y'' = -\Gamma_{11}^2(y')^3 + (\Gamma_{22}^2 - 2\Gamma_{12}^1)(y')^2 + (2\Gamma_{12}^2 - \Gamma_{11}^1)y' + \Gamma_{11}^2.$ 

- Two affine connections are said to be projectively equivalent if their unparametrized geodesics coincide. A class of projectively equivalent affine connections is called *a projective* connection.
- There is a 1-1 correspondence between projective connections and second order ODEs:

$$y'' = A_0(x, y) + A_1(x, y)y' + A_2(x, y)(y')^2 + A_3(x, y)(y')^3.$$

Generic solution of a 2nd order ODE depends on two constants of integration:

$$g(x, y, a, b) = 0.$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

 Generic solution of a 2nd order ODE depends on two constants of integration:

$$g(x,y,a,b)=0.$$

This relation defines also a two-dimensional family of curves on the parameter plane (a, b), where (x, y) serve as parameters.

 Generic solution of a 2nd order ODE depends on two constants of integration:

g(x, y, a, b) = 0.

- This relation defines also a two-dimensional family of curves on the parameter plane (a, b), where (x, y) serve as parameters.
- Define these curves as, for example, graphs of functions b = b(a), differentiate the above relation two times by a and exclude "parameters" x and y. We get so-called *dual* 2nd order ODE:

$$b''=G(a,b,b').$$

It is defined modulo point transofrmations in the (a, b) space.

 Generic solution of a 2nd order ODE depends on two constants of integration:

g(x, y, a, b) = 0.

- This relation defines also a two-dimensional family of curves on the parameter plane (a, b), where (x, y) serve as parameters.
- Define these curves as, for example, graphs of functions b = b(a), differentiate the above relation two times by a and exclude "parameters" x and y. We get so-called *dual* 2nd order ODE:

$$b''=G(a,b,b').$$

It is defined modulo point transofrmations in the (a, b) space.
It can be shown that the dual equation is cubic with respect to b' if and only if the initial equation y" = F(x, y, y') satisfies:

$$\frac{d^2 F_{y'}}{dx^2} - 4 \frac{d^2 F_{yy'}}{dx} - F_{y'} \frac{dF_{y'}}{dx} + 4F_{y'}F_{yy'} - 3F_yF_{y'} + 6F_y = 0.$$

Consider any 2nd order ODE satisfying the above relation, so that its dual equation is qubic with respect to b'.

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

- Consider any 2nd order ODE satisfying the above relation, so that its dual equation is qubic with respect to b'.
- Then we get a natural projective connection on the parameter space (a, b). Any invariants of this connection are functions on (a, b) space and, thus, are first integrals of the initial equation.

- Consider any 2nd order ODE satisfying the above relation, so that its dual equation is qubic with respect to b'.
- Then we get a natural projective connection on the parameter space (a, b). Any invariants of this connection are functions on (a, b) space and, thus, are first integrals of the initial equation.
- (Weyl) These invariants can be found as functions of the curvature of the projective connection and its covariant derivatives. This is tedious, but very explicit and constructive process.

- Consider any 2nd order ODE satisfying the above relation, so that its dual equation is qubic with respect to b'.
- Then we get a natural projective connection on the parameter space (a, b). Any invariants of this connection are functions on (a, b) space and, thus, are first integrals of the initial equation.
- (Weyl) These invariants can be found as functions of the curvature of the projective connection and its covariant derivatives. This is tedious, but very explicit and constructive process.
- However, one needs to integrate the equation first to find its dual equations and the corresponding projective connection. Can we do better and construct this projective connection without integrating the equation?

## Cartan connections for general 2nd order ODEs

► Given an arbitrary 2nd order ODE y'' = F(x, y, y'), Élie Cartan constructs a so-called *Cartan connection* on the space (x, y, z = y') of contact elements on on the plane (= jet space J<sup>1</sup>(ℝ, ℝ)).

