

Lecture 2: Numerical Methods for Hopf bifurcations and periodic orbits in large systems

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- 1 Introduction
- 2 Calculation of Hopf points
- 3 Hopf detection using bifurcation theory
- 4 Hopf detection using Complex Analysis
- 5 Hopf detection using the Cayley Transform
- 6 Stable and unstable periodic orbits

Outline

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Recap and plan for today

- **Lecture 1:**

- 1 Compute paths of $F(x, \lambda) = 0$ using pseudo-arclength
- 2 Detect singular points $\text{Det}(F_x(x, \lambda)) = 0$
- 3 Compute paths of singular points in two-parameter problems
- 4 bordered systems
- 5 4-6 cell interchange in the Taylor problem

- **Lecture 2:**

- Accurate calculation of Hopf points
- Detection of Hopf bifurcations (find **pure imaginary** eigenvalues in a **large sparse** parameter-dependent matrix)
 - 1 Bifurcation theory
 - 2 Complex analysis
 - 3 Cayley transform
- Stable and unstable periodic orbits

Lecture 1: Compute singular points

- Seek (x, λ) such that $F_x(x, \lambda)$ is singular
- Consider

$$\begin{bmatrix} F_x(x, \lambda) & F_\lambda(x, \lambda) \\ c^T & d \end{bmatrix} \begin{bmatrix} * \\ g \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

- $\text{Det}(F_x) = 0 \iff g = 0$.
- Accurate calculation: Consider the pair

$$F(x, \lambda) = 0, \quad g(x, \lambda) = 0$$

- Newton's Method:

$$\begin{bmatrix} F_x(x, \lambda) & F_\lambda(x, \lambda) \\ g_x(x, \lambda)^T & g_\lambda(x, \lambda) \end{bmatrix} \begin{bmatrix} \Delta x \\ \Delta \lambda \end{bmatrix} = - \begin{bmatrix} F \\ g \end{bmatrix}$$

- System nonsingular if $\frac{d}{dt}\mu \neq 0$ at singular point

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Accurate calculation of Hopf points

- Assume $A(\lambda) = F_x(x, \lambda)$ is real and nonsingular
- At Hopf point: $A(\lambda)$ has eigenvalues $\pm i\omega$
- $\text{Rank}(A(\lambda)^2 + \omega^2 I) = n - 2$

Accurate calculation of Hopf points

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- At Hopf point: $A(\lambda)$ has eigenvalues $\pm i\omega$
- $\text{Rank}(A(\lambda)^2 + \omega^2 I) = n - 2$
- Calculate Hopf point using 2-bordered matrix: set up

$$F(x, \lambda) = 0, \quad g(x, \lambda, \omega) = 0, \quad h(x, \lambda, \omega) = 0$$

where

$$\begin{bmatrix} A(\lambda)^2 + \omega^2 I & B \\ C^T & D \end{bmatrix} \begin{bmatrix} * \\ g \\ h \end{bmatrix} = \begin{bmatrix} 0 \\ r_1 \\ r_2 \end{bmatrix}$$

- Newton system, $(n + 2) \times (n + 2)$, needs $g_x, g_\lambda, g_\omega, h_x, \dots$
- Block version of (D)+iterative refinement on (C)
- 2-bordered matrix is nonsingular if complex pair cross imaginary axis “smoothly”



Hopf continued

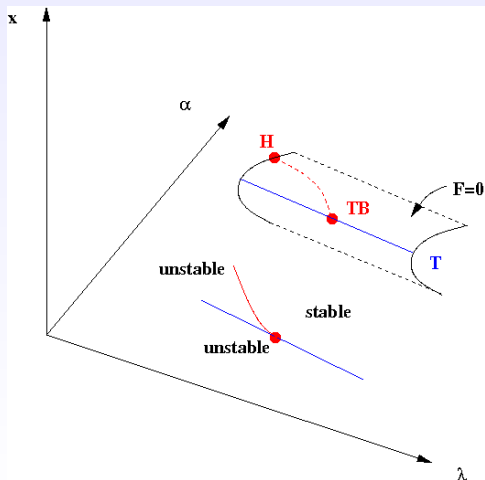
- $A(\lambda) = F_x(x, \lambda)$
- If you don't want to form $A(\lambda)^2$: split complex eigenvector/eigenvalue into Real and Imaginary parts and work with $(2n + 2) \times (2n + 2)$ matrices involving $A(\lambda)$
- Extensions for N-S: $A(\lambda)\phi = \mu B\phi$
- **BUT**: Whatever system is used, **accurate estimates** for λ and ω are needed
- Compute **paths** of Hopf points in 2-parameter problems (**3-bordered matrices**)
- Summary of methods: Govaerts (2000)

