

Collective coordinate approximation to the scattering of solitons in modified $(1+1)$ NLS and sine-Gordon models.

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- Quasi-integrability
- Collective coordinate approximation
- NLS model
- Modified NLS
- Modified sine-Gordon



EoM can be rewritten in terms of a zero curvature condition:

$$\partial_t A_x - \partial_x A_t + [A_x, A_t] = 0$$

Abelianization procedure...

$$\frac{dQ^{(n)}}{dt} = 0$$

IST \rightarrow exact soliton solutions



EoM can be rewritten in terms of an anomalous zero curvature condition:

$$\partial_t A_x - \partial_x A_t + [A_x, A_t] = XT_3^0$$

Abelianization procedure...

$$\partial_t a_x^{(3,-n)} - \partial_x a_t^{(3,-n)} = X\alpha^{(3,-n)} \quad n = 0, 1, 2, \dots$$

$$\Rightarrow \frac{dQ^{(n)}}{dt} = \beta_n; \quad \text{with} \quad Q^{(n)} = \int_{-\infty}^{\infty} dx a_x^{(3,-n)}$$
$$\text{where} \quad \beta_n = \int_{-\infty}^{\infty} dx X \alpha^{(3,-n)}.$$



When fields are eigenstates of the space-time parity

$$P : \quad (\tilde{x}, \tilde{t}) \rightarrow (-\tilde{x}, -\tilde{t}) \quad \tilde{x} = x - x_{\Delta} \quad \text{and} \quad \tilde{t} = t - t_{\Delta},$$

then the anomaly integrates over time to zero and the charges are asymptotically conserved

$$Q^{(n)}(t \rightarrow \infty) - Q^{(n)}(t \rightarrow -\infty) = \int_{-\infty}^{\infty} dt \beta_n = 0$$



Start with static solution $\psi(x, q_1, \dots, q_n)$ and allow the parameters to vary with time $q_i = q_i(t)$.

- Reduces the infinite-dimensional problem to a coupled set of ODEs
- but**
- Requires a sensible choice of coordinates
 - Does not account for deformations or energy radiation



(1+1) dim NLS eq:

$$i\partial_t\psi = -\partial_x^2\psi + \frac{\delta V}{\delta|\psi|^2}\psi$$

$$V_{\text{NLS}} = \eta|\psi|^4$$

1-soliton solution ($\eta = -1$):

$$\psi = \frac{b}{\cosh[b(x - vt - x_0)]} e^{i\left[\left(b^2 - \frac{v^2}{4}\right)t + \frac{v}{2}x + \delta\right]}$$



$$\psi = \psi_1 + \psi_2 = \varphi_1 e^{i\theta_1} + \varphi_2 e^{i\theta_2}$$

$$\varphi_1 = \frac{a}{\cosh [b(x + \xi)]}, \quad \theta_1 = -\mu(x + \xi) + b^2 t + \lambda + \delta_1,$$

$$\varphi_2 = \frac{a}{\cosh [b(x - \xi)]}, \quad \theta_2 = \mu(x - \xi) + b^2 t + \lambda + \delta_2,$$

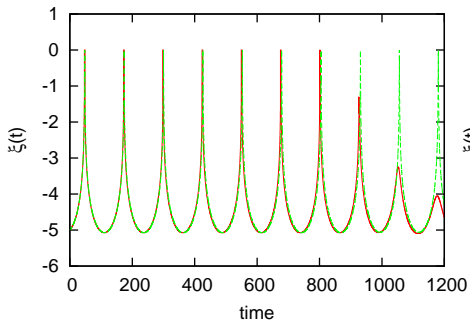
Take our collective coordinates:

$a(t) \sim$ height, $\xi(t) \sim$ position, $\mu(t) \sim$ velocity, $\lambda(t) \sim$ phase