# Cartan connections for general 2nd order ODEs

- ► Given an arbitrary 2nd order ODE y" = F(x, y, y'), Élie Cartan constructs a so-called *Cartan connection* on the space (x, y, z = y') of contact elements on on the plane (= jet space J<sup>1</sup>(ℝ, ℝ)).
- ► Cartan connection is modelled by a homogeneous space PSL(3, ℝ)/B and consists of the following data:
  - 1. principal *B*-bundle  $\pi \colon \mathcal{G} \to J^1(\mathbb{R}, \mathbb{R})$ ;
  - 2. 1-form  $\omega : T\mathcal{G} \to \mathfrak{sl}(3,\mathbb{R})$  that satisfies properties similar to the Maurer-Cartan form on the Lie group  $PSL(3,\mathbb{R})$ .
  - 3. the form  $d\omega + [\omega, \omega]$  is zero only on vertical vector fields and forms the curvature tensor  $\Omega$  of the Cartan connection.

# Cartan connections for general 2nd order ODEs

- ► Given an arbitrary 2nd order ODE y'' = F(x, y, y'), Élie Cartan constructs a so-called *Cartan connection* on the space (x, y, z = y') of contact elements on on the plane (= jet space J<sup>1</sup>(ℝ, ℝ)).
- ► Cartan connection is modelled by a homogeneous space PSL(3, ℝ)/B and consists of the following data:
  - 1. principal *B*-bundle  $\pi \colon \mathcal{G} \to J^1(\mathbb{R}, \mathbb{R})$ ;
  - 2. 1-form  $\omega \colon T\mathcal{G} \to \mathfrak{sl}(3,\mathbb{R})$  that satisfies properties similar to the Maurer-Cartan form on the Lie group  $PSL(3,\mathbb{R})$ .
  - 3. the form  $d\omega + [\omega, \omega]$  is zero only on vertical vector fields and forms the curvature tensor  $\Omega$  of the Cartan connection.
- The construction of this Cartan connection for a given 2nd order ODE is very explicit. All components ω<sub>ij</sub>, i, j = 1,..., 3 of the connection form ω are expressed explicitly in terms of the function F(x, y, z) and its partial derivatives.

► The connection descends to the parameter space (=space of solutions) if and only if the curvature of this connection vanishes identically on the direction ∂<sub>x</sub> + z∂<sub>y</sub> + F∂<sub>z</sub> along the solutions of the equation. This turns to be equivalent to the above "complicated" condition on F(x, y, z).

- ► The connection descends to the parameter space (=space of solutions) if and only if the curvature of this connection vanishes identically on the direction ∂<sub>x</sub> + z∂<sub>y</sub> + F∂<sub>z</sub> along the solutions of the equation. This turns to be equivalent to the above "complicated" condition on F(x, y, z).
- Thus, we are able to construct the projective connection on the solution space without integrating the equation.

- ► The connection descends to the parameter space (=space of solutions) if and only if the curvature of this connection vanishes identically on the direction ∂<sub>x</sub> + z∂<sub>y</sub> + F∂<sub>z</sub> along the solutions of the equation. This turns to be equivalent to the above "complicated" condition on F(x, y, z).
- Thus, we are able to construct the projective connection on the solution space without integrating the equation.

What if this projective connection on the solution space possesses no non-trivial invariants (is flat)?

- ► The connection descends to the parameter space (=space of solutions) if and only if the curvature of this connection vanishes identically on the direction ∂<sub>x</sub> + z∂<sub>y</sub> + F∂<sub>z</sub> along the solutions of the equation. This turns to be equivalent to the above "complicated" condition on F(x, y, z).
- Thus, we are able to construct the projective connection on the solution space without integrating the equation.
- What if this projective connection on the solution space possesses no non-trivial invariants (is flat)?
- In this case we end up with a Cartan connection with vanishing curvature. This happens if and only if the initial 2nd order ODE is *trivializable* (=equivalent to the trivial equation y" = 0).

