ON THE LONG-TIME STABILITY OF THE IMPLICIT EULER SCHEME FOR THE TWO-DIMENSIONAL NAVIER–STOKES EQUATIONS*

F. TONE[†] AND D. WIROSOETISNO^{†‡}

Abstract. In this paper we study the stability for all positive time of the fully implicit Euler scheme for the two-dimensional Navier–Stokes equations. More precisely, we consider the time discretization scheme and with the aid of the discrete Gronwall lemma and the discrete uniform Gronwall lemma we prove that the numerical scheme is stable.

Key words. Navier-Stokes equations, discrete Gronwall lemmas, implicit Euler scheme

AMS subject classifications. 65M12, 76D05

DOI. 10.1137/040618527

1. Introduction. Let $\Omega \subset \mathbb{R}^2$ be an open bounded set with boundary $\partial\Omega$ of class C^2 . The Navier–Stokes equations of viscous incompressible fluids are

(1.1)
$$u_t + (u \cdot \nabla)u - \nu \Delta u + \nabla p = f,$$

$$div u = 0,$$

where $u = (u_1, u_2)$ is the velocity, p is the pressure, ν is the kinematic viscosity, and f represents body forces applied to the fluid. We complete these equations with the initial condition

$$(1.3) u(x,0) = u_0(x),$$

with $u_0: \Omega \to \mathbb{R}^2$ being given, and with the nonslip boundary condition

$$(1.4) u = 0 on \partial \Omega.$$

In the notation described below, system (1.1)–(1.4) can be written as the functional evolution equation

(1.5)
$$u_t + \nu A u + B(u, u) = f, \quad u(0) = u_0.$$

In the two-dimensional case under consideration, the solution to the Navier–Stokes equations is known to be smooth for all time (cf. [13]). The velocity u is bounded uniformly for all time by

$$(1.6) |u(t)|_{L^2(\Omega)^2}^2 \le e^{-\nu\lambda_1 t} |u_0|_{L^2(\Omega)^2}^2 + c(1 - e^{-\nu\lambda_1 t}) |f|_{L^{\infty}(\mathbb{R}_+; L^2(\Omega)^2)}^2,$$

where λ_1 is the first eigenvalue of the Stokes operator A, and we have assumed that $f \in L^{\infty}(\mathbb{R}_+; L^2(\Omega)^2)$. Furthermore, using techniques based on the uniform Gronwall

^{*}Received by the editors November 8, 2004; accepted for publication (in revised form) May 10, 2005; published electronically February 8, 2006. This work was partially supported by the National Science Foundation under grant NSF-DMS-0305110 and by the Research Fund of Indiana University. http://www.siam.org/journals/sinum/44-1/61852.html

[†]The Institute for Scientific Computing and Applied Mathematics, Indiana University, Bloomington, IN 47405-7106 (ftone@indiana.edu).

[‡]Current address: Department of Mathematical Sciences, University of Durham, Durham DH1 3LE, United Kingdom (djoko.wirosoetisno@durham.ac.uk).

lemma (cf. [12]), one can bound u uniformly in $H_0^1(\Omega)$ for all $t \geq 0$ by a function which depends on the initial condition

$$|u(t)|_{H_0^1(\Omega)^2}^2 \le K(|u_0|_{H_0^1(\Omega)^2}, |f|_{L^{\infty}(\mathbb{R}_+; L^2(\Omega)^2)}).$$

This dependence on the initial data can be dropped when one considers sufficiently large time, $t \geq T_c(|u_0|_{L^2(\Omega)^2}, |f|_{L^{\infty}(\mathbb{R}_+; L^2(\Omega)^2)})$, giving

(1.8)
$$|u(t)|_{H_0^1(\Omega)^2}^2 \le K(|f|_{L^{\infty}(\mathbb{R}_+;L^2(\Omega)^2)}) \quad \forall t \ge T_c.$$

In this paper we consider a time discretization of (1.5) using the fully implicit Euler scheme

(1.9)
$$\frac{u^n - u^{n-1}}{h} + \nu A u^n + B(u^n, u^n) = f^n, \quad u^0 = u_0,$$

where

$$(1.10) f_n = \frac{1}{\Delta t} \int_{(n-1)\Delta t}^{n\Delta t} f(t) dt,$$

and seek to obtain similar bounds on $|u^n|_{H^1_0(\Omega)^2}$.

Before we proceed further, we note that a related result for the linearized implicit Euler scheme

(1.11)
$$\frac{u^n - u^{n-1}}{k} + \nu A u^n + B(u^{n-1}, u^n) = f^n, \quad u^0 = u_0,$$

is proved in [7]. A different approach for the linearized implicit Euler scheme for the case without forcing term appears in [3].

Important background information on different computational methods can be found in some of the books and articles available in the literature. On finite elements, see, e.g., [4], [6]; on finite differences and finite elements, [9], [13]; on spectral methods, [1], [5].