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Bifurcation Theory: Takens-Bogdanov (TB) point

At a TB point, F_x has a 2-dim Jordan block, i.e. $\begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$. A typical picture is:



“Organising Centre” Algorithm

- Two parameter problem $F(x, \lambda, \alpha) = 0$
- Fix α . Compute a Turning point in (x, λ) (Easy!). Remember:

$$F_x \phi = 0, \quad (F_x)^T \psi = 0$$

- For the 2-parameter problem: Compute path of Turning points looking for $\psi^T \phi = 0$ (TB point) (Easy)
- Jump onto path of Hopf points (symmetry-breaking) (Easy)
- Compute path of Hopf points (pseudo-arclength) (Easy)
- In parameter space the paths of Hopf and Turning points are **tangential** at TB

5 cell anomalous flows in the Taylor Problem

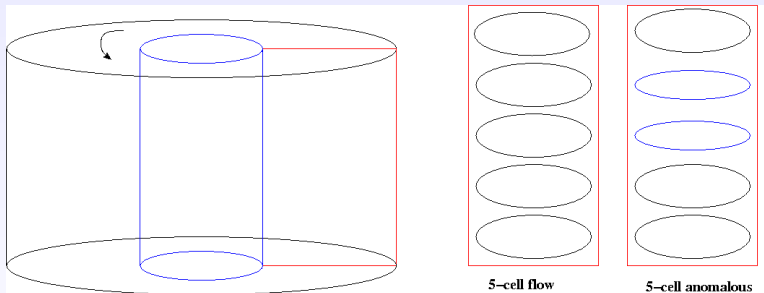


Figure: Two different 5-cell flows

5-cell flows experimental results

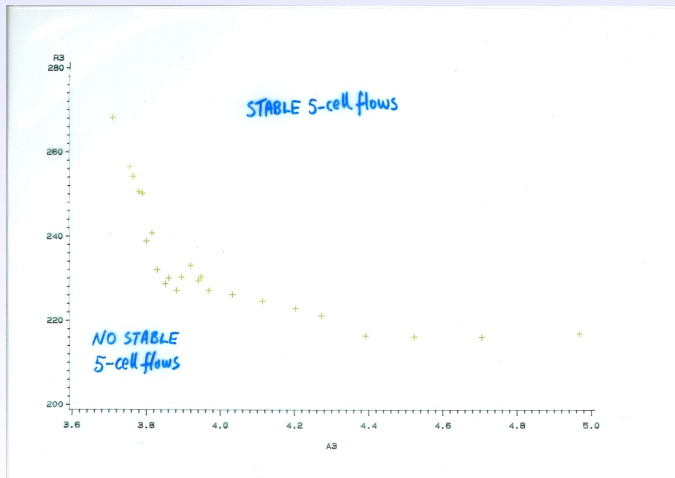


Figure: parameter space plots of 5-cell flows

5-cell flows numerical results (Anson)