- ► The connection descends to the parameter space (=space of solutions) if and only if the curvature of this connection vanishes identically on the direction ∂<sub>x</sub> + z∂<sub>y</sub> + F∂<sub>z</sub> along the solutions of the equation. This turns to be equivalent to the above "complicated" condition on F(x, y, z).
- Thus, we are able to construct the projective connection on the solution space without integrating the equation.
- What if this projective connection on the solution space possesses no non-trivial invariants (is flat)?
- In this case we end up with a Cartan connection with vanishing curvature. This happens if and only if the initial 2nd order ODE is *trivializable* (=equivalent to the trivial equation y" = 0).
- Integrating the equation is then equivalent to finding this trivialization transformation. And this is equivalent to integrating the form ω itself, i.e. constructing the map f: G → PSL(3, ℝ) such that f<sup>-1</sup>df = ω. Such problems are known are Lie type equations.

# Example 2: 3rd order ODEs

$$\frac{d^2 F_{y''}}{dx^2} - 2F_{y''}\frac{dF_{y''}}{dx} - 3\frac{dF_{y'}}{dx} + \frac{4}{9}(F_{y''})^3 + 3F_{y'}F_{y''} + 6F_y = 0.$$

#### Example 2: 3rd order ODEs

$$\frac{d^2 F_{y''}}{dx^2} - 2F_{y''}\frac{dF_{y''}}{dx} - 3\frac{dF_{y'}}{dx} + \frac{4}{9}(F_{y''})^3 + 3F_{y'}F_{y''} + 6F_y = 0.$$

► The pseudogroup *G* consists of all contact transformations:

$$(x, y, y') \mapsto (A, B, C), \quad dB - C dA = \lambda (dy - y' dx).$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

It preserves the class of 3rd order ODEs under the above restriction.

#### Example 2: 3rd order ODEs

$$\frac{d^2 F_{y''}}{dx^2} - 2F_{y''}\frac{dF_{y''}}{dx} - 3\frac{dF_{y'}}{dx} + \frac{4}{9}(F_{y''})^3 + 3F_{y'}F_{y''} + 6F_y = 0.$$

► The pseudogroup *G* consists of all contact transformations:

$$(x, y, y') \mapsto (A, B, C), \quad dB - C dA = \lambda (dy - y' dx).$$

It preserves the class of 3rd order ODEs under the above restriction.

There is a natural conformal structure on the (3-dimensional) solution space of ODEs of a given class.

#### More recent examples

Scalar ODEs of 4th order (R. Bryant, 1991): under certain (explicit) conditions on the right hand side of the equation y<sup>IV</sup> = F(x, y, y', y'', y''') the solution space carries a natural torsion-free affine connection with GL(2, ℝ) holonomy.

#### More recent examples

- Scalar ODEs of 4th order (R. Bryant, 1991): under certain (explicit) conditions on the right hand side of the equation y<sup>IV</sup> = F(x, y, y', y'', y''') the solution space carries a natural torsion-free affine connection with GL(2, ℝ) holonomy.
- Systems of *m* ODEs of 2nd order (D. Grossman, 2000): if the associated Cartan connection has vanishing torsion, the solution space carries a natural Segre (or Grassmanian) structure defined as a decomposition of the tangent space as a tensor product of ℝ<sup>2</sup> ⊗ ℝ<sup>m</sup>. If *m* = 2, this is equivalent to the conformal structure of split signature (2, 2).

#### More recent examples

- Scalar ODEs of 4th order (R. Bryant, 1991): under certain (explicit) conditions on the right hand side of the equation y<sup>IV</sup> = F(x, y, y', y'', y''') the solution space carries a natural torsion-free affine connection with GL(2, ℝ) holonomy.
- Systems of *m* ODEs of 2nd order (D. Grossman, 2000): if the associated Cartan connection has vanishing torsion, the solution space carries a natural Segre (or Grassmanian) structure defined as a decomposition of the tangent space as a tensor product of ℝ<sup>2</sup> ⊗ ℝ<sup>m</sup>. If *m* = 2, this is equivalent to the conformal structure of split signature (2, 2).
- In both cases it is shown that the natural Cartan connections for these classes of equations descend to the solution spaces and can be used to construct first integrals and integrate the equation.

