

Soler model: stability, bi-frequency solitons, and $SU(1,1)$

Andrew Comech, Texas A&M University and I.I.T.P. (Moscow)
Nabile Boussaïd, Université Bourgogne – Franche-Comté, Besançon, France

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$$\text{Einstein: } E^2 = p^2 + m^2$$

$$\text{Schrödinger: } (i\partial_t)^2\psi = (-i\nabla)^2\psi + m^2\psi$$

$$E = \sqrt{m^2 + p^2} \approx m + \frac{p^2}{2m}$$

$$[\text{Schrödinger}^{26}]: i\partial_t\psi = \frac{1}{2m}(-i\nabla)^2\psi$$

$$[\text{Dirac}^{28}]: E = \sqrt{p^2 + m^2} = \alpha \cdot p + \beta m,$$

$$i\partial_t\psi = \underbrace{(-i\alpha \cdot \nabla + \beta m)}_{D_m} \psi, \quad \psi(x, t) \in \mathbb{C}^4, \quad x \in \mathbb{R}^3$$

Dirac matrices α_j ($1 \leq j \leq 3$) and β : self-adjoint, $(-i\alpha \cdot \nabla + \beta m)^2 = -\Delta + m^2$

$$\text{Standard choice: } \alpha_j = \begin{bmatrix} 0 & \sigma_j \\ \sigma_j & 0 \end{bmatrix} \text{ (Pauli matrices), } \beta = \begin{bmatrix} I_2 & 0 \\ 0 & -I_2 \end{bmatrix}$$

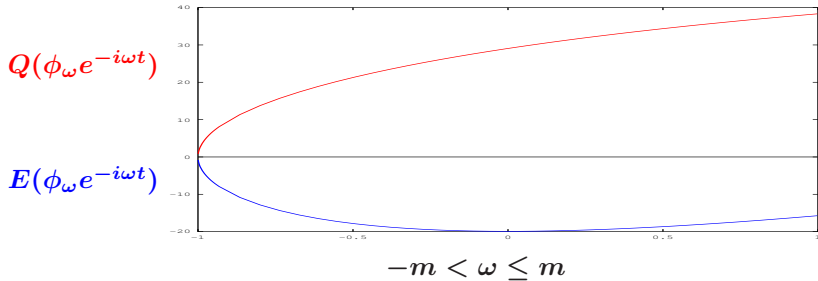
Dirac–Maxwell system

$$\bar{\psi} := \psi^* \beta$$

$$\gamma^\mu (i\partial_\mu - eA_\mu)\psi = m\psi, \quad \square A^\mu = e\bar{\psi}\gamma^\mu\psi, \quad \partial_\mu A^\mu = 0.$$

[*Esteban et al.*⁹⁶, *Abenda*⁹⁸]: solitary waves $\psi(x, t) = \phi(x)e^{-i\omega t}$, $\omega \in (-m, m)$

Dirac–Maxwell and Dirac–Coulomb systems



[*Comech*¹⁵]: solitary waves with $\omega = m$ for Dirac–Coulomb; $\phi(x) \sim e^{-\text{const}\sqrt{|x|}}$.

$$\gamma^0 = \beta, \quad \gamma^j = \beta\alpha^j, \quad 1 \leq j \leq 3$$

Self-interacting spinors

$$\bar{\psi} := \psi^* \beta$$

Models of self-interacting spinor field:

[*Ivanenko*³⁸, *Finkelstein et al.*⁵¹, *Finkelstein et al.*⁵⁶, *Heisenberg*⁵⁷] [...]

Soler model [*Ivanenko*³⁸, *Soler*⁷⁰], scalar self-interaction:

$$\mathcal{L}_{\text{Soler}} = \bar{\psi}(i\gamma^\mu \partial_\mu - m)\psi + |\bar{\psi}\psi|^{k+1} \quad k > 0$$

in (1+1)D: massive Gross-Neveu model [*Gross & Neveu*⁷⁴]

Massive Thirring model [*Thirring*⁵⁸], vector self-interaction:

$$\mathcal{L}_{\text{MTM}} = \bar{\psi}(i\gamma^\mu \partial_\mu - m)\psi + |\bar{\psi}\gamma_\mu\psi \bar{\psi}\gamma^\mu\psi|^{\frac{k+1}{2}} \quad k > 0$$

$\mathcal{J}^\mu(x) = \bar{\psi}(x)\gamma^\mu\psi(x)$: “charge-current density”