For the mathematical setting of the problem, we consider the following spaces:

(1.12)
$$V = \{ v \in H_0^1(\Omega)^2, \operatorname{div} v = 0 \},$$

(1.13)
$$H = \{ v \in L^2(\Omega)^2, \operatorname{div} v = 0, v \cdot n = 0 \text{ on } \partial \Omega \},$$

where n is the unit outward normal on $\partial\Omega$. The space V is endowed with the scalar product

(1.14)
$$((u,v)) = \sum_{i,j=1}^{2} \int_{\Omega} \frac{\partial u_i}{\partial x_j}(x) \frac{\partial v_i}{\partial x_j}(x) dx$$

and with the corresponding norm

(1.15)
$$||u|| = ((u,u))^{1/2},$$

and H is endowed with the scalar product and the norm of $L^2(\Omega)^2$, denoted by (\cdot, \cdot) and $|\cdot|$.

We denote by A the linear continuous operator from V into V' such that

$$(1.16) \langle Au, v \rangle_{V',V} = ((u,v)) \quad \forall u, v \in V.$$

The domain of A in H is denoted by D(A) and, using the regularity theory for the Stokes equation (see, for instance, [13]), one can show that

$$(1.17) D(A) = H^2(\Omega)^2 \cap V.$$

We have the following inclusions:

$$(1.18) D(A) \subset V \subset H,$$

and the so-called Poincaré inequality holds true:

$$(1.19) |u| \le \frac{1}{\sqrt{\lambda_1}} ||u| \quad \forall u \in V,$$

where $\lambda_1 > 0$ is the first eigenvalue of the Stokes operator A.

As is well known, the form (1.5) of the Navier–Stokes equations was derived by Leray [8], using the weak formulation of the Navier–Stokes equations. The latter is obtained by multiplying (1.1) by a test function $v \in V$ and integrating by parts over Ω , using Green's formula, viz.,

(1.20)
$$\frac{d}{dt}(u(t),v) + \nu((u(t),v)) + b(u(t),u(t),v) = (f(t),v) \quad \forall v \in V,$$

where

(1.21)
$$b(u, v, w) = \sum_{i, j=1, 2} \int_{\Omega} u_i(x) \frac{\partial v_j}{\partial x_i}(x) w_j(x) dx.$$

The form b is trilinear continuous on $H^1(\Omega)^2$ and enjoys the following properties:

$$(1.23) |b(u, v, w)| \le c_b |u|^{1/2} ||u||^{1/2} ||v|| ||w|^{1/2} ||w||^{1/2} \forall u, v, w \in V,$$

$$(1.24) b(u, v, v) = 0 \forall u, v \in V,$$

the last equation implying

$$(1.25) b(u, v, w) = -b(u, w, v) \quad \forall u, v, w \in V.$$

Using b, we define the bilinear operator B from $V \times V$ into V' by

$$\langle B(u,v), w \rangle_{V',V} = b(u,v,w) \quad \forall u,v,w \in V.$$

For more details about the functional spaces D(A), V, and H as well as the operators A, B, and b, the reader is referred to, e.g., [2], [11], and [13].

2. H^1 stability and the main result. Throughout the paper, we assume that $f \in L^{\infty}(\mathbb{R}_+; H)$ and we set $|f|_{\infty} := |f|_{L^{\infty}(\mathbb{R}_+; H)}$. We adopt the following convention: c_i denotes constants that depend only on the parameters such as λ_1, ν , etc.; K_i depend in addition on $u(t_*)$ at some specified time t_* and on the forcing f; κ_i are bounds on the timestep k and may depend on u_0 and f.

In proving the main result, we will need a couple of preliminary lemmas. We begin with an analogue of (1.6), proved in almost the same way (see, e.g., [12, p. 109]).

LEMMA 2.1. For every k > 0, we have

$$(2.1) |u^n|^2 \le (1 + \nu \lambda_1 k)^{-n} |u_0|^2 + [1 - (1 + \nu \lambda_1 k)^{-n}] \frac{|f|_{\infty}^2}{\nu^2 \lambda_1^2} \quad \forall n \ge 0,$$

and there exists $K_1 = K_1(|u_0|, |f|_{\infty})$ such that

$$(2.2) |u^n|^2 \le K_1 \quad \forall n \ge 0,$$

and

(2.3)
$$\nu \sum_{j=i}^{n} k \|u^{j}\|^{2} \le K_{1} + (n-i+1)k \frac{|f|_{\infty}^{2}}{\nu \lambda_{1}} \quad \forall i = 1, \dots, n.$$

Proof. Taking the scalar product of (1.9) with $2ku^n$ in H and using the relation

(2.4)
$$2(\varphi - \psi, \varphi) = |\varphi|^2 - |\psi|^2 + |\varphi - \psi|^2 \quad \forall \varphi, \psi \in H,$$

and the skew property (1.24), we obtain

$$(2.5) |u^n|^2 - |u^{n-1}|^2 + |u^n - u^{n-1}|^2 + 2\nu k ||u^n||^2 = 2k(f^n, u^n).$$