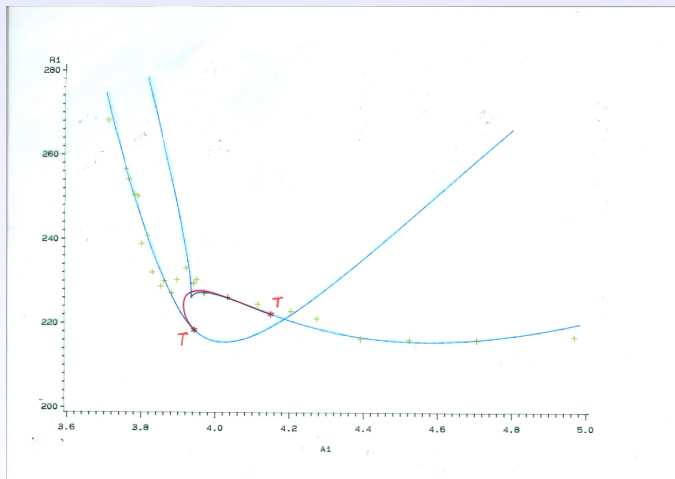


Figure: parameter space plots of 5-cell flows

“Organising Centre” approach

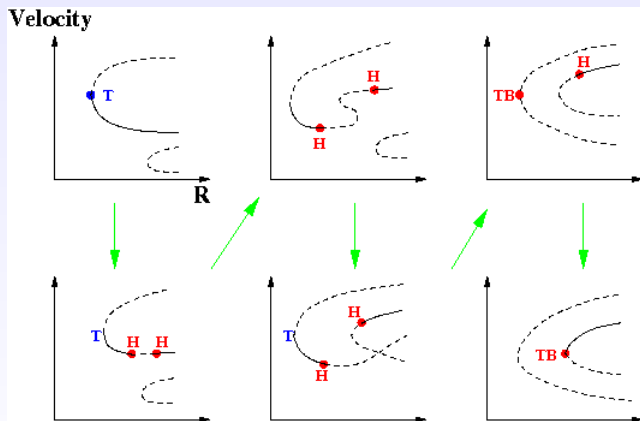


Figure: 5-cell flows: Sequence of Bifurcation diagrams as aspect ratio changes

This understanding wouldn't be possible without the numerical results

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The “idea”: Govaerts/Spence (1996)

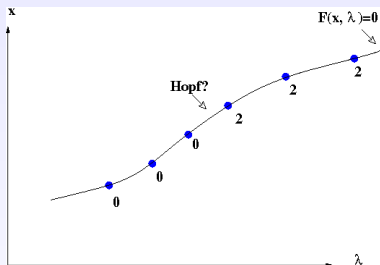


Figure: For each point on $F(x, \lambda) = 0$ can we calculate the **number** of eigenvalues in the unstable half plane?

Why Nice?

(a) Seek an integer, and (b) Estimate for $\text{Im}(\mu)$ not needed.

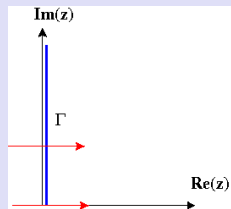
Complex Analysis

Winding number

If $g(z)$ is analytic in Γ

$$\begin{aligned} N - P &= \frac{1}{2\pi} [\arg g(z)]_{\Gamma} \\ &= \text{Winding Number} \\ &= W(g) \end{aligned}$$

Contour for real matrices



Algorithm

- “Counting Sectors”: Ying/Katz (1988) (based on Henrici (1974))

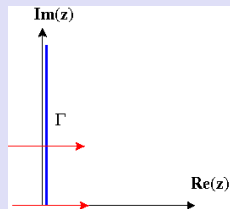
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Algorithm

- “Counting Sectors”: Ying/Katz (1988) (based on Henrici (1974))
- If g changes so that a **real pole** crosses Left to Right, $W(g)$ **decreases** by π . (**real zero** crosses L to R then $W(g)$ **increases**)
- If g changes so that a complex pole crosses Left to Right, $W(g)$ decreases by 2π

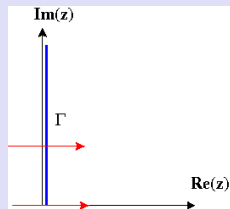
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- If g changes so that a complex pole crosses Left to Right, $W(g)$ decreases by 2π
- Need to evaluate $g(iy)$ on Γ

How to choose $g(z)$?

- Don't choose $g(z) = \text{Det}(A(\lambda) - zI)$
- $g(z) = c^T(A(\lambda) - zI)^{-1}b$
- Schur complement of $M = \begin{bmatrix} A(\lambda) - zI & b \\ c^T & 0 \end{bmatrix}$
- poles are eigenvalues of $A(\lambda)$; zeros depend on choices of b and c . Choose b and c so that the zeros “cancel” the poles to keep $W(g)$ “small”
- Need to evaluate

$$g(iy) = c^T(A(\lambda) - iyI)^{-1}b$$

as y moves up Imaginary axis (Ying/Katz algorithm chooses y 's)

The Tubular Reactor problem (Govaerts/Spence, 1996)

- Coupled pair of nonlinear parabolic PDEs for Temperature and Concentration
- Scaling: for a complex pole crossing Imag axis $W(g)$ reduces by 4

The Tubular Reactor problem (Govaerts/Spence, 1996)

- Coupled pair of nonlinear parabolic PDEs for Temperature and Concentration
- Scaling: for a complex pole crossing Imag axis $W(g)$ reduces by 4
- Winding numbers for 3 choices of g

point on path	$W(g_1)$	$W(g_2)$	$W(g_3)$
1	3	5	1
2	3	5	1
3	3	5	3*
4	3	5	3
5	-1^\dagger	1^\dagger	-1^\dagger
6	-1	3^\ddagger	1^\ddagger

- ❶ * = zero of g_3
- ❷ † = Hopf!
- ❸ ‡ = zero of g_2 and g_3 .