Soler model: NLD with scalar self-interaction

$$i\partial_t\psi = \underbrace{(-i\alpha \cdot \nabla + m\beta)}_{D_m} \psi - (\bar{\psi}\psi)^k \beta\psi, \quad \psi(x, t) \in \mathbb{C}^N, \quad x \in \mathbb{R}^n$$

- [Soler⁷⁰, Cazenave & Vázquez⁸⁶]: existence of solitary waves in \mathbb{R}^3 ,

$$\psi(x, t) = \phi_\omega(x)e^{-i\omega t}, \quad \omega \in (0, m), \quad \phi_\omega \in H^1(\mathbb{R}^3)$$

- Attempts at stability: [Bogolubsky⁷⁹, Alvarez & Soler⁸⁶, Strauss & Vázquez⁸⁶] ...
- Numerics suggest that (all?) solitary waves in 1D cubic Soler model are stable: [Alvarez & Carreras⁸¹, Alvarez & Soler⁸³, Cooper et al.¹⁰, Berkolaiko & Comech¹²], [Mertens et al.¹², Shao et al.¹⁴, Cuevas-Maraver et al.¹⁴]
- Assuming linear stability, one tries to prove asymptotic stability [Pelinovsky & Stefanov¹²] [Boussaid & Cuccagna¹²] [Comech, Phan, Stefanov¹⁴]

Nonrelativistic limit of NLD: $\omega \lesssim m$

[Ounaies⁰⁰, Guan⁰⁸]

Solitary wave: $\psi(x, t) = \begin{bmatrix} v(x) \\ u(x) \end{bmatrix} e^{-i\omega t}, \quad v, u \in \mathbb{C}^2 \quad x \in \mathbb{R}^3$

$$i\dot{\psi} = \left\{ -i \begin{bmatrix} 0 & \sigma \cdot \nabla \\ \sigma \cdot \nabla & 0 \end{bmatrix} + (m - (\bar{\psi}\psi)^k) \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \right\} \psi,$$

$$\omega \begin{bmatrix} v \\ u \end{bmatrix} \approx -i\sigma \cdot \nabla \begin{bmatrix} u \\ v \end{bmatrix} + (m - |v|^{2k}) \begin{bmatrix} v \\ -u \end{bmatrix}$$

If $\omega \lesssim m$, $|u| \ll |v| \ll 1$: $2mu \approx -i\sigma \cdot \nabla v$, v satisfies NLS:

$$-(m - \omega)v = -\frac{1}{2m}\Delta v - |v|^{2k}v$$

Scaling: $v(x) = \epsilon^{1/k}\Phi(\epsilon x)$, $\epsilon = \sqrt{m - \omega}$,

$$-\Phi = -\frac{1}{2m}\Delta\Phi - |\Phi|^{2k}\Phi, \quad \Phi(x) \in \mathbb{R}, \quad x \in \mathbb{R}^n.$$

NLD: linearization at a solitary wave

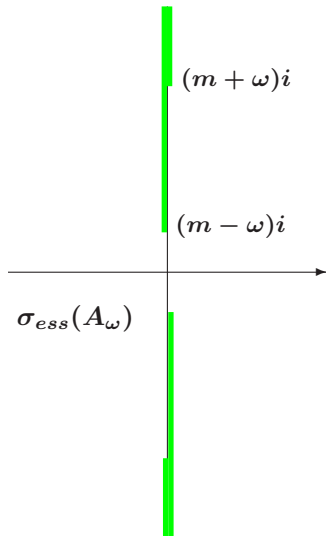
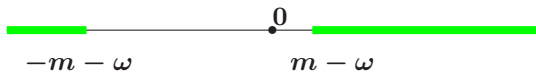
Given $\phi_\omega(x)e^{-i\omega t}$, $\omega \in (-m, m)$, consider $\psi(x, t) = (\phi_\omega(x) + r(x, t))e^{-i\omega t}$

Linearized eqn on $r(x, t)$, $i\partial_t r = D_m r - \omega r + \dots$

$$\partial_t \begin{bmatrix} \operatorname{Re} r \\ \operatorname{Im} r \end{bmatrix} = \underbrace{\begin{bmatrix} 0 & D_m - \omega + \dots \\ -D_m + \omega + \dots & 0 \end{bmatrix}}_{A_\omega} \begin{bmatrix} \operatorname{Re} r \\ \operatorname{Im} r \end{bmatrix}$$

$\partial_t R = A_\omega R$; $\sigma(A_\omega) \subset i\mathbb{R}$ (“spectral stability”) ???

