

Localisation and delocalisation in the parabolic Anderson model

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joint with Stephen Muirhead (Oxford) and Richard Pymar (Birkbeck)

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Parabolic Anderson model

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$$\frac{\partial u}{\partial t} = \Delta u + \xi u$$

with **independent identically distributed random potential** $\{\xi(z) : z \in \mathbb{Z}^d\}$
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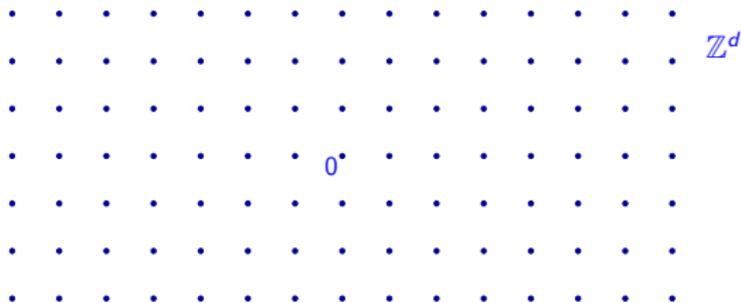
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How does $u(t, \cdot)$ behave as $t \rightarrow \infty$?

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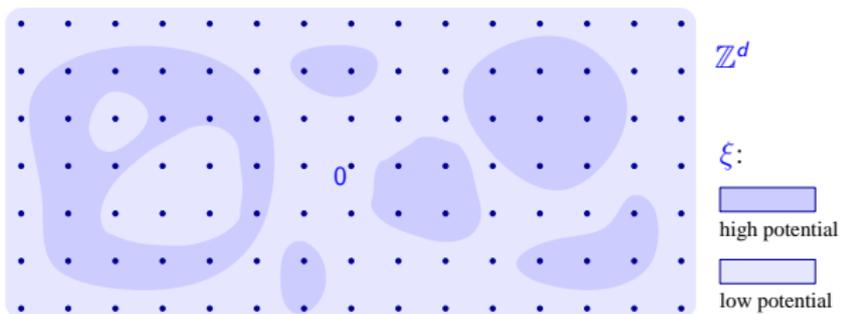
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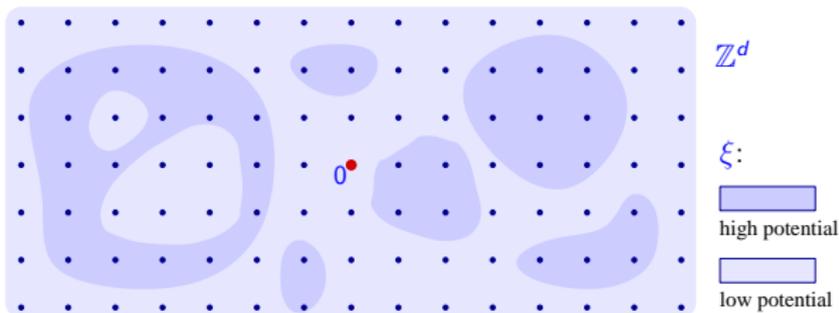
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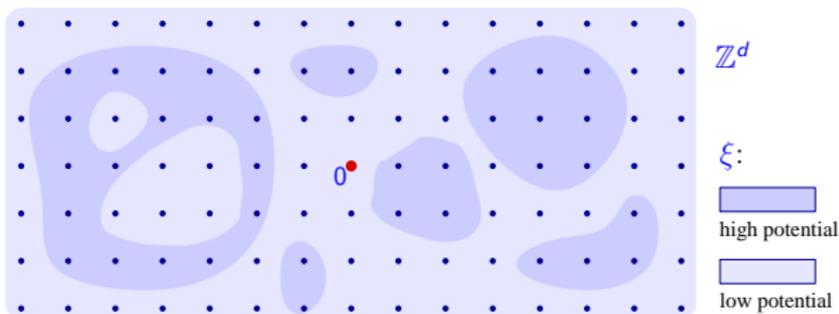
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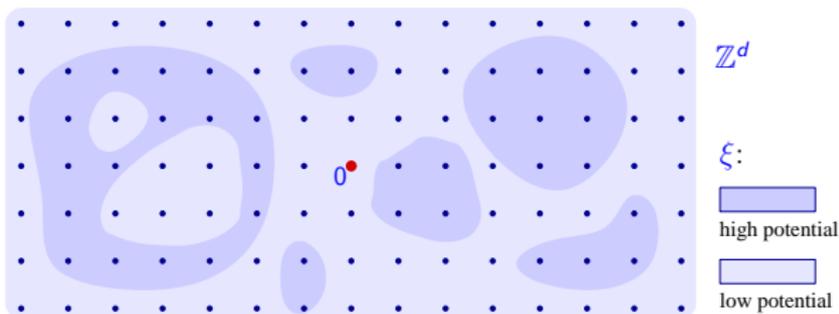