Final comments on “Winding Number” algorithm

- Govaerts/Spence was “proof of concept”: tested on a “not too difficult” problem
- Work is to evaluate

$$g(iy) = c^T (A(\lambda) - iyI)^{-1} b$$

as y moves up Imaginary axis

- For PDE matrices - Krylov solvers/model reduction?
- Ideas from yesterday’s lectures by [Strakos](#) (scattering amplitude) and [Ernst](#) (frequency domain).
- Also: [Stoll, Golub, Wathen \(2007\)](#)
- Note: you choose b and c !

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The Cayley Transform

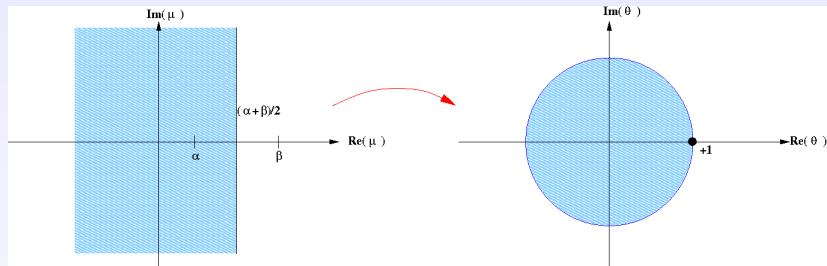


Figure: The mapping of μ to θ

- $A\phi = \mu B\phi$
- Choose α and β and form:

$$C = (A - \alpha B)^{-1}(A - \beta B) \quad \text{The Cayley transform}$$

- As λ varies, if μ crosses the line $\text{Re}(\alpha + \beta)/2$ then θ moves outside the unit ball



Hopf detection using the Cayley Transform

- Mapping

$$\theta = (\mu - \alpha)^{-1}(\mu - \beta)$$

- So $\beta = -\alpha$ maps left-half plane (“stable”) into unit circle
- Algorithm: At each point on $F(x, \lambda) = 0$:
 - 1 Choose α, β
 - 2 monitor dominant eigenvalues of $C = (A - \alpha B)^{-1}(A - \beta B)$
- Don’t need to know $\text{Im}(\mu)$
- Successfully computed Hopf bifurcations in [Taylor problem](#) and [Double-diffusive convection](#)
- BUT: “large” eigenvalues, μ , “cluster” at $\theta = 1$

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Periodic orbits

Theory

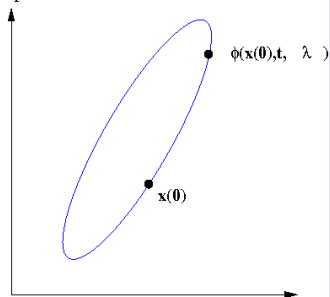
- $\dot{x} = F(x, \lambda), \quad x(t) \in \mathbb{R}^n$
- $x(0) = x(T), \quad T = \text{period}$
- Solution (“flow”): $\phi(x(0), t, \lambda)$
- Periodic: $\phi(x(0), T, \lambda) = x(0)$
- **Phase condition**: $s(x(0)) = 0$
- Stability: Monodromy matrix

$$\phi_x = \frac{\partial \phi}{\partial x(0)}(x(0), T, \lambda)$$

- $\mu_i \in \sigma(\phi_x)$: Floquet multipliers
- Stability: $|\mu_i| < 1, i = 2 \dots n$
($\mu_1 = 1$)
- **Monodromy matrix is FULL**

Phase plane

Phase plane



Stability of periodic orbits

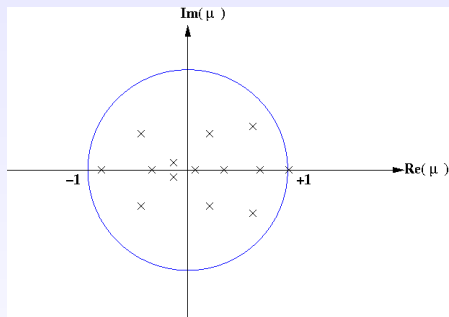


Figure: Plot of Floquet multipliers for a stable periodic orbit

- Loss of stability: multiplier crosses unit circle (e.g. real eigenvalue crosses through -1 then “period-doubling bifurcation”)
- If solution is stable just integrate in time: OK if μ_i not near unit circle
- “Integrate in time” is no good for unstable orbits

Newton-Picard Method for periodic orbits (Lust et. al.)