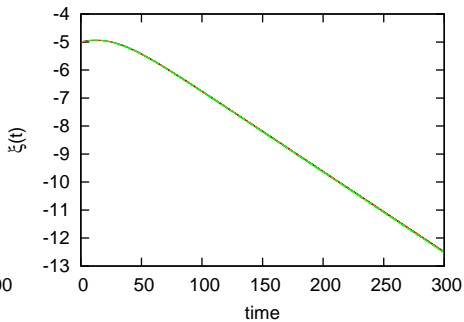
Initial phase difference:

$$\delta \equiv \delta_1 - \delta_2$$

F. Zou and J. Yan, *Chin. Phys. Lett.* **11**, 265 (1994)

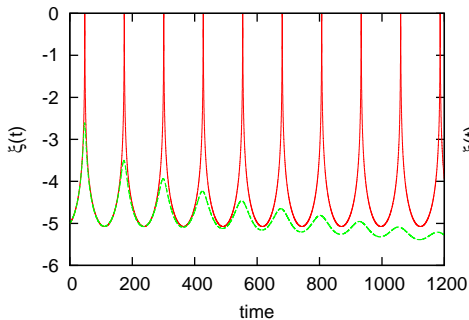


$$\delta = 0$$

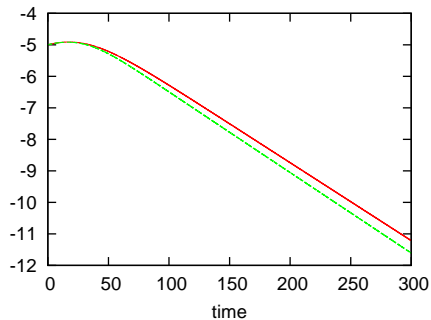


$$\delta = \pi$$

red: full simulation, green: approximation



$$\delta = \frac{\pi}{32}$$



$$\delta = \frac{3\pi}{4}$$

red: full simulation, green: approximation



$$\psi = \psi_1 + \psi_2 = \varphi_1 e^{i\theta_1} + \varphi_2 e^{i\theta_2}$$

$$\varphi_1 = \frac{a_1}{\cosh[a_1(x + \xi_1)]}, \quad \theta_1 = -\mu_1 \left(x + \frac{\xi_1}{2}\right) + a_1^2 t + \lambda_1$$

$$\varphi_2 = \frac{a_2}{\cosh[a_2(x + \xi_2)]}, \quad \theta_2 = -\mu_2 \left(x + \frac{\xi_2}{2}\right) + a_2^2 t + \lambda_2$$

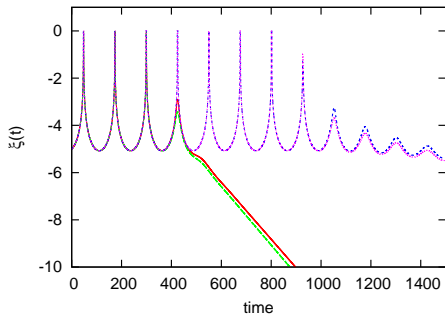
New collective coordinates:

$a_{1,2}(t) \sim$ heights, $\xi_{1,2}(t) \sim$ positions, $\mu_{1,2}(t) \sim$ velocities,

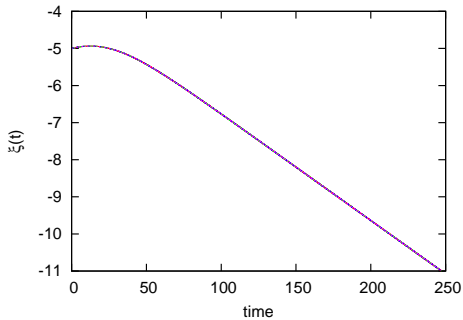
$\lambda_{1,2}(t) \sim$ phases

Phase difference:

$$\delta \equiv \lambda_1 - \lambda_2$$

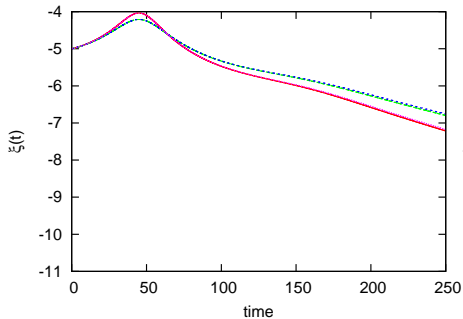


$$\delta = 0$$

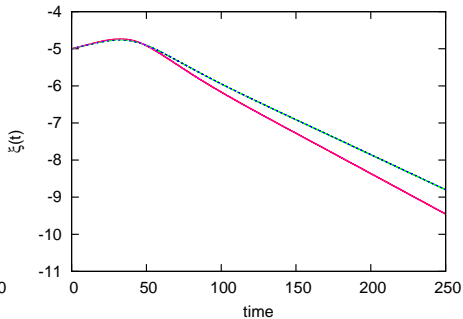


$$\delta = \pi$$

red/green: LH/RH approximation, pink/blue: LH/RH full simulation



$$\delta = \frac{\pi}{4}$$



$$\delta = \frac{\pi}{2}$$

red/green: LH/RH approximation, pink/blue: LH/RH full simulation



(1+1) dim NLS eq:

$$i\partial_t\psi = -\partial_x^2\psi + \frac{\delta V}{\delta|\psi|^2}\psi$$

$$V = \frac{2}{2+\epsilon}\eta (|\psi|^2)^{2+\epsilon}$$

1-soliton solution ($\eta < 0$):

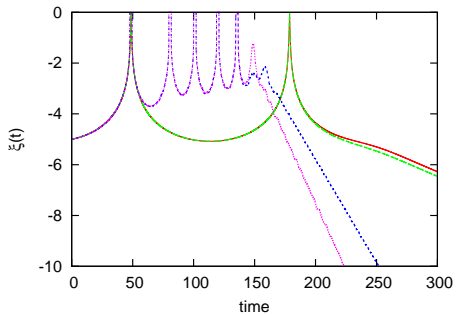
$$\psi = \left(\sqrt{\frac{2+\epsilon}{2|\eta|}} \frac{b}{\cosh[(1+\epsilon)b(x-vt-x_0)]} \right)^{\frac{1}{1+\epsilon}} e^{i\left[\left(b^2 - \frac{v^2}{4}\right)t + \frac{v}{2}x\right]}$$



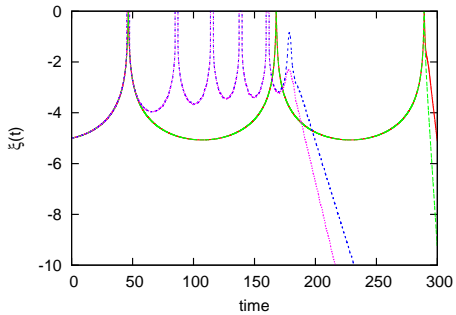
$$\psi = \psi_1 + \psi_2 = \varphi_1 e^{i\theta_1} + \varphi_2 e^{i\theta_2}$$

$$\varphi_1 = \left(\sqrt{\frac{2+\epsilon}{2}} \frac{a_1}{\cosh [(1+\epsilon) a_1 (x + \xi_1)]} \right)^{\frac{1}{1+\epsilon}}$$

$$\varphi_2 = \left(\sqrt{\frac{2+\epsilon}{2}} \frac{a_2}{\cosh [(1+\epsilon) a_2 (x + \xi_2)]} \right)^{\frac{1}{1+\epsilon}}$$