#### Questions

Is there an analog of projective structure on the solution space for higher order ODEs or systems of ODEs?

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

#### Questions

- Is there an analog of projective structure on the solution space for higher order ODEs or systems of ODEs?
- What are the conditions on the equation which guarantee that such structures exist?

#### Questions

- Is there an analog of projective structure on the solution space for higher order ODEs or systems of ODEs?
- What are the conditions on the equation which guarantee that such structures exist?
- Is there a way to construct connections for these structures without integrating an equation?

# Generalized Wilczynski invariants

• Consider a linear system on  $y(x) \in \mathbb{R}^m$ :

$$y^{(k)} + P_{k-1}(x)y^{(k-1)} + \dots + P_0(x)y(x) = 0$$

up to transformations  $(x, y) \mapsto (\lambda(x), \mu(x)y)$ ,  $\mu(x) \in GL(m)$ .

< ロ > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

## Generalized Wilczynski invariants

• Consider a linear system on  $y(x) \in \mathbb{R}^m$ :

$$y^{(k)} + P_{k-1}(x)y^{(k-1)} + \dots + P_0(x)y(x) = 0$$

up to transformations  $(x,y)\mapsto (\lambda(x),\mu(x)y)$ ,  $\mu(x)\in GL(m)$ .

► The canonical Laguerre-Forsyth form is defined by conditions:  $P_{k-1} = 0$  and tr  $P_{k-2} = 0$ .

#### Generalized Wilczynski invariants

• Consider a linear system on  $y(x) \in \mathbb{R}^m$ :

$$y^{(k)} + P_{k-1}(x)y^{(k-1)} + \dots + P_0(x)y(x) = 0$$

up to transformations  $(x,y)\mapsto (\lambda(x),\mu(x)y)$ ,  $\mu(x)\in GL(m)$ .

- The canonical Laguerre-Forsyth form is defined by conditions:  $P_{k-1} = 0$  and tr  $P_{k-2} = 0$ .
- Then the following expressions become fundamental invariants for the class of linear equations:

$$\Theta_r = \sum_{j=1}^{r-1} (-1)^{j+1} \frac{(2r-j-1)!(k-r+j-1)!}{(r-j)!(j-1)!} P_{k-r+j-1}^{(j-1)},$$

for r = 2, ..., k.

## Generalized Wilczynski invariants

• Consider a linear system on  $y(x) \in \mathbb{R}^m$ :

$$y^{(k)} + P_{k-1}(x)y^{(k-1)} + \dots + P_0(x)y(x) = 0$$

up to transformations  $(x,y)\mapsto (\lambda(x),\mu(x)y)$ ,  $\mu(x)\in GL(m)$ .

- The canonical Laguerre-Forsyth form is defined by conditions:  $P_{k-1} = 0$  and tr  $P_{k-2} = 0$ .
- Then the following expressions become fundamental invariants for the class of linear equations:

$$\Theta_r = \sum_{j=1}^{r-1} (-1)^{j+1} \frac{(2r-j-1)!(k-r+j-1)!}{(r-j)!(j-1)!} P_{k-r+j-1}^{(j-1)},$$

for r = 2, ..., k.

Generalized Wilczynski invariants W<sub>r</sub>, r = 2,..., k for a non-linear system are defined as invariants Θ<sub>r</sub> evaluated at the linearization of the system.

## Structures on the solution spaces for scalar ODEs

It is known (Wilczynski, Se-ashi) that a linear system of ODEs is equivalent to the trivial one if and only if all its Wilczynski invariants vanish identically.