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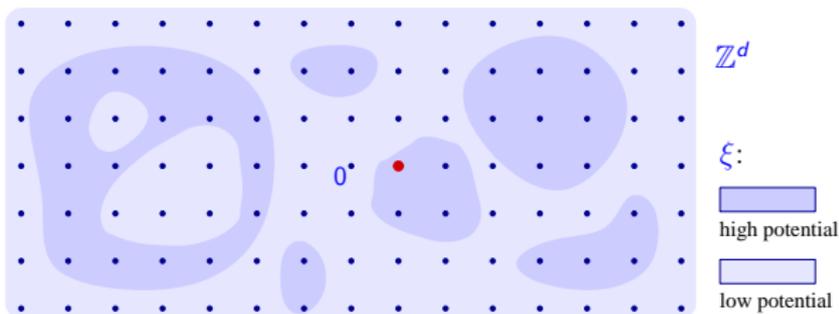
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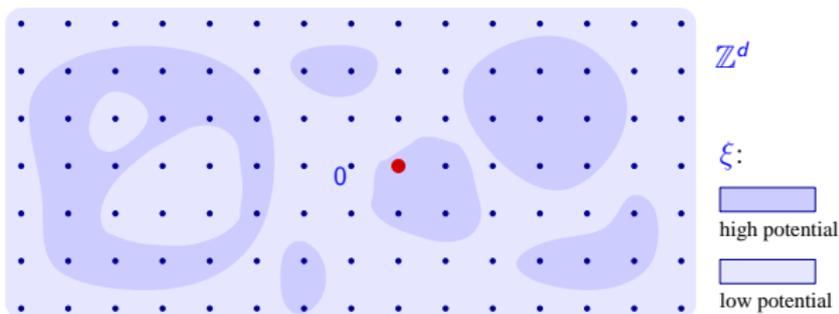
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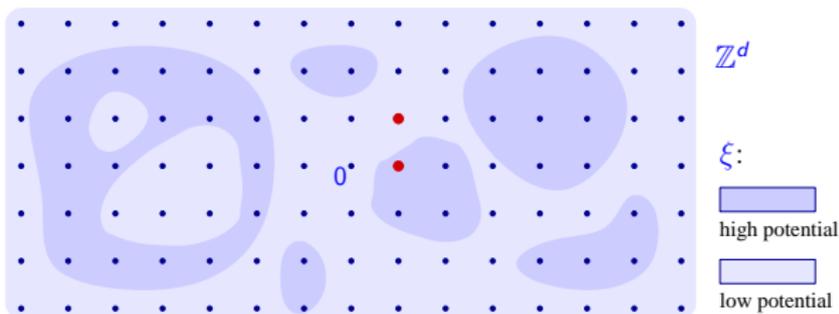
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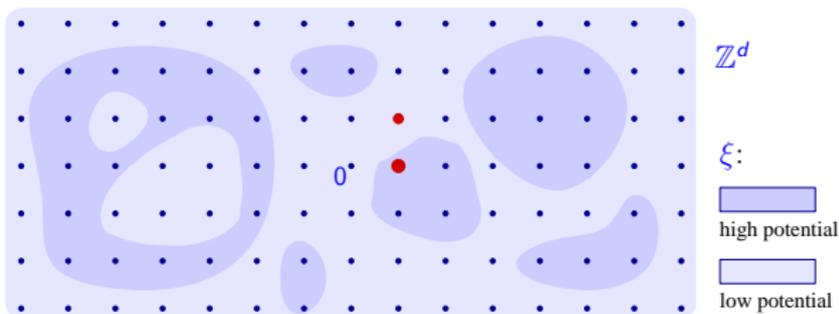
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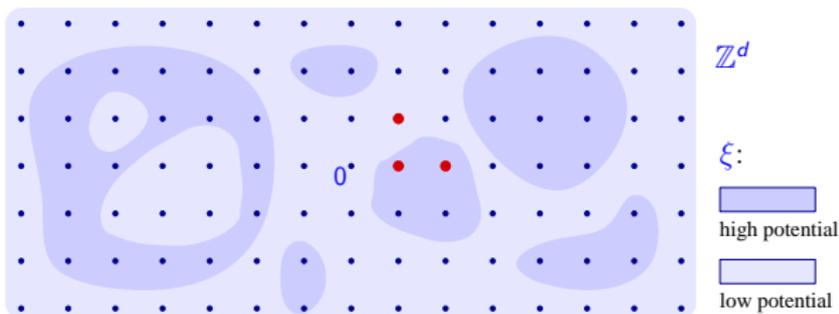
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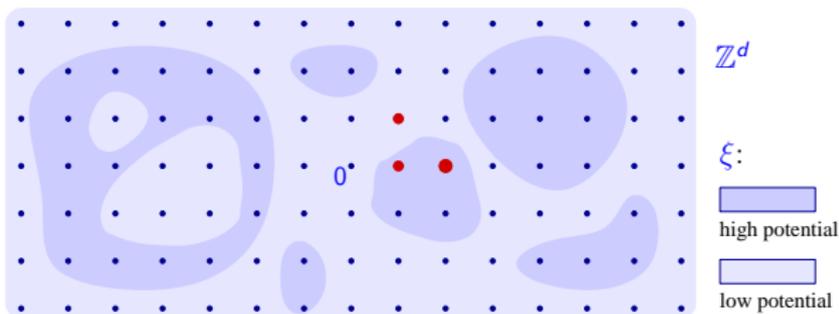
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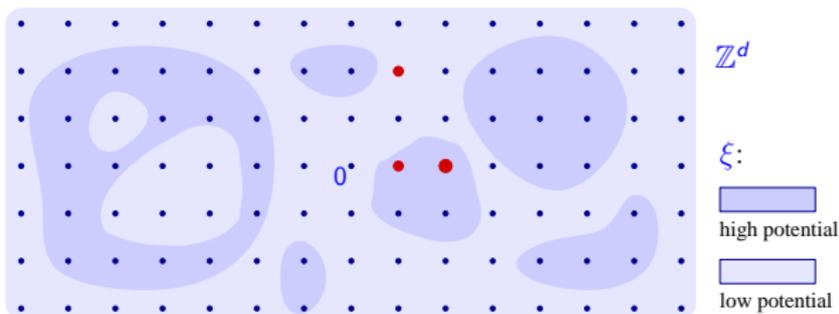
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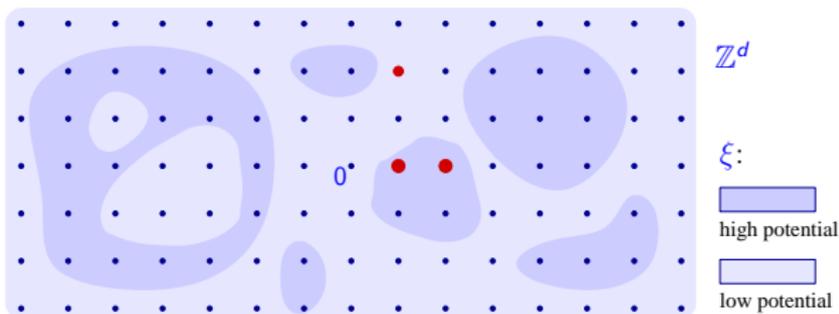
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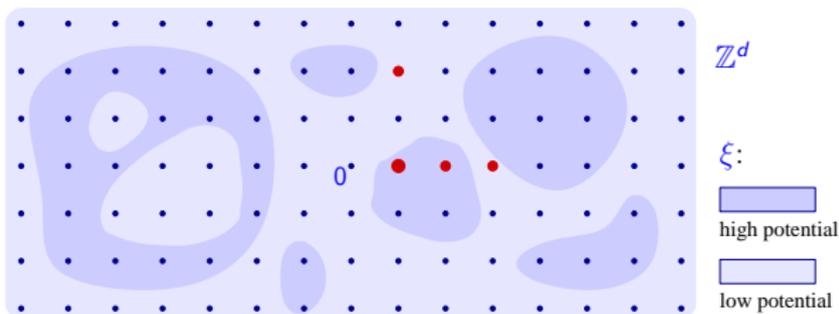
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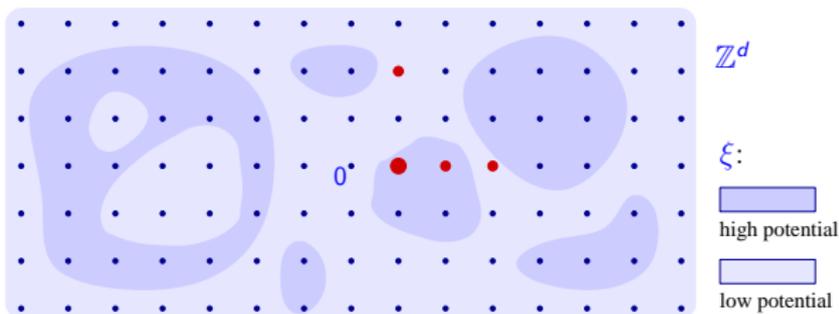
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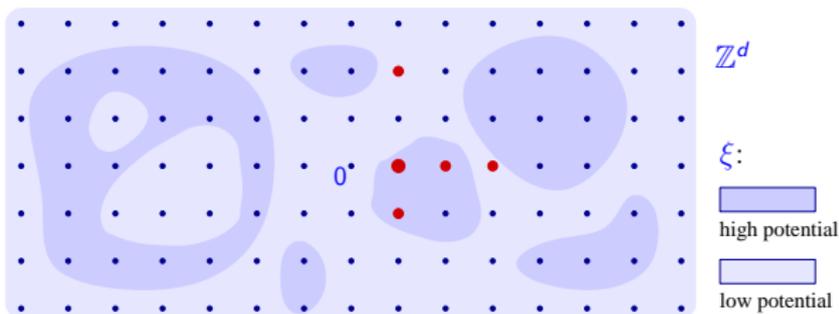
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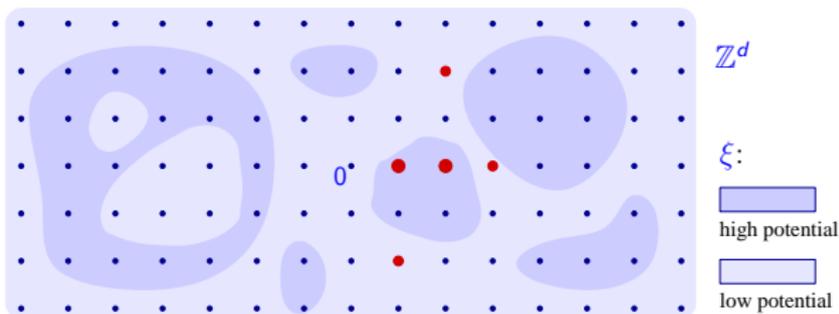