- Unknowns: **initial condition**, $x(0)$, and **period**, T , (drop λ)
- Fixed point problem + phase condition

$$\phi(x(0), T) = x(0), \quad s(x(0)) = 0$$

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- **Picard Iteration**: Guess $(x^{(0)}(0), T^{(0)})$ and compute $x^{(1)}(0)$

$$\phi(x^{(0)}(0), T^{(0)}) = x^{(1)}(0)$$

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$$\begin{bmatrix} \phi_x - I & \phi_T \\ s_x & 0 \end{bmatrix} \begin{bmatrix} \Delta x(0) \\ \Delta T \end{bmatrix} = - \begin{bmatrix} r_1 \\ r_2 \end{bmatrix}$$

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- Newton-Picard Method**: Split \mathbb{R}^n into “stable” and “unstable” subspaces. Convergence? - **Modified Newton**

- Picard** on “stable” subspace (large)
- Newton** on “unstable” subspace (small)
- Schroff&Keller**: “Recursive Projection Method” - computing stable and unstable steady states using initial value codes



Newton-Picard Method for periodic orbits

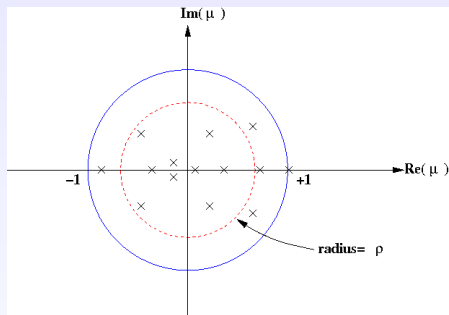
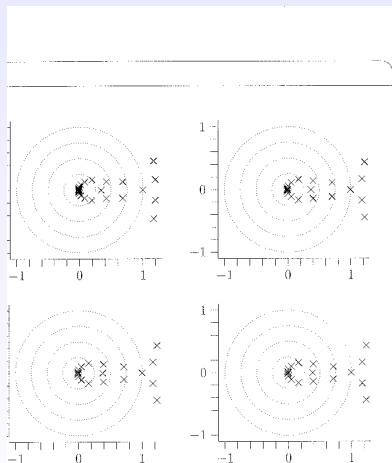


Figure: Splitting of Floquet multipliers into “stable” and “unstable” subsets

- Pick $\rho < 1$
- “Stable”: $|\mu_i| < \rho$ (hopefully dimension $\approx n$)
- “Unstable”: $|\mu_i| \geq \rho$ (hopefully dimension very small)

Floquet multipliers for the Brusselator

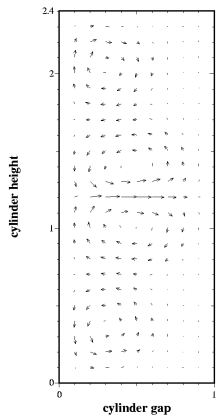
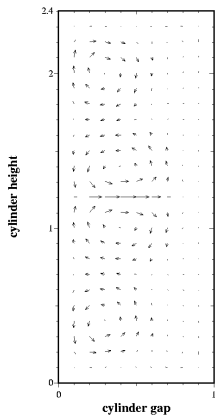


Four plots of the Floquet multipliers for the Brusselator model, solution on branch VI at $L = 1.9$, with 31 (upper left), 63 (lower left) and 127 (lower right) multipliers.

Lots of Numerical Linear Algebra!

- 1 Find (orthogonal) basis for “unstable” space, called V
- 2 Construct projectors onto “unstable” and “stable” spaces
- 3 need the **action** of ϕ_x on V (implemented by a small number of ODE solves)
- 4 need to increase /decrease dimension of V as Floquet multipliers enter or leave the “unstable” space
- 5 need to compute paths of periodic orbits: use pseudo-arclength (bordered matrices)

Taylor problem with counter-rotating cylinders:
Grande/Tavener/Thomas (2008)



Conclusions

- An efficient method to roughly “detect” a Hopf bifurcation in large systems is still an **open problem**
- Methods exist for accurate calculation once good starting values are known
- Look again at the winding number algorithm?
- Computation of stable and unstable periodic solutions for discretised PDEs (e.g. Navier-Stokes) is **wide open!**
- Software:
 - 1 LOCA “Library of Continuation Algorithms” Sandia (PDEs)
 - 2 MATCONT “Continuation software in Matlab”: W Govaerts
 - 3 AUTO