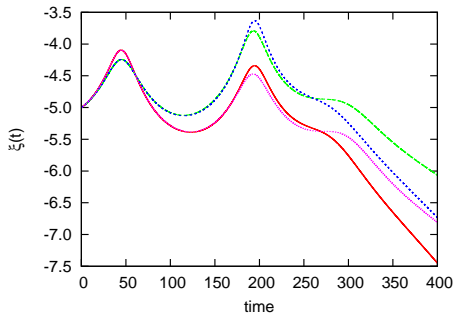


$$\epsilon = 0.06$$

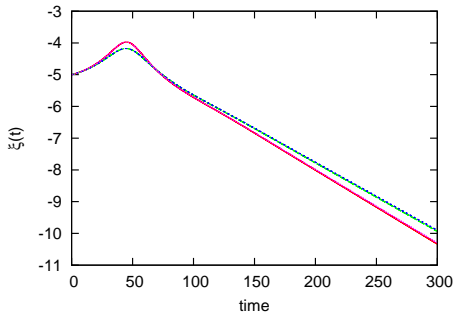


$$\epsilon = -0.06$$

red/green: LH/RH approximation, pink/blue: LH/RH full simulation

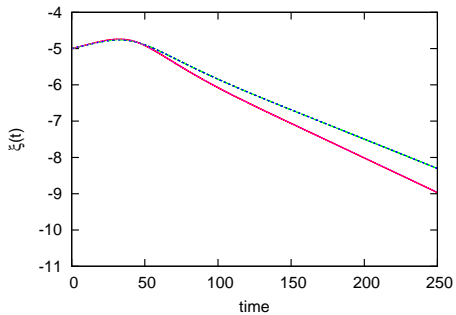


$$\epsilon = 0.06$$

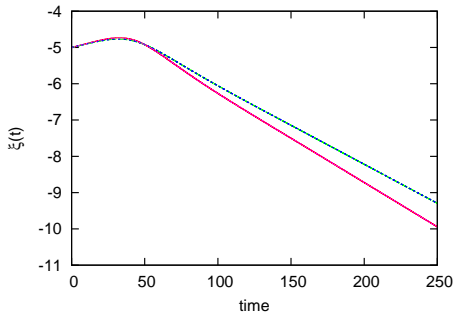


$$\epsilon = -0.06$$

red/green: LH/RH approximation, pink/blue: LH/RH full simulation

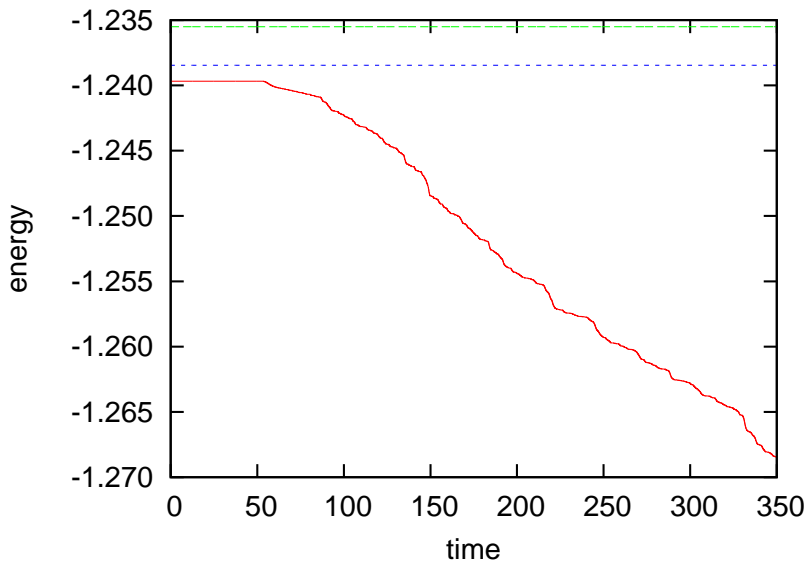


$$\epsilon = 0.06$$



$$\epsilon = -0.06$$

red/green: LH/RH approximation, pink/blue: LH/RH full simulation



red: $\delta = 0$, blue: $\delta = \frac{\pi}{4}$, green: $\delta = \frac{\pi}{2}$



Anomalous curvature condition:

$$\partial_t A_x - \partial_x A_t + [A_x, A_t] = \chi T_3^0$$

$$\chi \equiv -i \partial_x \left(\frac{\delta V}{\delta |\psi|^2} - 2\eta |\psi|^2 \right)$$