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

## Structures on the solution spaces for scalar ODEs

- It is known (Wilczynski, Se-ashi) that a linear system of ODEs is equivalent to the trivial one if and only if all its Wilczynski invariants vanish identically.
- It is no longer true for non-linear ODEs, as they possess other invariants independent of the generalized Wilczynski ones. However, ODEs with vanishing generalized Wilczynski invariants carry a very special geometric structure on the solution space S.

## Structures on the solution spaces for scalar ODEs

- It is known (Wilczynski, Se-ashi) that a linear system of ODEs is equivalent to the trivial one if and only if all its Wilczynski invariants vanish identically.
- It is no longer true for non-linear ODEs, as they possess other invariants independent of the generalized Wilczynski ones. However, ODEs with vanishing generalized Wilczynski invariants carry a very special geometric structure on the solution space S.
- ► In case of a scalar ODE we get a GL<sub>2</sub> structure on the solution space iff W<sub>3</sub> = ··· = W<sub>k</sub> = 0. The structure can be defined as a rational curve P<sup>1</sup> embedded into each (projectivized) tangent space:

$$[1:t:\cdots:t^{k-1}].$$

In other words, the tangent space  $T\gamma S$  of the solution space at a "point"  $\gamma$  is identified with an irreducible representation  $V_{k-1}$  of the  $\mathfrak{sl}_2$ .

In case of systems of *n* ODEs we get a GL<sub>m</sub> ⊗ SL<sub>2</sub> structure iff W<sub>2</sub> = W<sub>3</sub> = ··· = W<sub>k</sub> = 0. The structure can be defined as a projective variety P<sup>1</sup> × P<sup>m-1</sup> embedded into each (projectivized) tangent space:

$$[z_1:t\,z_1:\cdots:t^{k-1}z_1: \quad z_2:t\,z_2:\cdots:t^{k-1}z_2:\cdots z_m:t\,z_n:\cdots:t^{k-1}z_m].$$

The tangent space  $T\gamma S$  is identified with  $V_{k-1} \otimes \mathbb{R}^m$ .

In case of systems of *n* ODEs we get a GL<sub>m</sub> ⊗ SL<sub>2</sub> structure iff W<sub>2</sub> = W<sub>3</sub> = ··· = W<sub>k</sub> = 0. The structure can be defined as a projective variety P<sup>1</sup> × P<sup>m-1</sup> embedded into each (projectivized) tangent space:

$$[z_1:t z_1:\cdots:t^{k-1}z_1: z_2:t z_2:\cdots:t^{k-1}z_2:\cdots z_m:t z_n:\cdots:t^{k-1}z_m].$$

The tangent space  $T\gamma S$  is identified with  $V_{k-1} \otimes \mathbb{R}^m$ .

 All these structures admit natural connections and have a families of invariants. In fact, these structures always come with additional properties.

In case of systems of *n* ODEs we get a GL<sub>m</sub> ⊗ SL<sub>2</sub> structure iff W<sub>2</sub> = W<sub>3</sub> = ··· = W<sub>k</sub> = 0. The structure can be defined as a projective variety P<sup>1</sup> × P<sup>m-1</sup> embedded into each (projectivized) tangent space:

$$[z_1:t z_1:\cdots:t^{k-1}z_1: z_2:t z_2:\cdots:t^{k-1}z_2:\cdots z_m:t z_n:\cdots:t^{k-1}z_m].$$

The tangent space  $T\gamma S$  is identified with  $V_{k-1} \otimes \mathbb{R}^m$ .

- All these structures admit natural connections and have a families of invariants. In fact, these structures always come with additional properties.
- For scalar 3rd order ODEs we get conformal structures on 3-dimensional manifolds equipped with an Einstein-Weyl structure.