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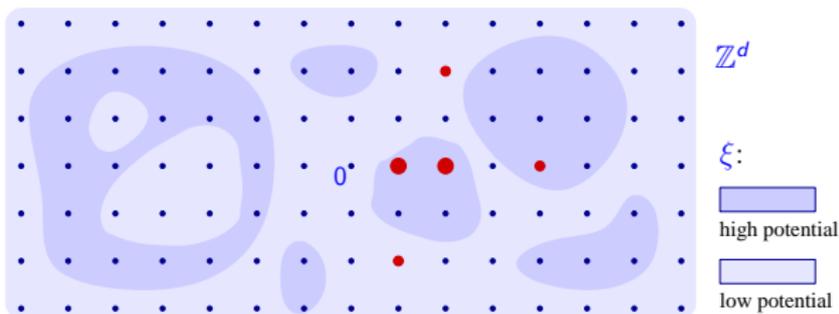
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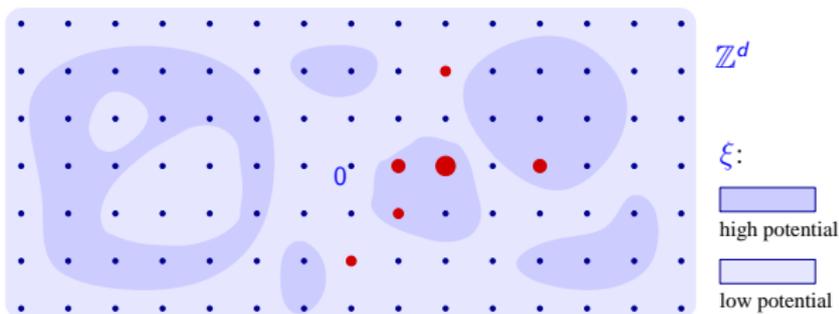
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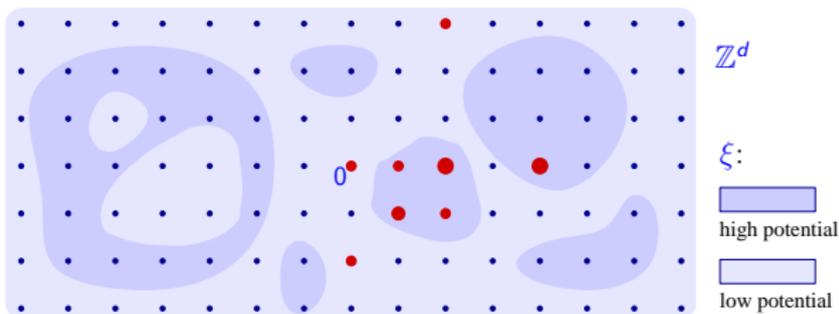
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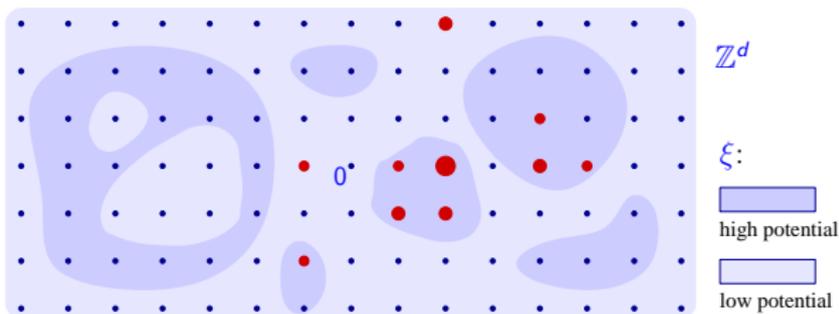
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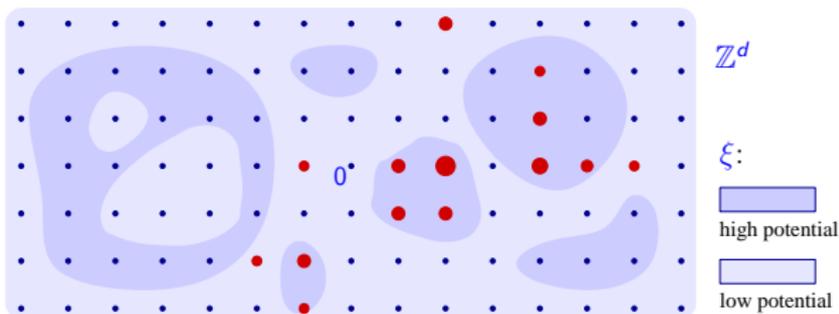
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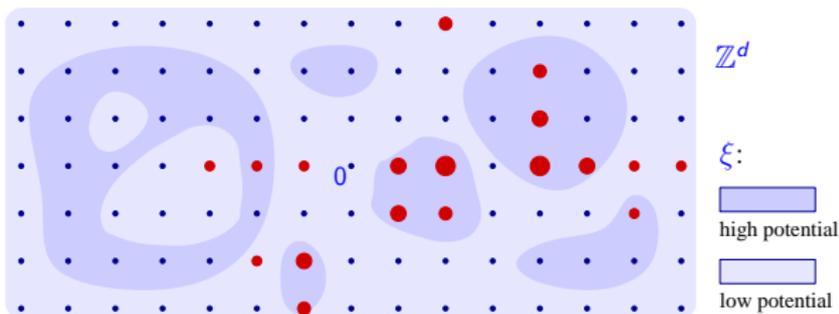
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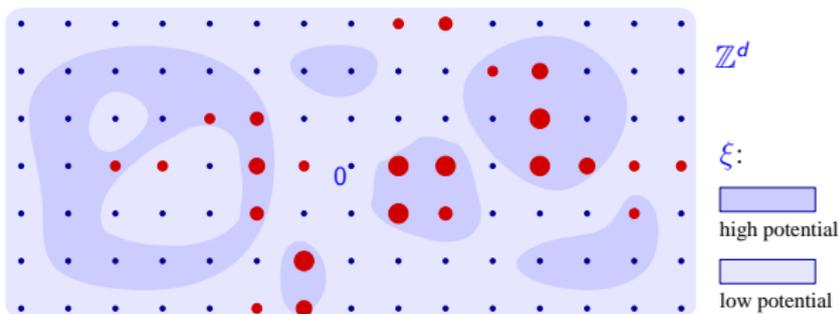
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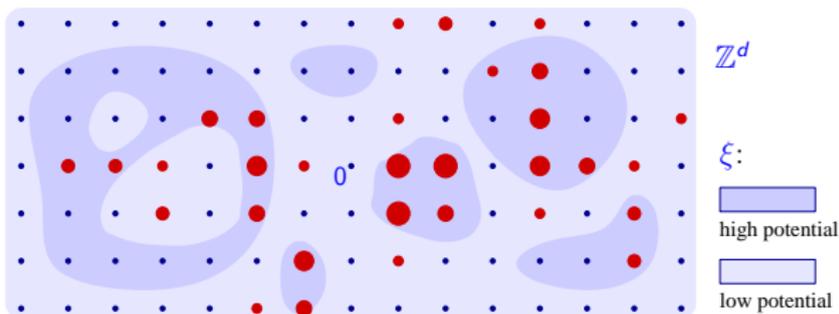
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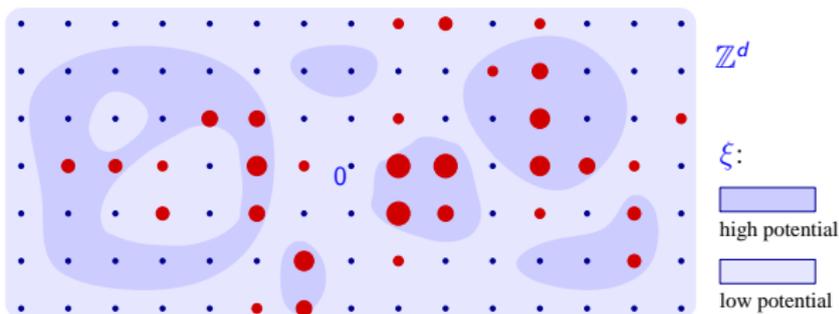
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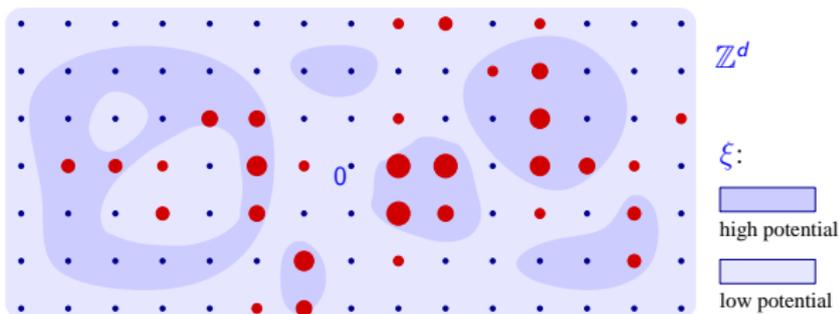
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$u(t, z) = \mathbb{E}N(t, z)$ is the **average** number of particles at time t at site z , **still random**.