$\sigma(D_m - \omega)$



Linear instability of NLD

$$i\partial_t \psi = D_m \psi - (\bar{\psi}\psi)^k \beta \psi, \quad x \in \mathbb{R}^n$$

$$\psi(x, t) = \phi_\omega(x) e^{-i\omega t}$$

Theorem 1 ([[Comech, Guan, Gustafson¹²](#)]).

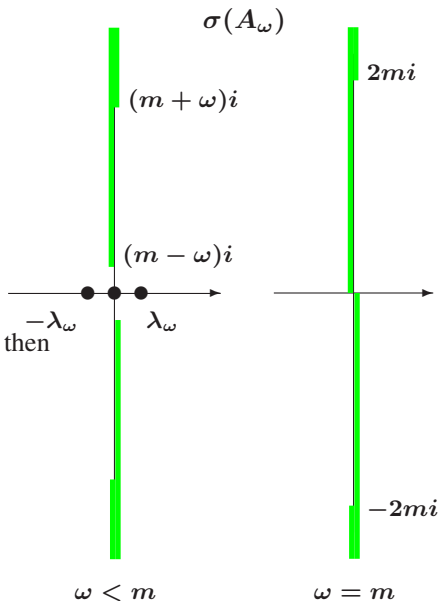
|| If $k > 2/n$, so that NLS_k is linearly unstable, then
 for $\omega \lesssim m$, $\exists \pm\lambda_\omega \in \sigma_d(A_\omega)$,
 $\text{Re } \lambda_\omega > 0$, $\lambda_\omega \xrightarrow{\omega \rightarrow m} 0$

Unstable for $n = 1, k > 2$ (above quintic)

Unstable for $n = 2, k > 1$ (above cubic)

Unstable for $n = 3, k > 2/3$

Proof: Rescale; use Schur reduction and Rayleigh-Schrödinger perturbation theory. \square



Bifurcations from σ_{ess}

Let $0 \leq \omega_0 \leq m$

Theorem 2 ([*Boussaid & Comech*¹⁶]).

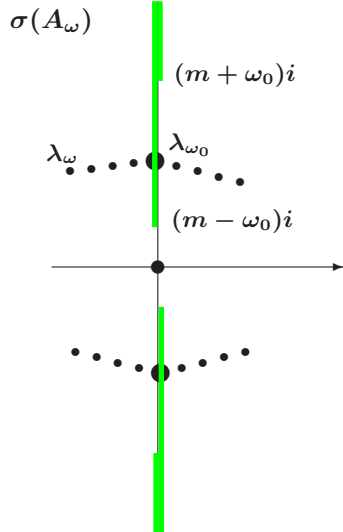
Assume: $\lambda_\omega \in \sigma_p(A_\omega)$, $\text{Re } \lambda_\omega \neq 0$,

$$\lambda_\omega \xrightarrow{\omega \rightarrow \omega_0} \lambda_{\omega_0} \in i\mathbb{R}$$

Then:

$$|\lambda_{\omega_0}| \leq m + \omega_0,$$

$$\lambda_{\omega_0} \in \sigma_p(A_{\omega_0}) \cup \{0\} \cup \{i(m + \omega_0)\}$$



LAP methods of [*Agmon*⁷⁵], [*Berthier & Georgescu*⁸⁷]:

$$\|u\|_{H_{-s}^1} \leq c_{z,s,\delta} \|(D_m - z)u\|_{L_s^2}, \quad z \in \mathbb{C} \setminus [-m, m], \quad s > 1/2, \quad |z \pm m| > \delta$$

$$\|u\|_{H_s^1} := \|(1 + |x|)^s u\|_{H^1}$$

Linear stability of NLD

$$i\partial_t\psi = D_m\psi - (\bar{\psi}\psi)^k\beta\psi, \quad x \in \mathbb{R}^n$$

$$\psi(x, t) = \phi_\omega(x)e^{-i\omega t}$$

Theorem 3 ([Comech¹¹]).