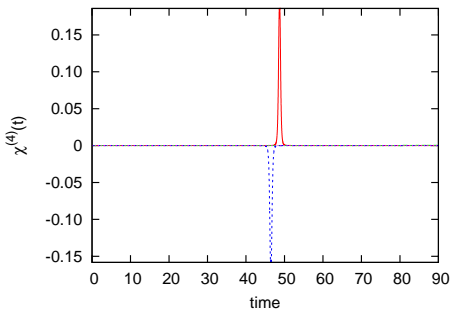
Anomalies:

$$\frac{dQ^{(n)}}{dt} = \beta_n(t)$$

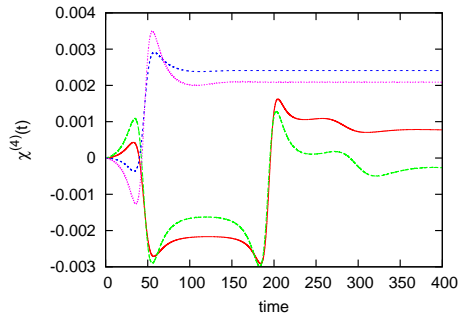
Time-integrated anomaly:

$$\chi^{(4)}(t) \equiv \int_{-\infty}^t dt' \beta_4$$

$$\text{If } P(\psi) = (-1)^n \psi^* \Rightarrow Q^{(n)}(t = +\infty) = Q^{(n)}(t = -\infty)$$

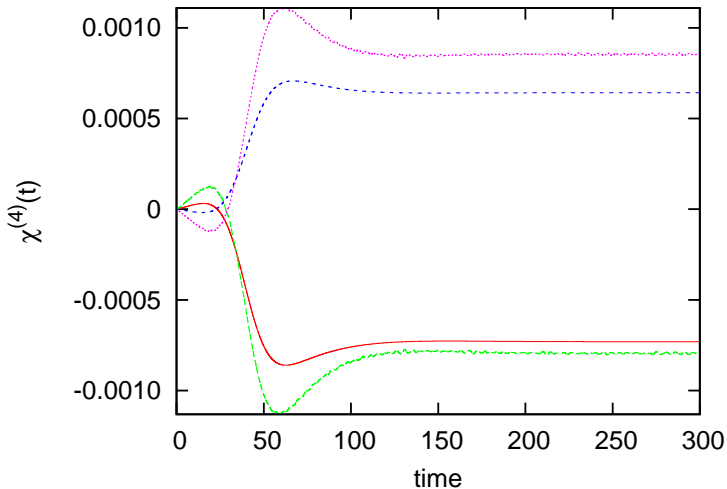


$$\delta = 0$$



$$\delta = \frac{\pi}{4}$$

red/blue: $\epsilon = \pm 0.06$ approximation, green/pink: $\epsilon = \pm 0.06$ full simulation



$$\delta = \frac{\pi}{2}$$

red/blue: $\epsilon = \pm 0.06$ approximation, green/pink: $\epsilon = \pm 0.06$ full simulation



(1+1) dim sine-Gordon eq:

$$L = \int dx \frac{1}{2} \left((\partial_t \psi)^2 - (\partial_x \psi)^2 \right) - V(\psi)$$

$$V_{SG} = \frac{1}{8} \sin^2(2\psi)$$

1-soliton solution:

$$\psi = \text{ArcTan} \left(e^{\pm(x-x_0)} \right)$$



Take a change of variable $\psi \rightarrow \phi$

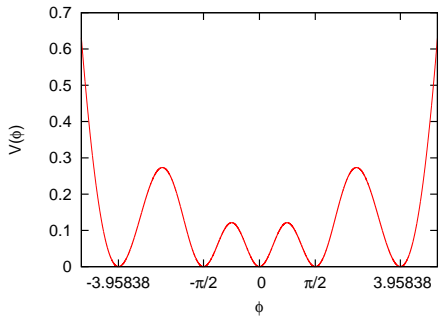
$$\psi(\phi) = \frac{c\phi}{\sqrt{1 + \epsilon\phi(\phi - 2\gamma)}}$$