Using the Cauchy–Schwarz inequality and the Poincaré inequality (1.19), we majorize the right-hand side of (2.5) by

(2.6)
$$2k|f^n||u^n| \le \frac{2k}{\sqrt{\lambda_1}}|f^n|||u^n|| \le \nu k||u^n||^2 + \frac{k}{\nu \lambda_1}|f^n|^2.$$

Relations (2.5) and (2.6) imply

$$(2.7) |u^n|^2 - |u^{n-1}|^2 + |u^n - u^{n-1}|^2 + \nu k ||u^n||^2 \le \frac{k}{\nu \lambda_1} |f^n|^2.$$

Using again the Poincaré inequality (1.19), we find from (2.7)

$$|u^n|^2 \le \frac{1}{\alpha} |u^{n-1}|^2 + \frac{k}{\alpha \nu \lambda_1} |f^n|^2,$$

where

$$(2.9) \alpha = 1 + \nu \lambda_1 k.$$

Using (2.8) recursively, we find

$$|u^{n}|^{2} \leq \frac{1}{\alpha^{n}}|u^{0}|^{2} + \frac{k}{\nu\lambda_{1}}\sum_{i=1}^{n}\frac{1}{\alpha^{i}}|f^{n+1-i}|^{2}$$

$$\leq (1+\nu\lambda_{1}k)^{-n}|u_{0}|^{2} + \frac{|f|_{\infty}^{2}}{\nu^{2}\lambda_{1}^{2}}[1-(1+\nu\lambda_{1}k)^{-n}],$$

which proves (2.1); (2.1) easily implies (2.2) with

(2.11)
$$K_1(|u_0|,|f|_{\infty}) := |u_0|^2 + \frac{1}{\nu^2 \lambda_1^2} |f|_{\infty}^2.$$

Now adding up (2.7) with n from i to m and dropping some terms, we find

(2.12)
$$\nu k \sum_{j=i}^{m} ||u^{j}||^{2} \leq |u^{i-1}|^{2} + \frac{k}{\nu \lambda_{1}} \sum_{j=i}^{m} |f^{j}|^{2}$$
$$\leq K_{1} + \frac{|f|_{\infty}^{2}}{\nu \lambda_{1}} (m - i + 1)k,$$

which is just (2.3) with n in place of m. COROLLARY 2.2. If

$$(2.13) 0 < k \le \frac{1}{\nu \lambda_1} =: \kappa_1,$$

then

$$(2.14) |u^n|^2 \le 2\rho_0^2 \quad \forall \, nk \ge T_0(|u_0|, |f|_\infty) := \frac{4}{\nu \lambda_1} \ln\left(\frac{|u_0|}{\rho_0}\right),$$

where $\rho_0 := |f|_{\infty}/(\nu \lambda_1)$.

Proof. From the bound (2.1) on $|u^n|^2$, we infer that

$$|u^n|^2 \le (1 + \nu \lambda_1 k)^{-n} |u_0|^2 + \rho_0^2$$

and using assumption (2.13) on k and the fact that $1 + x \ge \exp(x/2)$ if $x \in (0,1)$, we obtain

$$|u^n|^2 \le \exp\left(-nk\frac{\nu\lambda_1}{2}\right)|u_0|^2 + \rho_0^2.$$

For $nk \geq T_0$, the above inequality implies conclusion (2.14) of the corollary.

We now seek to obtain uniform bounds on u^n in V similar to those obtained in H (see (2.2)). To this end, we first derive bounds on a finite interval of time (see Proposition 2.5). We then repeatedly use these together with (a discrete uniform Gronwall) Lemma 2.6 on successive intervals to arrive at the desired uniform bounds.

We begin with some preliminary inequalities. Taking the scalar product of (1.9) with $2kAu^n$ in H, we obtain

(2.15)
$$||u^n||^2 - ||u^{n-1}||^2 + ||u^n - u^{n-1}||^2 + 2\nu k |Au^n|^2 + 2kb(u^n, u^n, Au^n) = 2k(f^n, Au^n).$$

Using property (1.22) of the trilinear form b and recalling (2.2), we have the following bound of the nonlinear term:

(2.16)
$$2kb(u^{n}, u^{n}, Au^{n}) \leq 2 c_{b} k |u^{n}|^{1/2} ||u^{n}|| |Au^{n}||^{3/2}$$
$$\leq \frac{\nu k}{2} |Au^{n}|^{2} + \frac{27c_{b}^{4}}{2\nu^{3}} K_{1} k ||u^{n}||^{4}.$$

We bound the right-hand side of (2.15) by Cauchy-Schwarz,

(2.17)
$$2k(f^{n}, Au^{n}) \leq 2k|f^{n}||Au^{n}|$$
$$\leq \frac{\nu k}{2}|Au^{n}|^{2} + \frac{2}{\nu}k|f^{n}|^{2}.$$