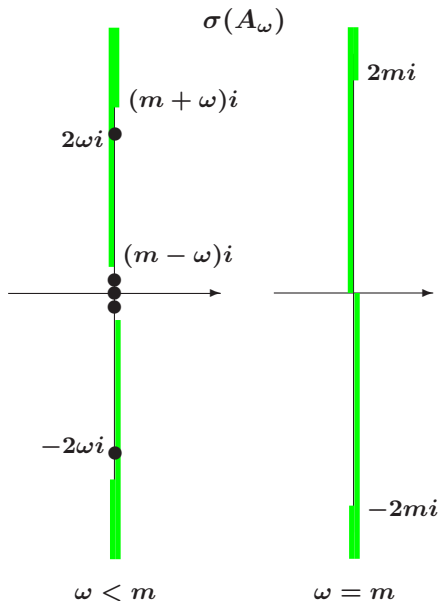
|| If $n \leq 3$, then $\lambda = \pm 2\omega i \in \sigma_p(A_\omega)$.

Theorem 4 ([Boussaid & Comech¹⁶]).

|| $n \leq 3, \quad k \lesssim 2/n:$

|| $\phi_\omega e^{-i\omega t}$ are linearly stable for $\omega \lesssim m$

|| Including the “charge-critical” case $k = 2/n$



Eigenvalues $\lambda = \pm 2\omega i$ in the Soler model

$\lambda = \pm 2\omega i$ is bad for proving asymptotic stability!

[*Boussaid & Cuccagna*¹²], [*Comech, Phan, Stefanov*¹⁴]

Asymptotic stability of solitons in Soler model for “radially symmetric” case:

1. perturbations orthogonal to translations;
2. perturbations orthogonal to $\pm 2\omega i$ eigenvectors

SU(1,1) invariance of the Soler model and the Soler charge

Theorem 5 ([*Boussaid & Comech*¹⁶]).

1. Soler model hamiltonian is invariant under transformations

$$\psi(x, t) \longrightarrow (a - ib\gamma^2\mathbb{C})\psi(x, t), \quad a, b \in \mathbb{C}, \quad |a|^2 - |b|^2 = 1$$

2. These transformations form the group $\text{SU}(1, 1) \supset \text{U}(1)$

$$\begin{bmatrix} a & b \\ \bar{b} & \bar{a} \end{bmatrix}$$

3. Conserved charges: $Q = \int \psi^* \psi dx > 0$

$$\Lambda = \int \psi^t (-i\gamma^2) \psi dx \in \mathbb{C}$$

4. Bi-frequency solutions: if $\phi(x)e^{-i\omega t}$ is a solitary wave, so is

$$\psi(x, t) = ae^{-i\omega t}\phi(x) - be^{i\omega t}i\gamma^2\mathbb{C}(\phi(x))$$

Note: if $|b| \ll 1$, then $\psi(x, t) \approx e^{-i\omega t}(\phi(x) - be^{2i\omega t}i\gamma^2\mathbb{C}(\phi(x)))$

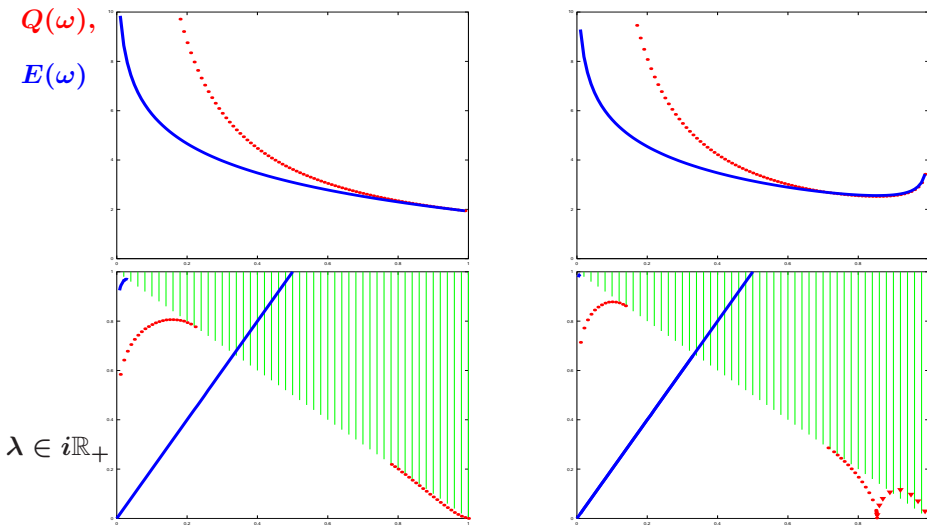


Figure 1: (1+1)D Soler model. LEFT: $k = 2$ (“charge-critical”); RIGHT: $k = 3$.
 TOP ROW: charge and energy of the solitary waves as functions of $\omega \in (0, 1)$.
 BOTTOM ROW: Spectrum on the upper half of the imaginary axis. Note the exact eigenvalue $\lambda = 2\omega i$.