Two approaches to study $u(t, z)$

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- **Analytical:** use **Spectral Theory** to analyse the parabolic Anderson equation

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- **Analytical:** use **Spectral Theory** to analyse the parabolic Anderson equation

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- **Probabilistic:** use path analysis to analyse the **Feynman–Kac Formula**

$$u(t, z) = \mathbb{E} \left\{ e^{\int_0^t \xi(X_s) ds} \mathbf{1}_{\{X_t=z\}} \right\},$$

where (X_s) is a continuous-time random walk starting at zero.

Heat equation

The propagation of temperature $u(t, x)$ at time t at the point $x \in \mathbb{R}$ is described by

$$\frac{\partial u}{\partial t} = \Delta u.$$

If the initial temperature is δ_0 then

$$u(t, x) = \frac{1}{\sqrt{4\pi t}} e^{-\frac{x^2}{4t}}.$$

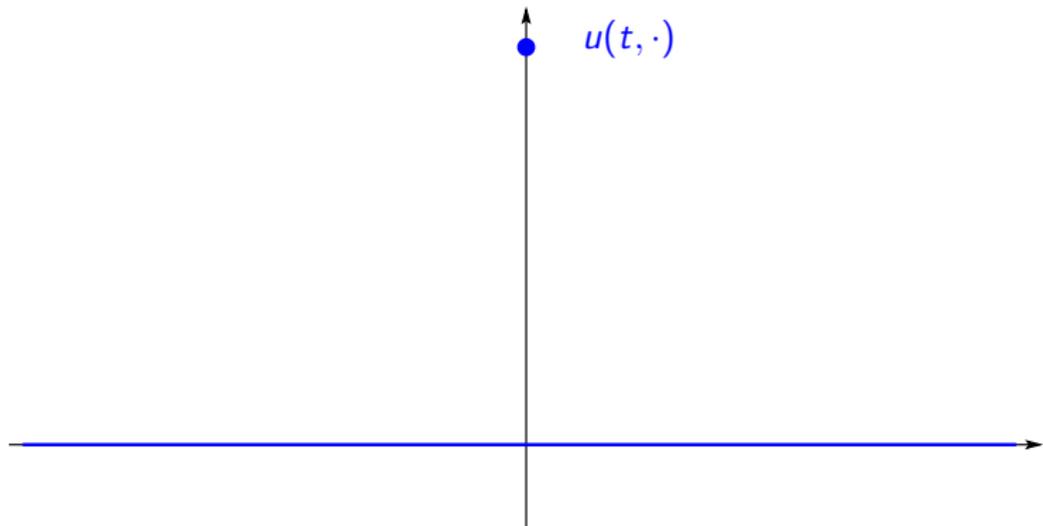
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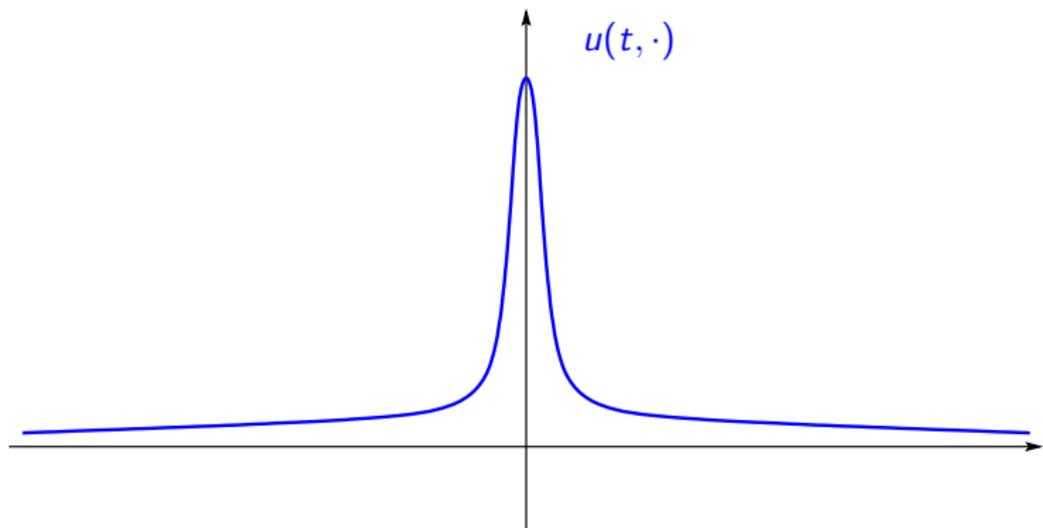
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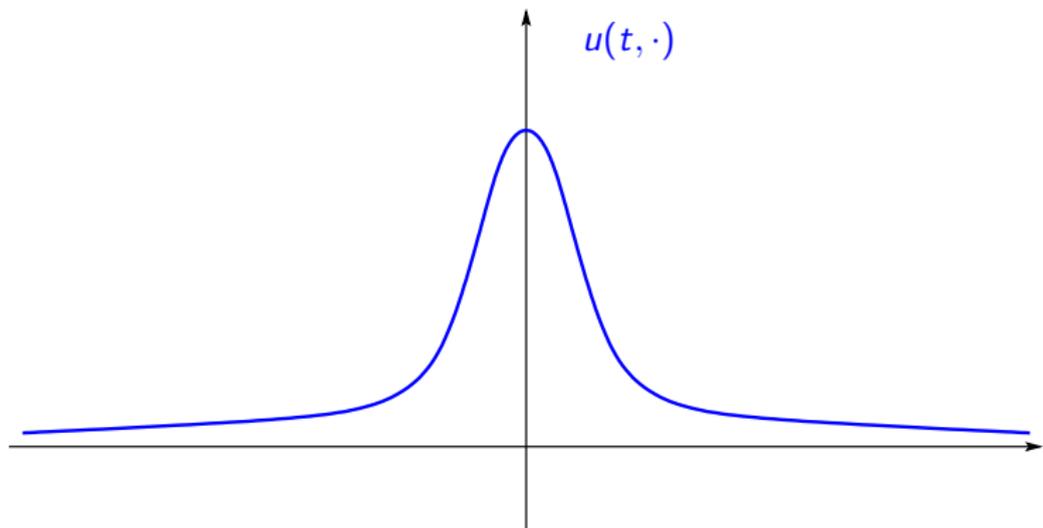
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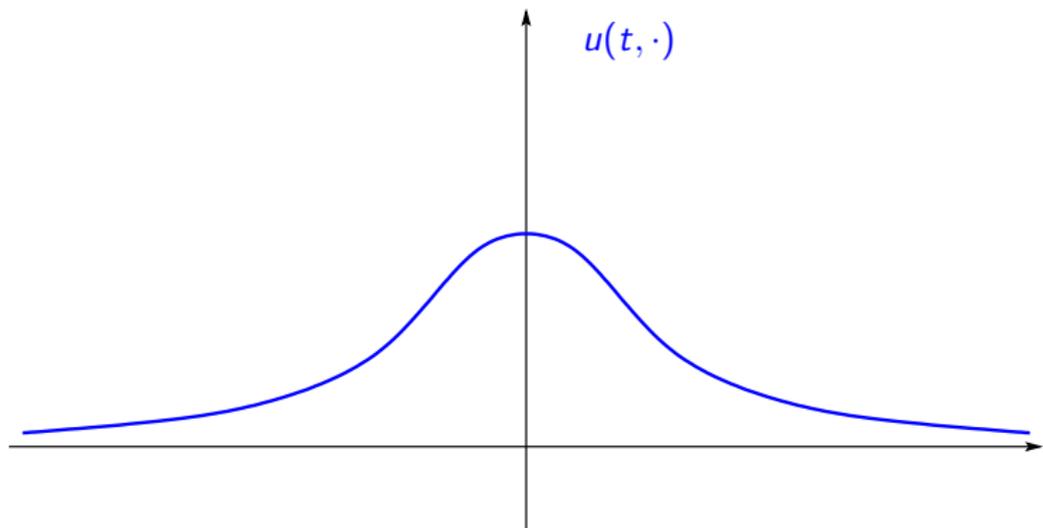
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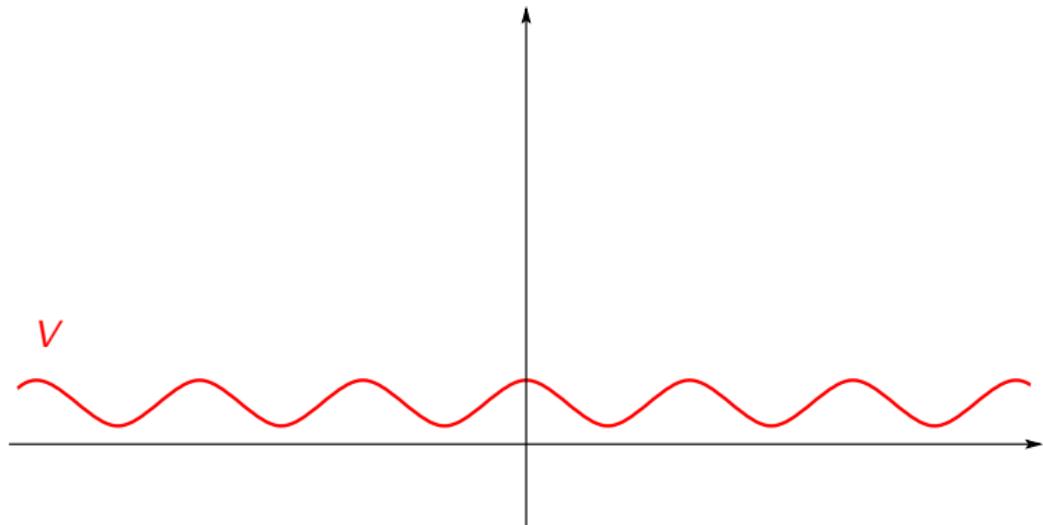


Heat equation with a potential

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where $V : \mathbb{R} \rightarrow \mathbb{R}$ is a reasonably nice **potential**.