Relations (2.15)–(2.17) imply

(2.18)
$$||u^{n}||^{2} - ||u^{n-1}||^{2} + ||u^{n} - u^{n-1}||^{2} + \nu k |Au^{n}|^{2}$$

$$\leq \frac{27c_{b}^{4}}{2\nu^{3}} K_{1}k||u^{n}||^{4} + \frac{2}{\nu}k|f^{n}|^{2},$$

from which we obtain

$$(2.19) 0 \le c_2 K_1 k \|u^n\|^4 - \|u^n\|^2 + \|u^{n-1}\|^2 + c_3 k |f|_{\infty}^2,$$

where

(2.20)
$$c_2 = \frac{27c_b^4}{2\nu^3} \text{ and } c_3 = \frac{2}{\nu}.$$

Unlike (2.7), (2.19) does not (directly) provide a useful bound for $||u^n||$, so we proceed to show that (2.19) does give a proper bound for $||u^n||$ if the timestep k is sufficiently small.

LEMMA 2.3. Suppose that $0 < k \le \kappa_1$ and assume that, for some n, we have

(2.21)
$$c_2 K_1 k \left(K_2 \| u^{n-1} \|^2 + c_4 |f|_{\infty}^2 \right) \le \frac{1}{5},$$

where $K_2(|u_0|, |f|_{\infty}) = 2 + 4c_b^2 K_1/\nu^2$ and $c_4 = 4/(\nu^2 \lambda_1)$. Then (2.19) implies

$$(2.22) ||u^n||^2 \le ||u^{n-1}||^2 \left[1 + c_5 K_1 k \left(||u^{n-1}||^2 + k|f|_{\infty}^2\right)\right] + c_6 k|f|_{\infty}^2$$

for some constants c_5 and c_6 .

Proof. Relation (2.19) implies either

(2.23)
$$||u^n||^2 \le \frac{1 - \sqrt{\Delta_{n-1}}}{2c_2 K_1 k}$$

or

(2.24)
$$||u^n||^2 \ge \frac{1 + \sqrt{\Delta_{n-1}}}{2c_2 K_1 k},$$

where

(2.25)
$$\Delta_{n-1} = 1 - 4c_2 K_1 k (\|u^{n-1}\|^2 + c_3 k |f|_{\infty}^2) > 0$$
 by (2.13) and (2.21).

We now show that (2.21) excludes (2.24). Indeed, taking the scalar product of (1.9) with $2k(u^n - u^{n-1})$ in H, we obtain

(2.26)
$$2|u^{n} - u^{n-1}|^{2} + \nu k||u^{n}||^{2} - \nu k||u^{n-1}||^{2} + \nu k||u^{n} - u^{n-1}||^{2} + 2k b(u^{n}, u^{n}, u^{n} - u^{n-1}) = 2k (f^{n}, u^{n} - u^{n-1}).$$

Using properties (1.23), (1.24), and (1.25) of the trilinear form b and recalling (2.2), we bound the nonlinear term as

(2.27)
$$2kb(u^{n}, u^{n}, u^{n} - u^{n-1}) = -2kb(u^{n}, u^{n}, u^{n-1})$$

$$\leq 2c_{b}k|u^{n}||u^{n}|||u^{n-1}||$$

$$\leq \frac{\nu}{2}k||u^{n}||^{2} + \frac{2c_{b}^{2}}{\nu}K_{1}k||u^{n-1}||^{2}.$$

We bound the right-hand side of (2.26) using Cauchy-Schwarz,

Relations (2.26)–(2.28) imply

(2.29)
$$2|u^{n} - u^{n-1}|^{2} + \frac{\nu}{2}k\|u^{n}\|^{2} - \left(\nu + \frac{2c_{b}^{2}}{\nu}K_{1}\right)k\|u^{n-1}\|^{2} + \frac{\nu}{2}k\|u^{n} - u^{n-1}\|^{2} \le \frac{2}{\nu\lambda_{1}}k|f^{n}|^{2},$$

from which we obtain

$$||u^n||^2 \le K_2 ||u^{n-1}||^2 + c_4 |f|_{\infty}^2,$$

and using hypothesis (2.21) we find

$$(2.31) 2c_2K_1k||u^n||^2 \le 2c_2K_1k\left(K_2||u^{n-1}||^2 + c_4|f|_{\infty}^2\right) < 1,$$

which contradicts (2.24). Therefore, (2.19) implies (2.23) and hence

(2.32)
$$||u^{n}||^{2} \leq \frac{1 - \left[1 - 4c_{2}K_{1}k\left(||u^{n-1}||^{2} + c_{3}k|f|_{\infty}^{2}\right)\right]^{1/2}}{2c_{2}K_{1}k}$$

$$= 2\frac{||u^{n-1}||^{2} + c_{3}k|f|_{\infty}^{2}}{1 + \sqrt{1 - x}},$$

where

$$x = 4c_2 K_1 k(\|u^{n-1}\|^2 + c_3 k|f|_{\infty}^2).$$

Since $x \leq 4/5$ by (2.21) and

$$\frac{2}{1+\sqrt{1-x}} \le 1 + \frac{x}{2}$$
 if $0 \le x \le \frac{4}{5}$,

relation (2.32) implies, under assumption (2.21), that

$$(2.33) ||u^n||^2 \le \left(||u^{n-1}||^2 + c_3k|f|_{\infty}^2\right) \left[1 + 2c_2K_1k\left(||u^{n-1}||^2 + c_3k|f|_{\infty}^2\right)\right].$$

Using (2.21) once again, (2.33) immediately implies (2.22).