In case of systems of *n* ODEs we get a GL<sub>m</sub> ⊗ SL<sub>2</sub> structure iff W<sub>2</sub> = W<sub>3</sub> = ··· = W<sub>k</sub> = 0. The structure can be defined as a projective variety P<sup>1</sup> × P<sup>m-1</sup> embedded into each (projectivized) tangent space:

$$[z_1:t z_1:\cdots:t^{k-1}z_1: z_2:t z_2:\cdots:t^{k-1}z_2:\cdots z_m:t z_n:\cdots:t^{k-1}z_m].$$

The tangent space  $T\gamma S$  is identified with  $V_{k-1} \otimes \mathbb{R}^m$ .

- All these structures admit natural connections and have a families of invariants. In fact, these structures always come with additional properties.
- For scalar 3rd order ODEs we get conformal structures on 3-dimensional manifolds equipped with an Einstein-Weyl structure.
- For systems of two equations of 2nd order we get an ASD conformal structure on a 4-dimensional manifold.

• Generic system of *m* equations of order *k*:

$$y_i(x)^{(k)} = F_i(x, y_j^{(l)}), \ 1 \le i, j \le m; \ 0 \le l \le k-1.$$

• Generic system of *m* equations of order *k*:

$$y_i(x)^{(k)} = F_i(x, y_j^{(l)}), \ 1 \le i, j \le m; \ 0 \le l \le k-1.$$

Geometrically it is a section σ: J<sup>k-1</sup>(ℝ, ℝ<sup>m</sup>) → J<sup>k</sup>(ℝ, ℝ<sup>m</sup>) or just a submanifold ε ⊂ J<sup>k</sup> with locally diffeomorphic projection to J<sup>k-1</sup>.

• Generic system of *m* equations of order *k*:

$$y_i(x)^{(k)} = F_i(x, y_j^{(l)}), \ 1 \le i, j \le m; \ 0 \le l \le k-1.$$

- Geometrically it is a section σ: J<sup>k-1</sup>(ℝ, ℝ<sup>m</sup>) → J<sup>k</sup>(ℝ, ℝ<sup>m</sup>) or just a submanifold ε ⊂ J<sup>k</sup> with locally diffeomorphic projection to J<sup>k-1</sup>.
- $\blacktriangleright$  As a pseudogroup  ${\cal G}$  we take
  - contact transformations for scalar ODEs ( $m = 1, k \ge 3$ )
  - point transformations for systems of ODEs ( $m \ge 2, k \ge 2$ ).

• Generic system of *m* equations of order *k*:

$$y_i(x)^{(k)} = F_i(x, y_j^{(l)}), \ 1 \le i, j \le m; \ 0 \le l \le k-1.$$

- Geometrically it is a section σ: J<sup>k-1</sup>(ℝ, ℝ<sup>m</sup>) → J<sup>k</sup>(ℝ, ℝ<sup>m</sup>) or just a submanifold ε ⊂ J<sup>k</sup> with locally diffeomorphic projection to J<sup>k-1</sup>.
- $\blacktriangleright$  As a pseudogroup  ${\cal G}$  we take
  - contact transformations for scalar ODEs ( $m = 1, k \ge 3$ )
  - point transformations for systems of ODEs ( $m \ge 2, k \ge 2$ ).
- There are two natural foliations on *E*:
  - the foliation on solutions lifted to J<sup>k</sup>. Its tangent direction is given by total derivative.

(日) (同) (三) (三) (三) (○) (○)

• "vertical" foliation of fibers of projection  $\pi: \mathcal{E} \to J^{k-2}$ .

• Generic system of *m* equations of order *k*:

$$y_i(x)^{(k)} = F_i(x, y_j^{(l)}), \ 1 \le i, j \le m; \ 0 \le l \le k-1.$$

- Geometrically it is a section σ: J<sup>k-1</sup>(ℝ, ℝ<sup>m</sup>) → J<sup>k</sup>(ℝ, ℝ<sup>m</sup>) or just a submanifold ε ⊂ J<sup>k</sup> with locally diffeomorphic projection to J<sup>k-1</sup>.
- $\blacktriangleright$  As a pseudogroup  ${\cal G}$  we take
  - contact transformations for scalar ODEs ( $m = 1, k \ge 3$ )
  - point transformations for systems of ODEs ( $m \ge 2, k \ge 2$ ).
- There are two natural foliations on *E*:
  - the foliation on solutions lifted to J<sup>k</sup>. Its tangent direction is given by total derivative.
  - "vertical" foliation of fibers of projection  $\pi: \mathcal{E} \to J^{k-2}$ .
- This double fibration completely determines the extrinsic geometry of ODEs. Hence, all invariants of ODEs are exactly the invariants of this double fibration.