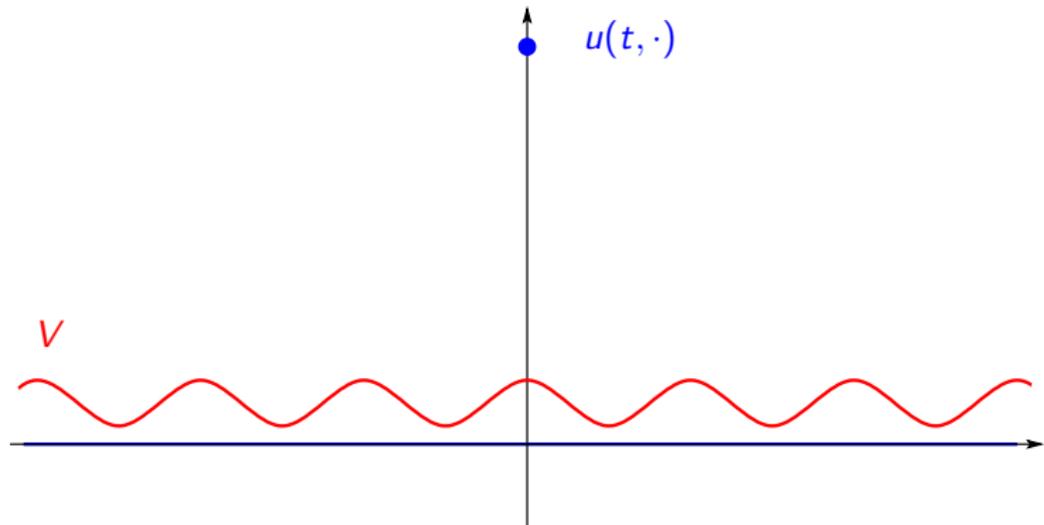
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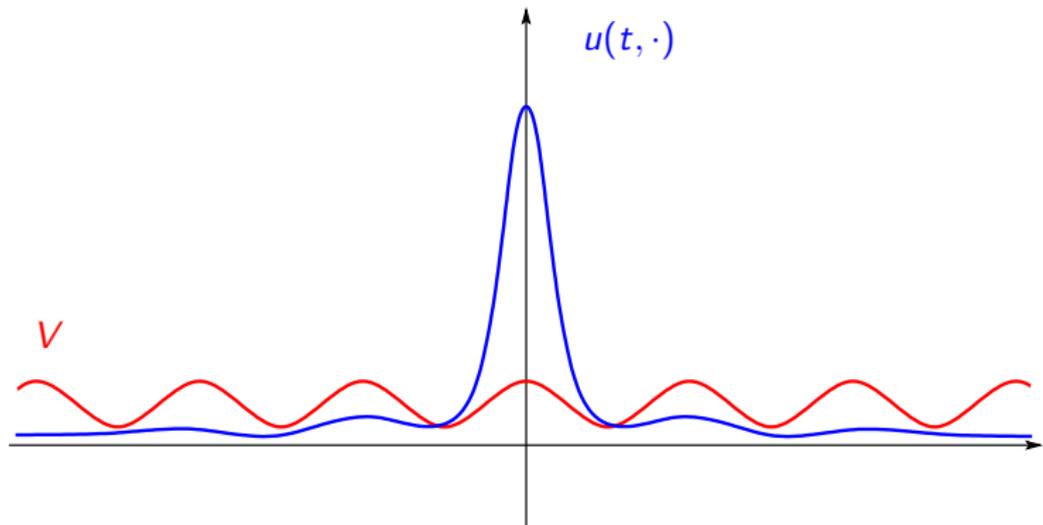
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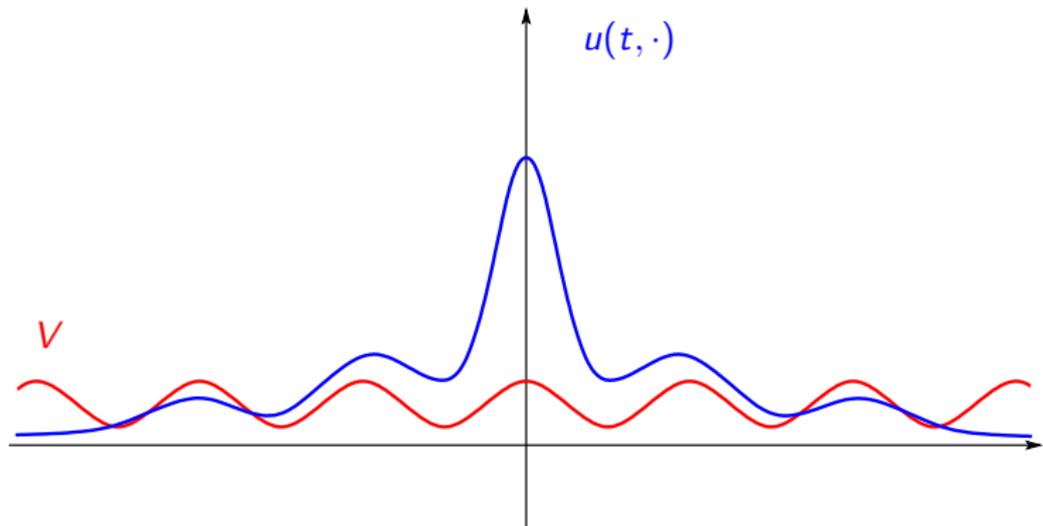
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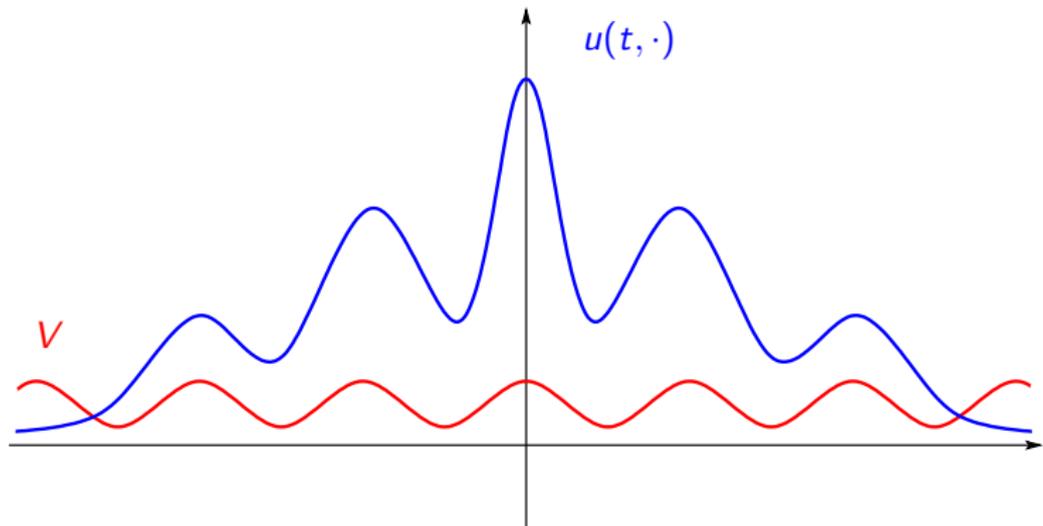
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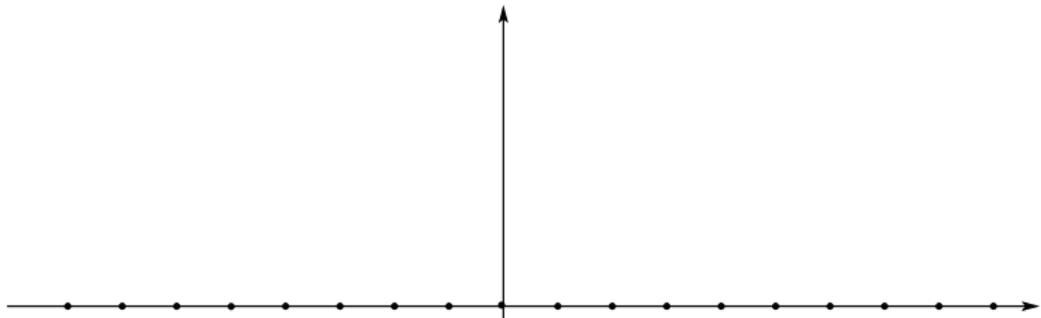
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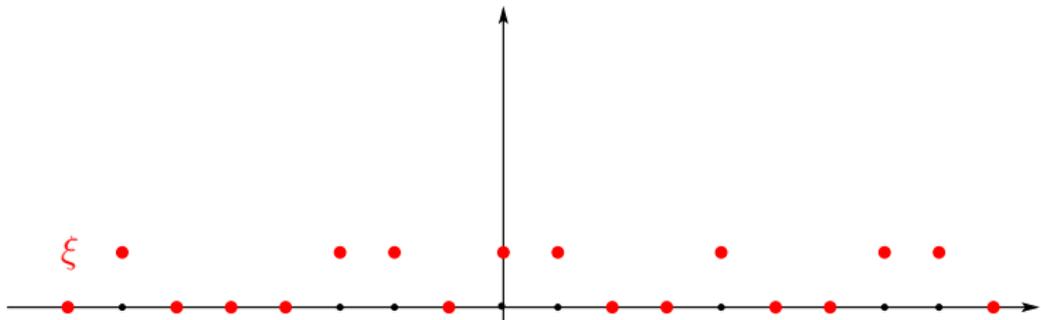