In order to obtain estimates on a finite interval of time, we will inductively use Lemma 2.3, together with the following result, which was proved in [10] and which we repeat here for convenience.

LEMMA 2.4. Given k > 0, an integer $n_* > 0$, and positive sequences ξ_n , η_n , and ζ_n such that

(2.34)
$$\xi_n \leq \xi_{n-1}(1+k\eta_{n-1}) + k\zeta_n \quad \text{for } n = 1, \dots, n_*,$$

we have, for any $n \in \{2, \ldots, n_*\}$,

(2.35)
$$\xi_n \le \xi_0 \exp\left(\sum_{i=0}^{n-1} k \eta_i\right) + \sum_{i=1}^{n-1} k \zeta_i \exp\left(\sum_{j=i}^{n-1} k \eta_j\right) + k \zeta_n.$$

Proof. Using (2.34) recursively, we derive

$$\xi_n \le \xi_0 \prod_{i=0}^{n-1} (1 + k\eta_i) + \sum_{i=1}^n k\zeta_i \prod_{j=i}^{n-1} (1 + k\eta_j)$$

with the convention that $\prod_{j=\alpha}^{\beta} r_j = 1$ for $\beta < \alpha$. Using the fact that $1 + x \leq e^x$ for all $x \in \mathbb{R}$, the conclusion of the lemma follows.

PROPOSITION 2.5 (estimates on a finite interval). Let T > 0 and let $K_3(\cdot, \cdot, \cdot)$ be the function, monotonically increasing in all its arguments, given in (2.47). Suppose the timestep k is such that

$$(2.36) k \le \min\{\kappa_1, \kappa_2(|u_0|, |f|_{\infty}), \kappa_3(||u^0||, |f|_{\infty}, T)\},\$$

where κ_1 is given by (2.13), and

(2.37)
$$\kappa_2(|u_0|,|f|_{\infty}) = \frac{1}{10c_2c_4K_1|f|_{\infty}^2},$$

(2.38)
$$\kappa_3(\|u^0\|, |f|_{\infty}, T) = \frac{1}{10c_2K_1K_2K_3(\|u^0\|, |f|_{\infty}, T)}.$$

Then (i) relation (2.22) holds for all n = 1, ..., N := |T/k|, and (ii)

$$(2.39) ||u^n||^2 \le K_3(||u^0||, |f|_{\infty}, nk) \forall n = 1, \dots, N := |T/k|.$$

Proof. Let T > 0 and k be such that hypothesis (2.36) is satisfied. We will use induction on n.

Since $||u^0||^2 \le K_3(||u^0||, |f|_{\infty}, 0)$, (2.37) and (2.38) imply that condition (2.21) of Lemma 2.3 is satisfied for n = 1,

$$(2.40) c_2 K_1 k(K_2 ||u^0||^2 + c_4 |f|_{\infty}^2) \le \frac{1}{10} + \frac{1}{10} \le \frac{1}{5}.$$

By the same lemma, we have

$$(2.41) ||u^1||^2 \le ||u^0||^2 \left[1 + c_5 K_1 k \left(||u^0||^2 + k|f|_{\infty}^2\right)\right] + c_6 k|f|_{\infty}^2.$$

Now assume that (2.21) holds for n = 1, ..., m for some $m \le N$. Then by Lemma 2.3, (2.22) holds for n = 1, ..., m; furthermore, we can bound $||u^m||$ as follows. We write the stepwise bound (2.22) in Lemma 2.3 in the form

$$(2.42) \xi_n \le \xi_{n-1} (1 + k\eta_{n-1}) + k\zeta,$$

where

(2.43)
$$\xi_n = ||u^n||^2, \quad \eta_n = c_5 K_1(||u^n||^2 + k|f|_{\infty}^2), \quad \text{and} \quad \zeta = c_6 |f|_{\infty}^2.$$

Our intention is to apply (the discrete Gronwall) Lemma 2.4. So we compute for i > 0, using (2.3),