where c is chosen to be

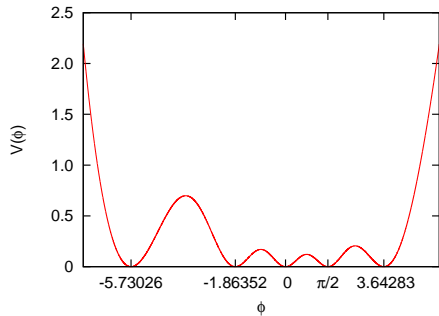
$$c = \sqrt{1 + \epsilon\pi \left(\frac{\pi}{4} - \gamma \right)}$$

then this solves the sine-Gordon equation with potential:

$$V(\phi) = \left(\frac{d\phi}{d\psi} \right)^2 V_{SG} = \frac{1}{8} \frac{(1 + \epsilon\phi(\phi - 2\gamma))^3}{c^2 (1 - \epsilon\gamma\phi)^2} \sin^2(2\psi(\phi)).$$



$$\epsilon = 0.05, \gamma = 0$$



$$\epsilon = 0.05, \gamma = 1$$

Modified sine-Gordon

Collective coordinate approximation

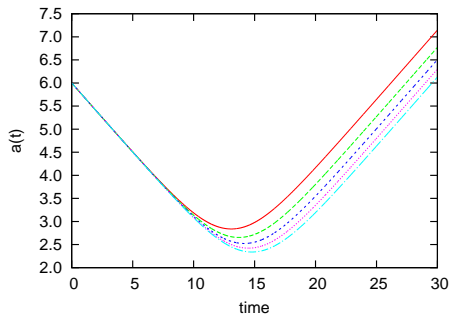


$$\psi = \psi_1 + \psi_2 = \text{ArcTan} \left(2 \text{Sinh}(x) e^{-a(t)} \right)$$

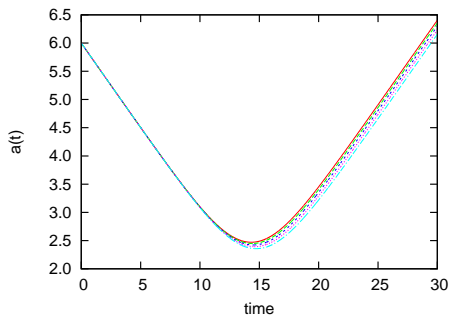
then take a change of variable $\psi \rightarrow \phi$

$$\phi = \frac{\psi^2 \epsilon \gamma + \sqrt{\psi^2 c^2 + \psi^4 \epsilon (-1 + \gamma^2 \epsilon)}}{\psi^2 \epsilon - c^2} \quad \text{for } x < 0$$

$$\phi = \frac{\psi^2 \epsilon \gamma - \sqrt{\psi^2 c^2 + \psi^4 \epsilon (-1 + \gamma^2 \epsilon)}}{\psi^2 \epsilon - c^2} \quad \text{for } x > 0$$



$\epsilon = -0.2, -0.1, 0, 0.1, 0.2$ and $\gamma = 0$



$\gamma = -0.4, -0.2, 0, 0.2, 0.4$ and $\epsilon = 0.1$



Anomalous curvature condition:

$$\partial_t A_x - \partial_x A_t + [A_x, A_t] = XT_3^0$$

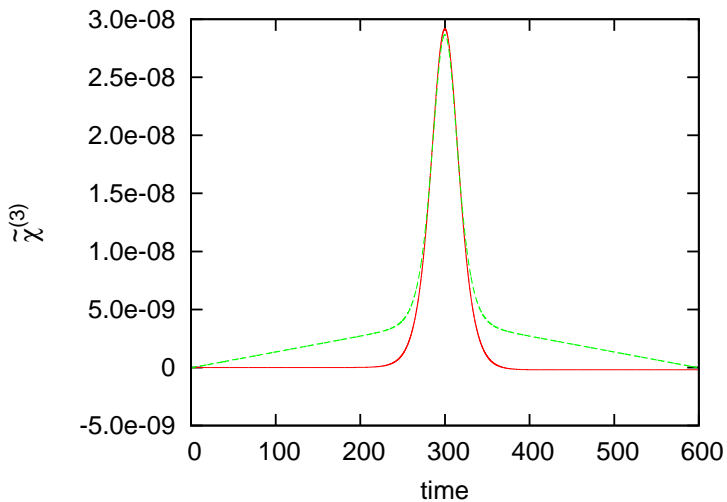
$$X \equiv \frac{iw}{2} \partial_- \phi \left[\frac{d^2 V}{d\phi^2} + w^2 V - m \right]$$

(X vanishes for sine-Gordon for $m = 1, w = 4$)

Time-integrated anomaly:

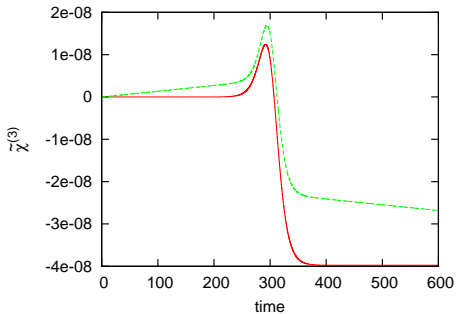
$$\chi^{(3)} = -\frac{1}{2} \int_{t_0}^t dt' \beta^{(3)}$$

If $P(\phi) = -\phi + \text{const}$ and $P(V) = V$

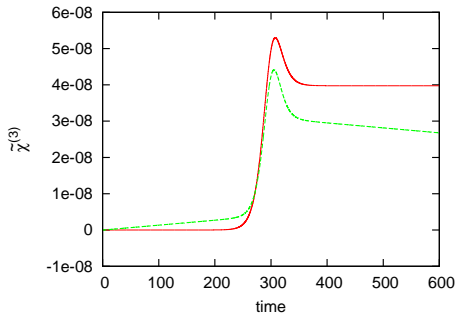


$$\epsilon = 0.000001, \gamma = 0.00001$$

red: approximation, green: full simulation

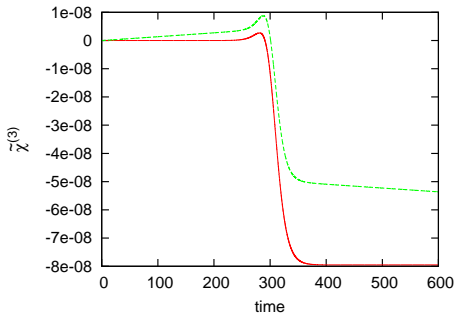


$$\epsilon = 0.000001, \gamma = 0.002$$

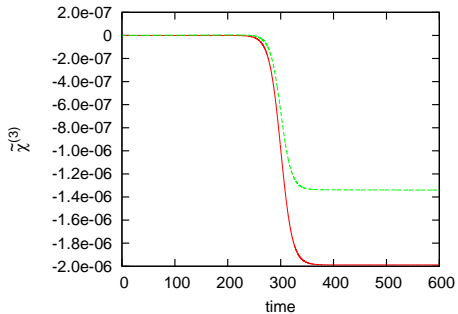


$$\epsilon = 0.000001, \gamma = -0.002$$

red: approximation, green: full simulation



$$\epsilon = 0.000001, \gamma = 0.004$$



$$\epsilon = 0.000001, \gamma = 0.1$$

red: approximation, green: full simulation



- Collective coordinate approximation generally accurate in all considered systems
- In our modified NLS approximation was less reliable the more time the solitons spent in close proximity
- In our modified sine-Gordon the values for the time-integrated anomaly are better when the fields possess suitable parity symmetry