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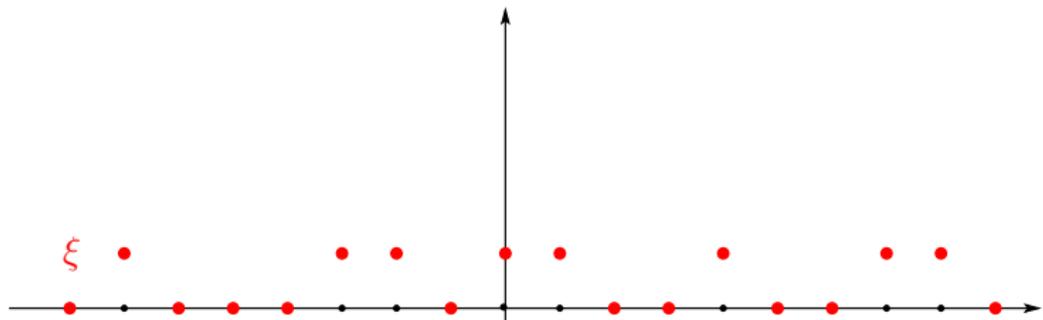


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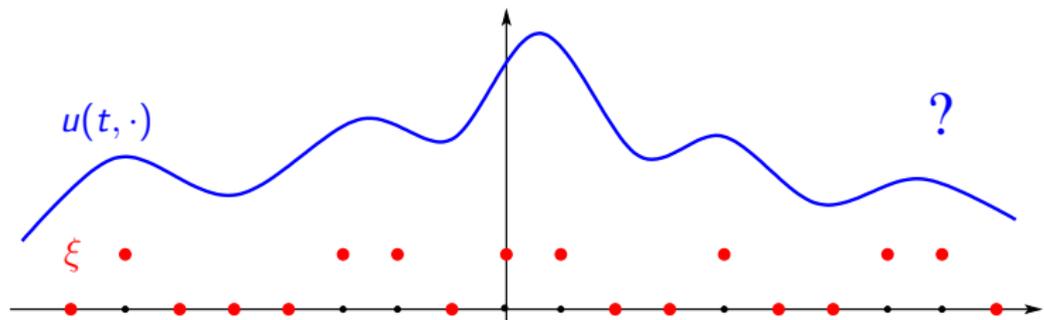
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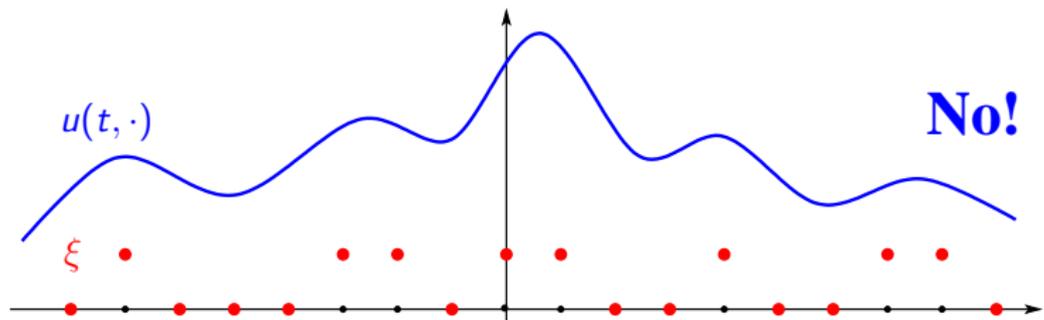
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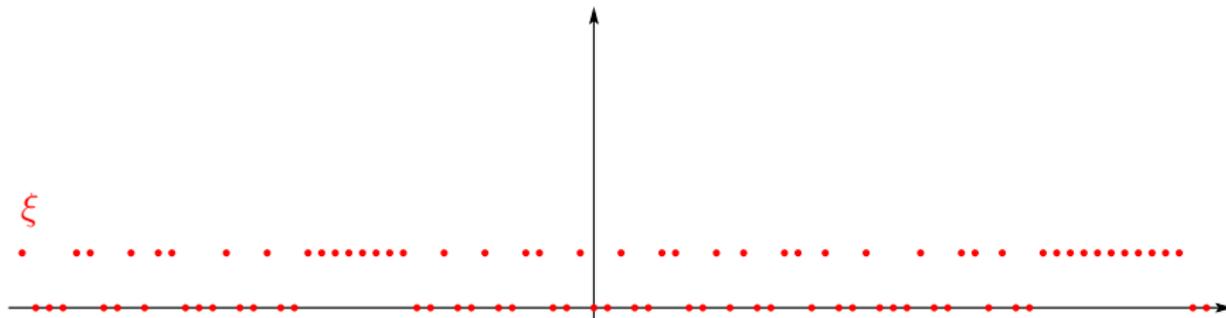
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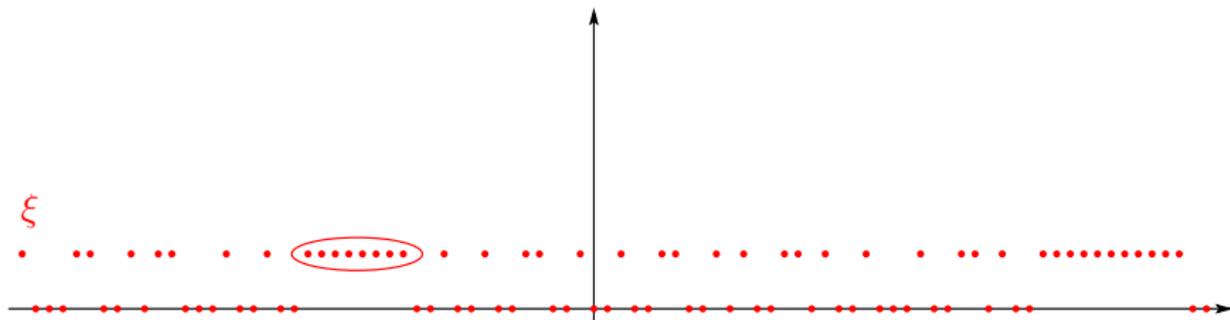
Why not?

Bernoulli potential:



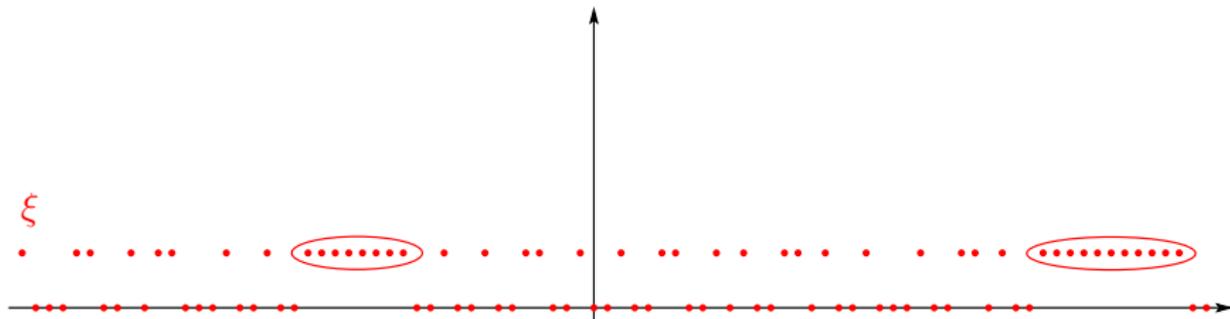
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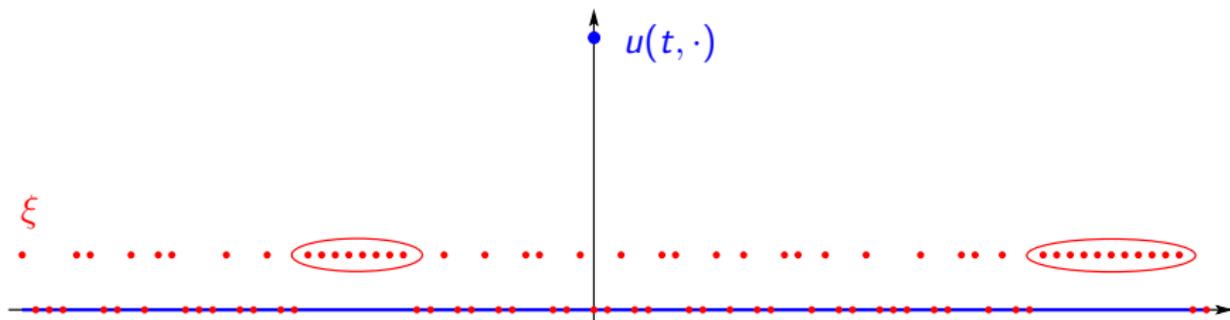
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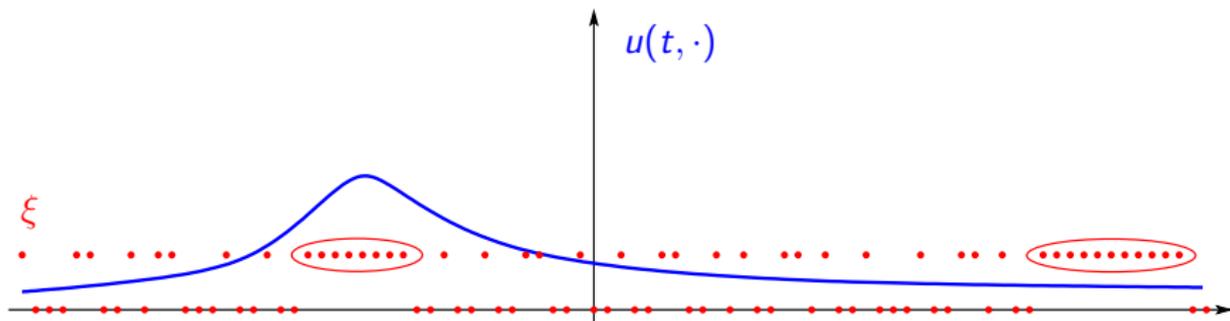
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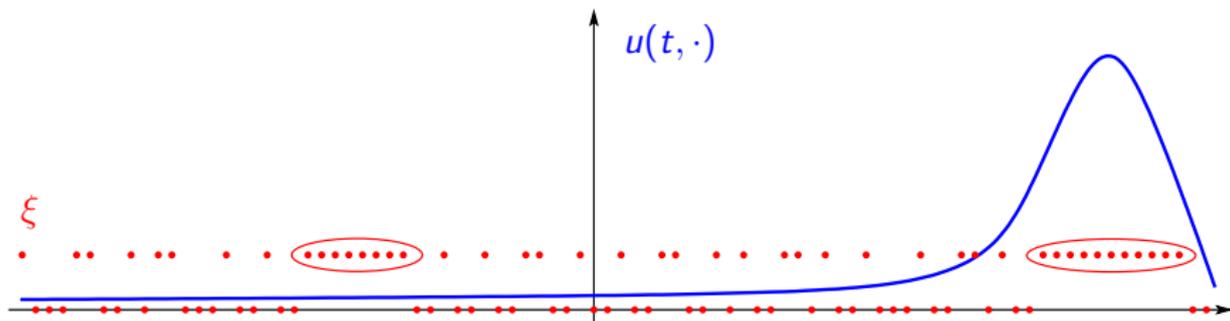
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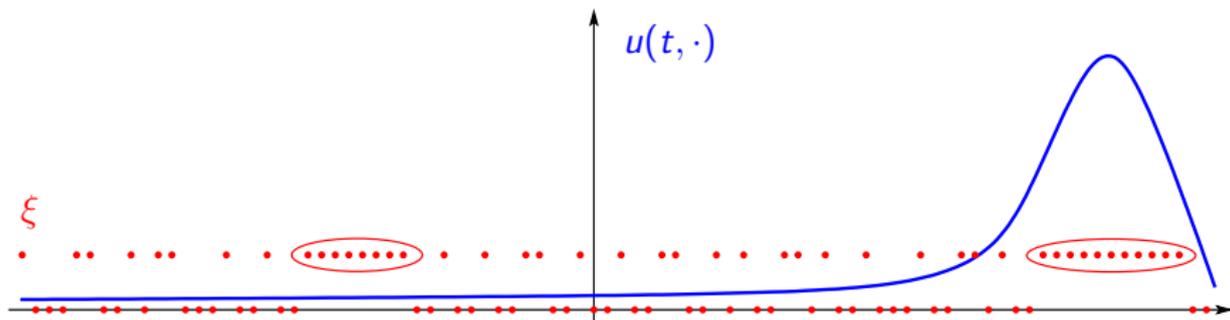
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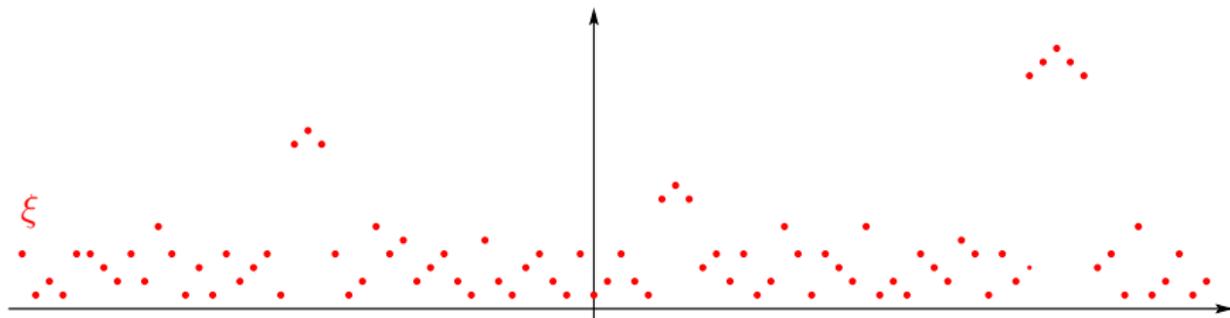


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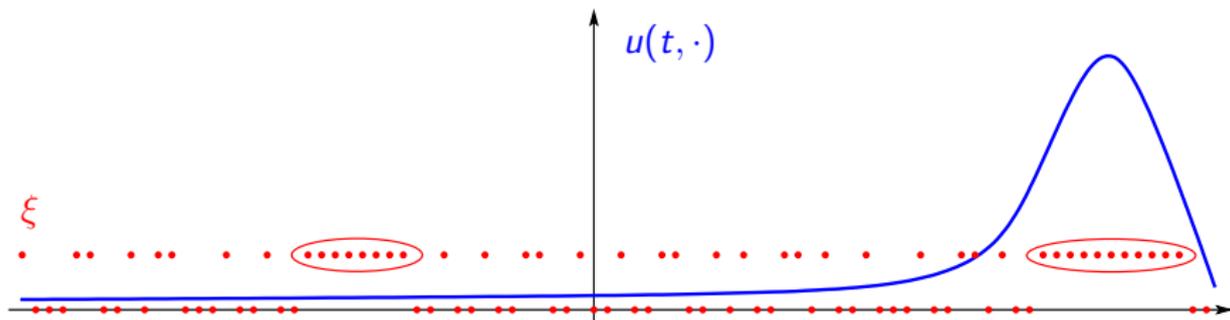


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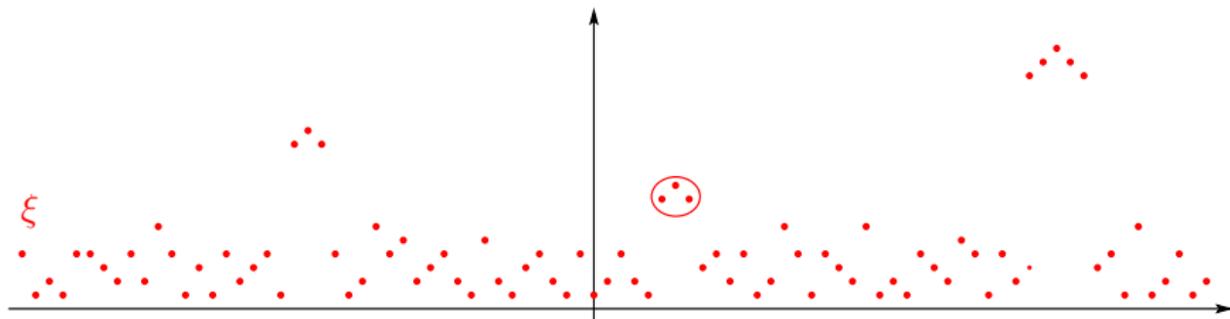