(2.44)
$$\sum_{j=i}^{m-1} k \eta_j = c_5 K_1 \sum_{j=i}^{m-1} k \left(\|u^j\|^2 + k|f|_{\infty}^2 \right) \\ \leq c_7 K_1 \left[K_1 + (m-i)k|f|_{\infty}^2 \right];$$

similarly, for i = 0,

(2.45)
$$\sum_{j=0}^{m-1} k \eta_j = c_5 K_1 \sum_{j=0}^{m-1} k \left(\|u^j\|^2 + k|f|_{\infty}^2 \right)$$

$$\leq c_7 K_1 \left(K_1 + mk|f|_{\infty}^2 \right) + c_5 K_1 k \|u^0\|^2.$$

We note that, using (2.38) and recalling that $K_2 \geq 2$, the last term can be bounded as

$$(2.46) c_5 K_1 k \|u^0\|^2 \le \frac{c_5 \|u^0\|^2}{10c_2 K_2 K_3 (\|u^0\|, |f|_{\infty}, T)}$$

$$\le \frac{c_5}{10c_2 K_2} \frac{\|u^0\|^2}{K_3 (\|u^0\|, |f|_{\infty}, 0)} \le \frac{c_5}{20c_2}.$$

The middle term in (2.35) here is

$$\sum_{i=1}^{m-1} k\zeta \exp\left(\sum_{j=i}^{m-1} k\eta_j\right) \le c_6 |f|_{\infty}^2 \sum_{i=1}^{m-1} k \exp(c_7 K_1^2 + c_7 K_1 (m-i)k |f|_{\infty}^2)$$

$$\le c_6 |f|_{\infty}^2 \exp(c_7 K_1^2) mk \exp(c_7 K_1 mk |f|_{\infty}^2).$$

The following bound on $||u^m||^2$ then follows from (2.35):

$$||u^{m}||^{2} \leq ||u^{0}||^{2} \exp(c_{7}K_{1}|f|_{\infty}^{2}mk) \exp(c_{7}K_{1}^{2} + c_{5}/(20c_{2}))$$

$$+ 2c_{6}|f|_{\infty}^{2} \exp(c_{7}K_{1}^{2})mk \exp(c_{7}K_{1}|f|_{\infty}^{2}mk)$$

$$=: K_{3}(||u^{0}||, |f|_{\infty}, mk).$$

We note that the bound K_3 depends on the initial discrete value through its norm $||u^0||$ and also on m, but this latter dependence is only through the time mk. We also note the dependence of K_3 on $|u_0|$ through K_1 , but K_1 bounds all $|u^n|^2$.

It is now clear that, given the hypothesis of the proposition, the timestep k satisfies condition (2.21) as long as $m \leq |T/k|$, completing the proof.

Now, since Proposition 2.5 gives a bound on $||u^n||^2$ that is valid on a finite time interval only, we are going to extend the result to infinite time by repeatedly applying it and the following (discrete uniform Gronwall) lemma, which is a slightly more general version of the discrete uniform Gronwall lemma of Shen [10].

LEMMA 2.6. Given k > 0, positive integers n_1, n_2, n_* such that $n_1 < n_*, n_1 + n_2 + 1 \le n_*$, positive sequences ξ_n, η_n , and ζ_n such that

(2.48)
$$\xi_n \le \xi_{n-1}(1 + k\eta_{n-1}) + k\zeta_n \quad \text{for } n = n_1, \dots, n_*,$$

and given the bounds

$$(2.49) \qquad \sum_{n=n'}^{n'+n_2} k \eta_n \le a_1(n_1, n_*), \quad \sum_{n=n'}^{n'+n_2} k \zeta_n \le a_2(n_1, n_*), \quad \sum_{n=n'}^{n'+n_2} k \xi_n \le a_3(n_1, n_*)$$

for any n' satisfying $n_1 \le n' \le n_* - n_2$, we have

(2.50)
$$\xi_n \le \left(\frac{a_3(n_1, n_*)}{kn_2} + a_2(n_1, n_*)\right) e^{a_1(n_1, n_*)}$$

for any n such that $n_1 + n_2 + 1 \le n \le n_*$.

Proof. Let n_3 and n_4 be such that $n_1 \leq n_3 - 1 \leq n_4 \leq n_2 + n_3 - 1 \leq n_* - 1$. Using (2.48) recursively, we derive

$$(2.51) \xi_{n_2+n_3} \le \xi_{n_4} \prod_{i=n_4}^{n_3+n_2-1} (1+k\eta_i) + \sum_{i=n_4+1}^{n_3+n_2} k\zeta_i \prod_{j=i}^{n_2+n_3-1} (1+k\eta_j)$$

with the convention that $\prod_{j=\alpha}^{\beta} r_j = 1$ for $\beta < \alpha$. Using the fact that $1 + x \leq e^x$ for all $x \in \mathbb{R}$, and recalling the first two assumptions in (2.49), we obtain

$$\xi_{n_2+n_3} \leq (\xi_{n_4} + a_2)e^{a_1}$$

Multiplying this inequality by k, summing n_4 from $n_3 - 1$ to $n_2 + n_3 - 2$, and using the third assumption in (2.49) gives the conclusion (2.50) of the lemma.

We are now in a position to give the main result, that is, to derive a uniform bound for $||u^n||$ for all $n \ge 1$.