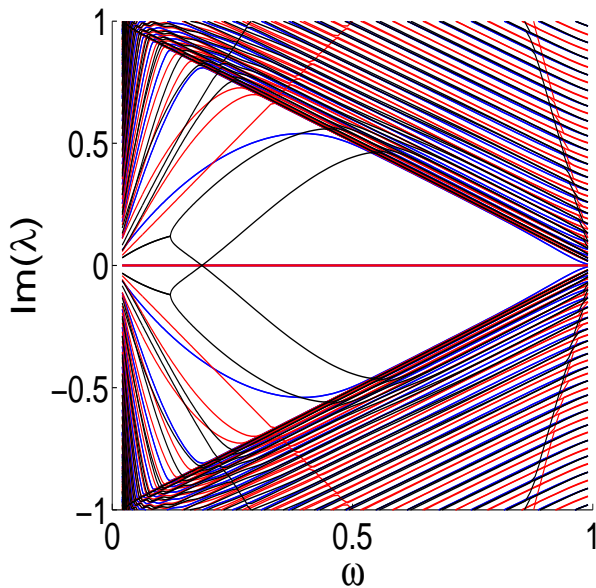


Figure 2: (2+1)D cubic “charge-critical” Soler model [*Cuevas-Maraver et al.*¹⁶]

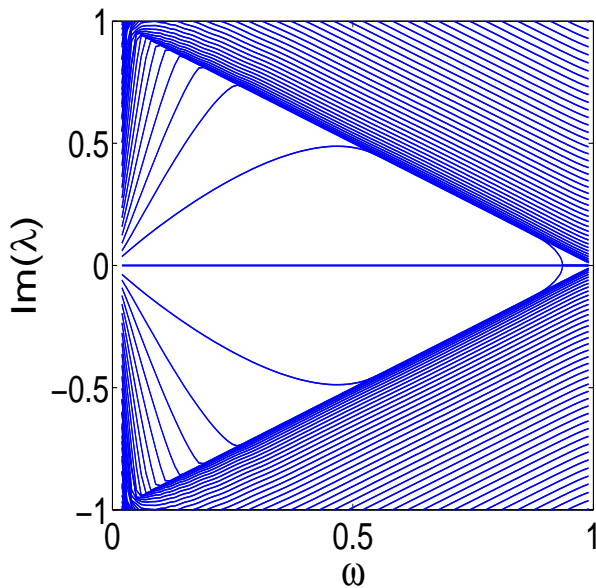


Figure 3: (3+1)D cubic “charge-supercritical” Soler model [*Cuevas-Maraver et al.*¹⁶]

Theorem 6 ([[Berkolaiko, Comech, Sukhtayev¹⁵](#)]).

|| Both $Q'(\omega) = 0$ and $E(\omega) = 0$ indicate collision of eigenvalues at $\lambda = 0$

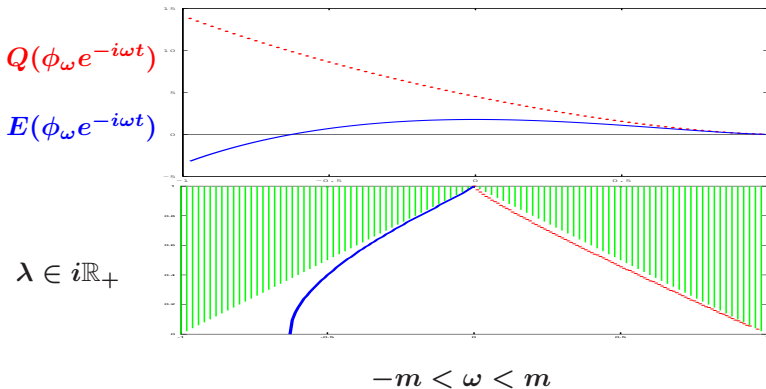


Figure 4: “Quadratic” massive Thirring model, $k = 1/2$.

TOP: energy and charge as functions of $\omega \in (-m, m)$.

BOTTOM: The spectrum of the linearization at a solitary wave.