 Double fibrations defined by systems of ODEs perfectly fit into notions of nilpotent differential geometry (Tanaka, pseudo-product structures).

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

- Double fibrations defined by systems of ODEs perfectly fit into notions of nilpotent differential geometry (Tanaka, pseudo-product structures).
- In particular, we get the characteristic Cartan connection associated with any system of ODEs E:

$$\pi\colon \mathcal{G}\to \mathcal{E}, \quad \omega\colon T\mathcal{G}\to \mathfrak{g}.$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

- Double fibrations defined by systems of ODEs perfectly fit into notions of nilpotent differential geometry (Tanaka, pseudo-product structures).
- In particular, we get the characteristic Cartan connection associated with any system of ODEs E:

$$\pi\colon \mathcal{G}\to \mathcal{E}, \quad \omega\colon T\mathcal{G}\to \mathfrak{g}.$$

For m = 1, k ≥ 4, or m ≥ 2, k ≥ 3 this Cartan connection is no longer modelled by parabolic homogeneous spaces.

- Double fibrations defined by systems of ODEs perfectly fit into notions of nilpotent differential geometry (Tanaka, pseudo-product structures).
- In particular, we get the characteristic Cartan connection associated with any system of ODEs E:

$$\pi\colon \mathcal{G}\to \mathcal{E}, \quad \omega\colon T\mathcal{G}\to \mathfrak{g}.$$

- For m = 1, k ≥ 4, or m ≥ 2, k ≥ 3 this Cartan connection is no longer modelled by parabolic homogeneous spaces.
- In these cases the symbol algebra g (=symmetry algebra of the trivial system) is a semi-direct product of sl(2) × gl(m) and an irreducible representation V = V<sub>k-1</sub> ⊗ ℝ<sup>m</sup>.

► All invariants of the ODEs under the action of the pseudogroup G ⇔ invariants of the corresponding double fibration ⇔ invariants of the associated Cartan connection

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

► All invariants of the ODEs under the action of the pseudogroup G ⇔ invariants of the corresponding double fibration ⇔ invariants of the associated Cartan connection

- All invariants of the Cartan connection consist of the coefficients of
  - its curvature tensor  $\Omega = d\omega + 1/2[\omega, \omega]$ ;
  - its total derivatives.

- ► All invariants of the ODEs under the action of the pseudogroup G ⇔ invariants of the corresponding double fibration ⇔ invariants of the associated Cartan connection
- All invariants of the Cartan connection consist of the coefficients of
  - its curvature tensor  $\Omega = d\omega + 1/2[\omega, \omega]$ ;
  - its total derivatives.
- Similar to the Weyl tensor for projective connections, only a part of the curvature tensor generates the differential algebra of all invariants.

- ► All invariants of the ODEs under the action of the pseudogroup G ⇔ invariants of the corresponding double fibration ⇔ invariants of the associated Cartan connection
- All invariants of the Cartan connection consist of the coefficients of
  - its curvature tensor  $\Omega = d\omega + 1/2[\omega, \omega]$ ;
  - its total derivatives.
- Similar to the Weyl tensor for projective connections, only a part of the curvature tensor generates the differential algebra of all invariants.

► This part of the curvature can be identified via a pure algebraic object: H<sup>2</sup><sub>+</sub>(𝔅<sub>−</sub>,𝔅). We call it the *fundamental invariants*.