THEOREM 2.7. Let $u_0 \in V$, $f \in L^{\infty}(\mathbb{R}_+; H)$, and u^n be the solution of the numerical scheme (1.9). Also, let $r \geq 4\kappa_1$ be arbitrarily fixed and let k be such that

$$(2.52) k \le \min\{\kappa_1, \kappa_2(|u_0|, |f|_{\infty}), \kappa_3(||u_0||, |f|_{\infty}, T_0 + r), \kappa_3(\rho_1, |f|_{\infty}, r)\},\$$

where $\kappa_1 = 1/(\nu\lambda_1)$ was defined in (2.13), $\kappa_2(\cdot, \cdot)$ and $\kappa_3(\cdot, \cdot, \cdot)$ are given in Proposition 2.5, T_0 , the time of entering an absorbing ball for $|u^n|$, is given by (2.14), and $\rho_1(|f|_{\infty}, r)$ is given in (2.57).

Then we have

$$||u^n||^2 \le K_5(||u_0||, |f|_{\infty}) \quad \forall n \ge 1,$$

where $K_5(\cdot,\cdot)$ is a continuous function defined on \mathbb{R}^2_+ , increasing in both arguments. Moreover,

$$(2.54) ||u^n||^2 \le K_4(|f|_{\infty}) \forall n \ge N_0 + N_r := |T_0/k| + |r/k|,$$

i.e., $||u^n||$ is bounded independently of u_0 beyond $N_0 + N_r$.

Proof. Let $r \geq 4\kappa_1$ be arbitrarily fixed and let k be such that (2.52) holds.

The idea for deriving a uniform bound for $||u^n||^2$ for all $n \ge 1$ is as follows:

- (i) Applying first Proposition 2.5 on $(0, T_0 + r)$ (that is, for $n = 1, ..., N_0 + N_r$), we get an upper bound for $||u^n||$ for $n = 1, ..., N_0 + N_r$; applying Lemma 2.6, we show that $||u^{N_0+N_r}||^2 \le \rho_1^2$, where $\rho_1(|f|_{\infty}, r)$ is defined in (2.57).
- (ii) Iterating Proposition 2.5 and Lemma 2.6, at each step $i \geq 2$, we show that for all $n = N_0 + (i-1)N_r + 1, \ldots, N_0 + iN_r$, $||u^n||^2$ is bounded by $K_3(||u^{N_0 + (i-1)N_r}||$,

 $|f|_{\infty}, r$); using the estimate on $||u^{N_0+(i-1)N_r}||$ from the previous step, we obtain that $||u^n||^2$ is bounded independently of the initial value for all $n = N_0 + (i-1)N_r + 1, \ldots, N_0 + iN_r$ for every $i \geq 2$ (and thus for all $n \geq N_0 + N_r$).

We now proceed to give a rigorous proof of the theorem.

Noting that, by hypothesis, k satisfies condition (2.36) of Proposition 2.5 with $T = T_0 + r$, we first apply Proposition 2.5 and obtain that (2.22) holds for all $n = 1, \ldots, N_0 + N_r$, and

$$(2.55) ||u^n||^2 \le K_3(||u^0||, |f|_{\infty}, nk) \forall n = 1, \dots, N_0 + N_r.$$

At this point we know that for k satisfying hypothesis (2.52),

(2.56)

$$||u^n||^2 \le ||u^{n-1}||^2 [1 + c_5 K_1 k(||u^{n-1}||^2 + k|f|_{\infty}^2)] + c_6 k|f|_{\infty}^2 \quad \forall n = 1, \dots, N_0 + N_r,$$

and we apply (the discrete uniform Gronwall) Lemma 2.6 with $\xi_n = ||u^n||^2$, $\eta_n = c_5 K_1(||u^n||^2 + k|f|_\infty^2)$, $\zeta_n = c_6|f|_\infty^2$, $n_1 = N_0 + 1$, $n_2 = N_r - 2$, and $n_* = N_0 + N_r$ to obtain a bound for $||u^{N_0 + N_r}||$. In computing the sums $a_1(n_1, n_*)$, $a_2(n_1, n_*)$, and $a_3(n_1, n_*)$ that appear there, we note that since all those sums are taken for $n \ge N_0$ and since, by hypothesis, k satisfies condition (2.13) of Corollary 2.2, we can replace K_1 , the bound on $|u^n|^2$, by $2\rho_0^2$, whenever the former appears. For every $n' = N_0 + 1$, $N_0 + 2$, we compute, using (2.3) and (2.14) for the first and last lines,

$$2c_5\rho_0^2 \sum_{n=n'}^{n'+n_2} \left(k\|u^n\|^2 + k^2|f|_{\infty}^2\right) \le c_8\rho_0^2(\rho_0^2 + r|f|_{\infty}^2),$$

$$c_6 \sum_{n=n'}^{n'+n_2} k|f|_{\infty}^2 \le c_6 r|f|_{\infty}^2,$$

$$\sum_{n=n'}^{n'+n_2} k\|u^n\|^2 \le c_9(\rho_0^2 + r|f|_{\infty}^2).$$

Using the conclusion (2.50) of Lemma 2.6 and the fact that $r \geq 4\kappa_1$, we obtain