Dotted eigenvalue collides with its opposite at the origin when $\omega_\star \approx -0.6276m$, where $E(\omega_\star) = 0$

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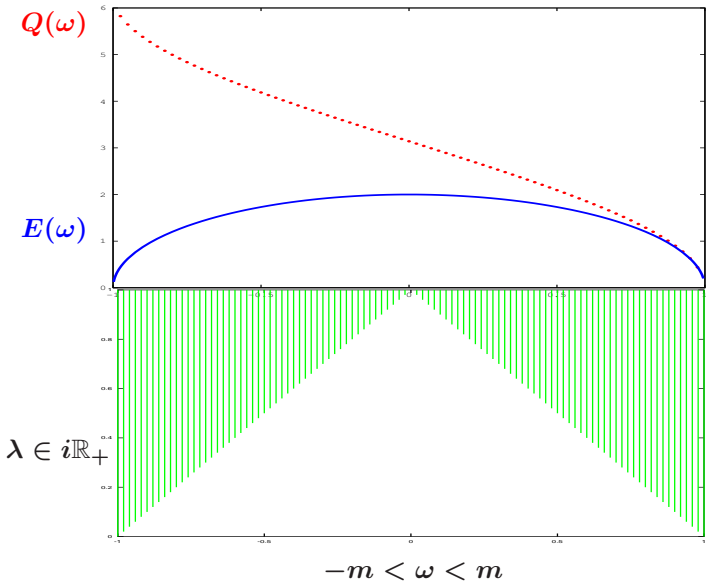


Figure 5: Massive Thirring model with $k = 1$.

TOP: energy and charge as functions of $\omega \in (-1, 1)$.

BOTTOM: The spectrum of the linearization at a solitary wave on the upper half of the imaginary axis.

No nonzero eigenvalues in completely integrable model.

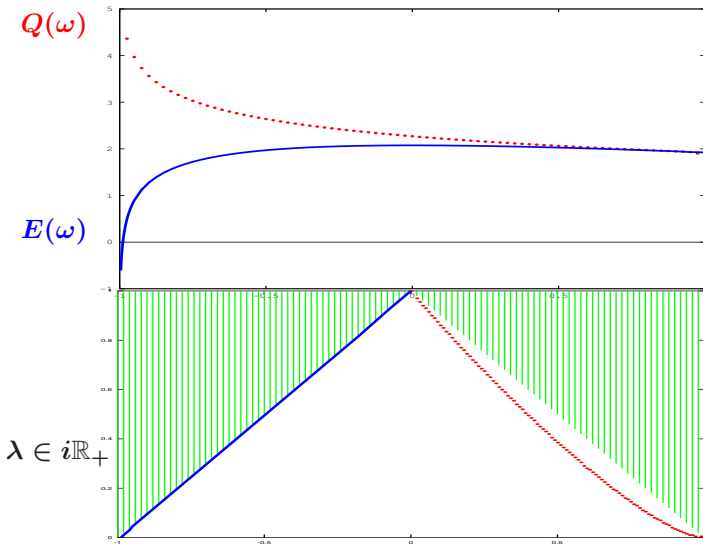


Figure 6: Massive Thirring model with $k = 2$, “charge critical”.
 TOP: energy and charge as functions of $\omega \in (-1, 1)$.
 BOTTOM: The spectrum of the linearization at a solitary wave.

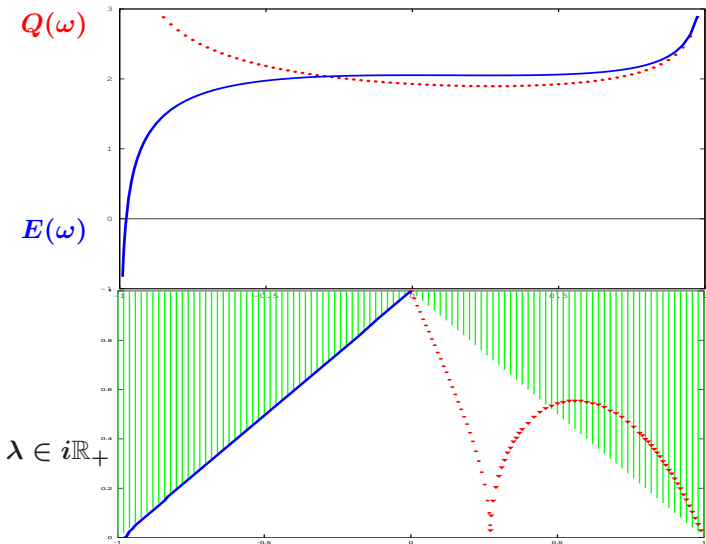


Figure 7: Massive Thirring model with $k = 3$.
 TOP: energy and charge as functions of $\omega \in (-1, 1)$.
 BOTTOM: The spectrum of the linearization at a solitary wave.

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