- ► All invariants of the ODEs under the action of the pseudogroup G ⇔ invariants of the corresponding double fibration ⇔ invariants of the associated Cartan connection
- All invariants of the Cartan connection consist of the coefficients of
  - its curvature tensor  $\Omega = d\omega + 1/2[\omega, \omega]$ ;
  - its total derivatives.
- Similar to the Weyl tensor for projective connections, only a part of the curvature tensor generates the differential algebra of all invariants.
- ► This part of the curvature can be identified via a pure algebraic object: H<sup>2</sup><sub>+</sub>(𝔅<sub>−</sub>,𝔅). We call it the *fundamental invariants*.
- Generalized Wilczynski invariants form a part of the fundamental invariants. But there are others!

# C-class equations of any order

Theorem

The following classes of equations:

- ► scalar ODEs of order ≥ 3 viewed up to contact transformations;
- ► systems of ODEs of order ≥ 2 viewed up to point transformations

with vanishing generalized Wilczynski invariants form (C) classes.

▲ロト ▲帰ト ▲ヨト ▲ヨト 三日 - の々ぐ

# C-class equations of any order

#### Theorem

The following classes of equations:

- ► scalar ODEs of order ≥ 3 viewed up to contact transformations;
- ► systems of ODEs of order ≥ 2 viewed up to point transformations

with vanishing generalized Wilczynski invariants form (C) classes.

Idea of the proof. Generalized Wilczynski invariants form part of the curvature of the Cartan connection associated with a given ODE. If they vanish, we can use Bianchi identity to prove that the curvature tensor vanishes on the direction tangent to the solutions and that the connection descends to a natural connection on a solution space.

# C-class equations of any order

#### Theorem

The following classes of equations:

- ► scalar ODEs of order ≥ 3 viewed up to contact transformations;
- ► systems of ODEs of order ≥ 2 viewed up to point transformations

#### with vanishing generalized Wilczynski invariants form (C) classes.

- Idea of the proof. Generalized Wilczynski invariants form part of the curvature of the Cartan connection associated with a given ODE. If they vanish, we can use Bianchi identity to prove that the curvature tensor vanishes on the direction tangent to the solutions and that the connection descends to a natural connection on a solution space.
- However, direct use of Bianchi identities is very messy. Smart algebraic techniques coming splitting operators in parabolic geometries are required to sort this out.

## Examples

Any Einstein-Weyl structure in 3D comes from a certain 3rd order ODE with vanishing Wilczynski (Wünschmann) invariant. Similarly, any pair of 2nd order ODEs with vanishing Wilczynski invariant comes from a ASD conformal structure.

・ロト・日本・モート モー うへぐ

## Examples

- Any Einstein-Weyl structure in 3D comes from a certain 3rd order ODE with vanishing Wilczynski (Wünschmann) invariant. Similarly, any pair of 2nd order ODEs with vanishing Wilczynski invariant comes from a ASD conformal structure.
- Universal example (Hitchin, LeBrun, Bryant). Take any rational curve  $P^1$  in an (m + 1)-dimensional complex manifold M with a normal bundle mO(k 1). Then the complete deformation family of this rational curve will form a solution space of a system of m ODEs of order k with vanishing Wilczynski invariants.

## Examples

- Any Einstein-Weyl structure in 3D comes from a certain 3rd order ODE with vanishing Wilczynski (Wünschmann) invariant. Similarly, any pair of 2nd order ODEs with vanishing Wilczynski invariant comes from a ASD conformal structure.
- ► Universal example (Hitchin, LeBrun, Bryant). Take any rational curve P<sup>1</sup> in an (m + 1)-dimensional complex manifold M with a normal bundle mO(k − 1). Then the complete deformation family of this rational curve will form a solution space of a system of m ODEs of order k with vanishing Wilczynski invariants.
- Complete deformation family of a non-degenerate conic in P<sup>2</sup> is the space of all conics and is given by the following 5th order ODE:

$$9(y'')^2 y^{(5)} - 45y'' y''' y^{(4)} + 40(y''')^3 = 0.$$