(2.57)
$$||u^{N_0+N_r}||^2 \le \left[2c_9(\rho_0^2/r + |f|_\infty^2) + c_6 r|f|_\infty^2\right] \exp\left(c_8\rho_0^2(\rho_0^2 + r|f|_\infty^2)\right)$$
$$=: \rho_1(|f|_\infty; r)^2.$$

Now, since by hypothesis $k \leq \kappa_3(\rho_1, |f|_{\infty}, r)$ and since $\kappa_3(\cdot, \cdot, \cdot)$ is a decreasing function of its arguments, we can regard $u^{N_0+N_r}$ as our initial data and apply Proposition 2.5 with T=r. We obtain that relation (2.22) holds for all $n=N_0+N_r+1,\ldots,N_0+2N_r$, and

$$(2.58) ||u^n||^2 \le K_3(||u^{N_0+N_r}||, |f|_{\infty}, N_r k) \forall n = N_0 + N_r + 1, \dots, N_0 + 2N_r.$$

Thanks to (2.57) and to the fact that $K_3(\cdot,\cdot,\cdot)$ is an increasing function of all its arguments, we have

$$(2.59) ||u^n||^2 \le K_3(\rho_1, |f|_{\infty}, N_r k) \forall n = N_0 + N_r + 1, \dots, N_0 + 2N_r.$$

Applying again Lemma 2.6 with $n_1 = N_0 + N_r + 1$, $n_2 = N_r - 2$, and $n_* = N_0 + 2N_r$, we obtain

$$||u^{N_0+2N_r}||^2 \le \rho_1^2.$$

Iterating Proposition 2.5 and Lemma 2.6 and reasoning as above, we arrive at

$$(2.61) ||u^n||^2 \le K_3(\rho_1, |f|_{\infty}, r) =: K_4(|f|_{\infty}) \quad \forall n \ge N_0 + N_r,$$

and recalling (2.55), we conclude

(2.62)
$$||u^n||^2 \le \max\{K_3(||u_0||, |f|_{\infty}, T_0 + r), K_4(|f|_{\infty})\}$$
$$=: K_5(||u_0||, |f|_{\infty}) \quad \forall n \ge 1,$$

thus proving the theorem.

Acknowledgments. The authors thank Prof. R. Temam for suggesting the problem and for his help in the course of this work; they also thank the referees for helpful comments.

REFERENCES

- C. Bernardi and Y. Maday, Approximations Spectrales de Problèmes aux Limites Elliptiques, Math. Appl. (Berlin) 10, Springer-Verlag, Paris, 1992.
- [2] C. Foias, O. Manley, R. Rosa, and R. Temam, Navier-Stokes Equations and Turbulence, Encyclopedia Math. Appl. 83, Cambridge University Press, Cambridge, UK, 2001.
- [3] T. GEVECI, On the convergence of a time discretization scheme for the Navier-Stokes equations, Math. Comp., 53 (1989), pp. 43-53.
- [4] V. GIRAULT AND P.-A. RAVIART, Finite Element Approximation of the Navier-Stokes Equations, Lecture Notes in Math. 749, Springer-Verlag, New York, 1979.
- [5] D. GOTTLIEB AND S. ORSZAG, Numerical Analysis of Spectral Methods, Theory and Applications, CBMS-NSF Regional Conf. Ser. in Appl. Math. 26, SIAM, Philadelphia, 1977.
- [6] J. G. HEYWOOD AND R. RANNACHER, Finite element approximation of the nonstationary Navier-Stokes problem. I. Regularity of solutions and second-order error estimates for spatial discretization, SIAM J. Numer. Anal., 19 (1982), pp. 275-311.
- [7] N. Ju, On the global stability of a temporal discretization scheme for the Navier-Stokes equations, IMA J. Numer. Anal., 22 (2002), pp. 577-597.
- [8] J. LERAY, Etude de diverses équations intégrales non linéaires et de quelques problèmes que pose l'hydrodynamique, J. Math. Pures Appl., 12 (1933), pp. 1–82.
- [9] M. MARION AND R. TEMAM, Navier-Stokes equations. Theory and approximation, in Handbook of Numerical Analysis, Vol. VI, North-Holland, Amsterdam, 1998, pp. 503-688.
- [10] J. Shen, Long time stabilities and convergences for the fully discrete nonlinear Galerkin methods, Appl. Anal., 38 (1990), pp. 201–229.
- [11] R. Temam, Navier-Stokes Equations and Nonlinear Functional Analysis, 2nd ed., CBMS-NSF Regional Conf. Ser. in Appl. Math. 66, SIAM, Philadelphia, 1995.
- [12] R. Temam, Infinite-Dimensional Dynamical Systems in Mechanics and Physics, 2nd ed., Appl. Math. Sci. 68, Springer-Verlag, New York, 1997.
- [13] R. Temam, Navier-Stokes Equations: Theory and Numerical Analysis, AMS Chelsea, Providence, RI, 2001 (reprint of the 1984 edition).