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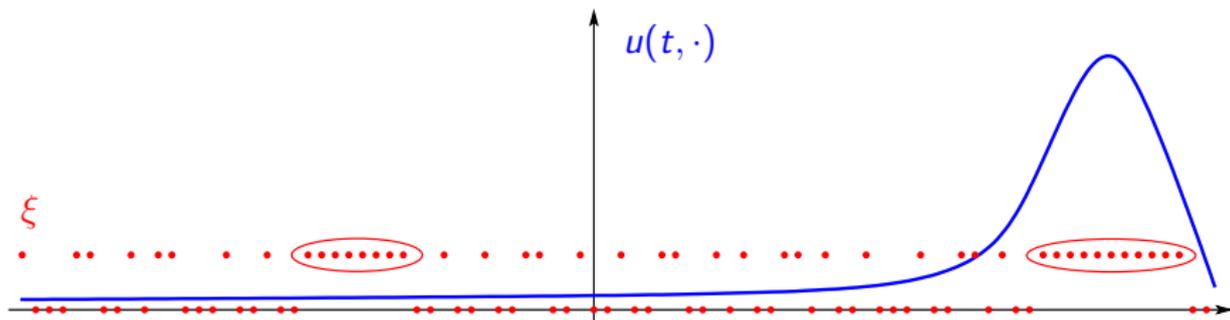


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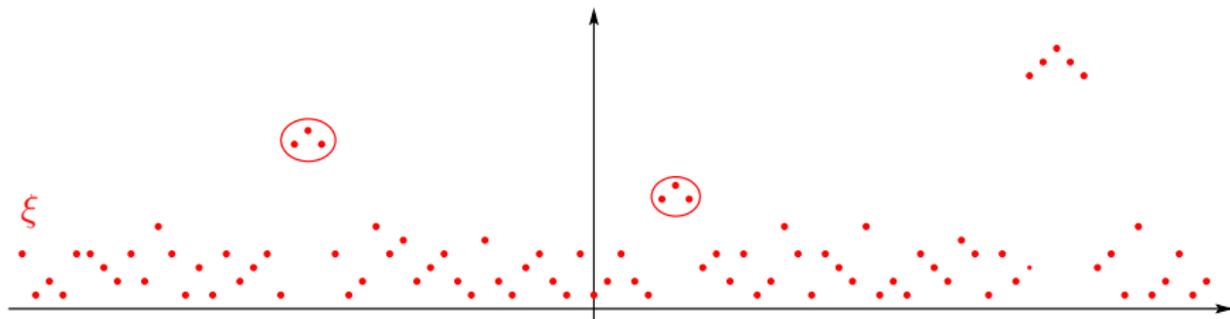


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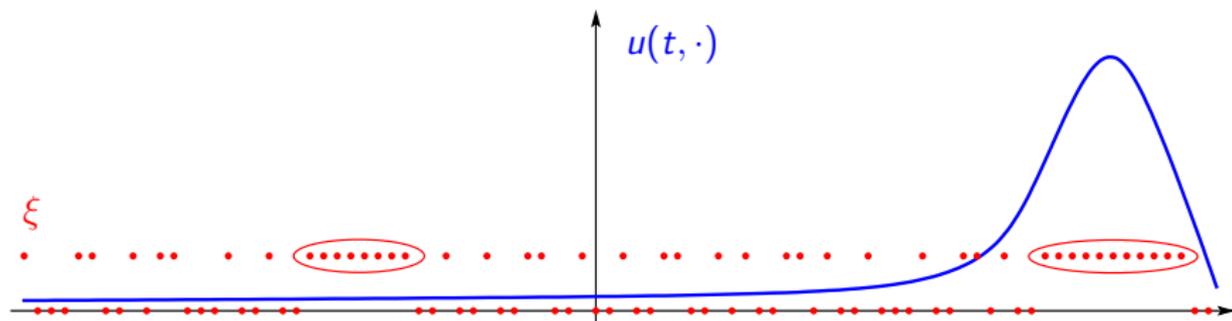


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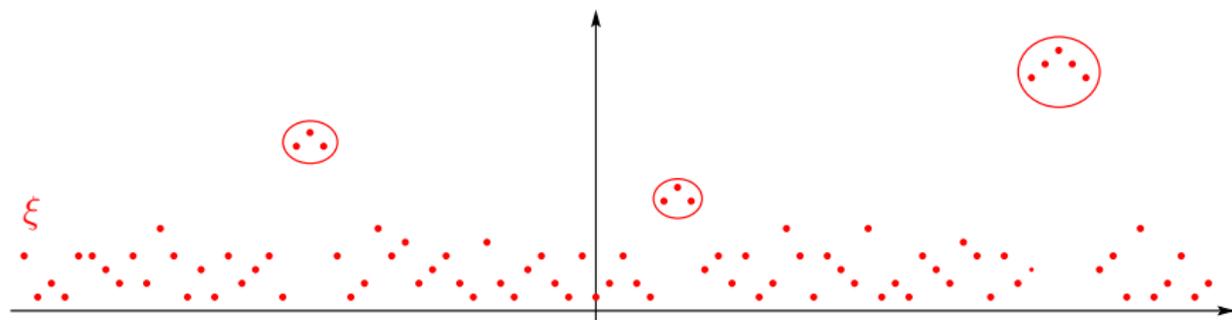


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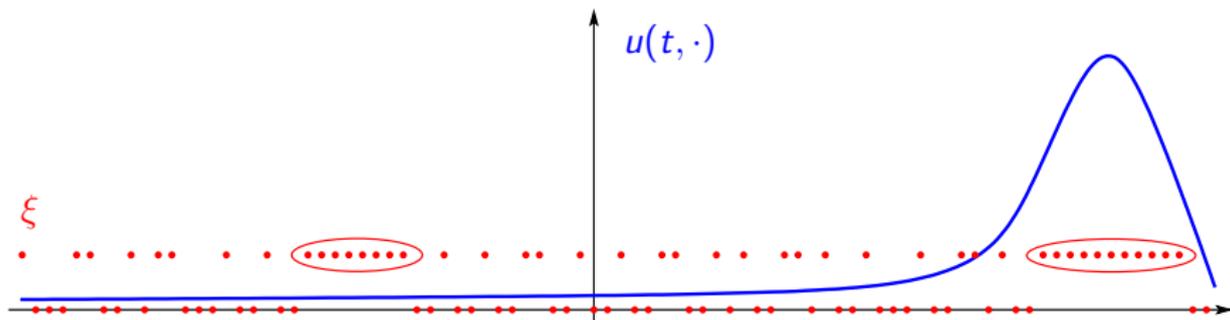


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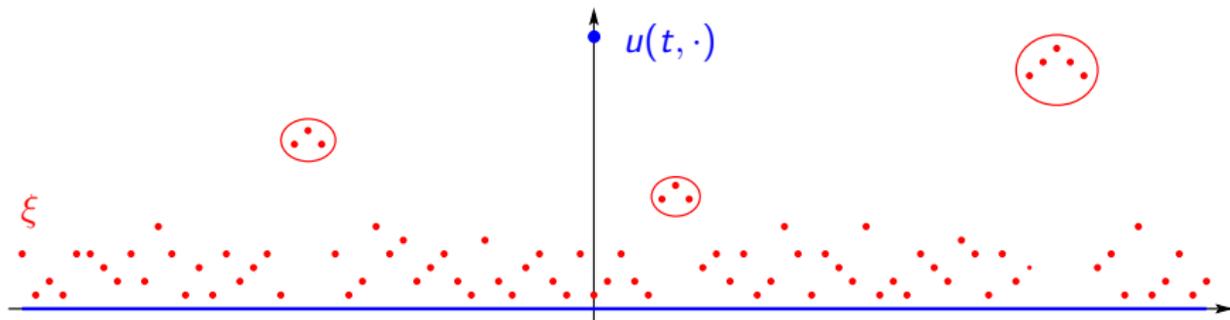


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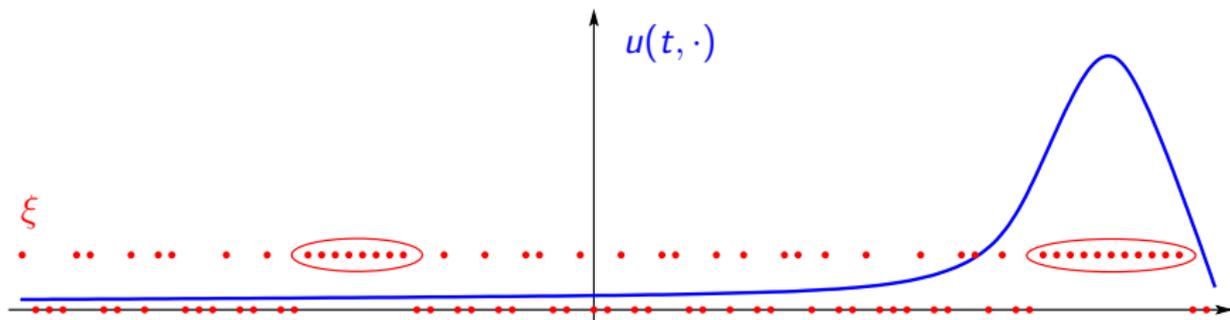


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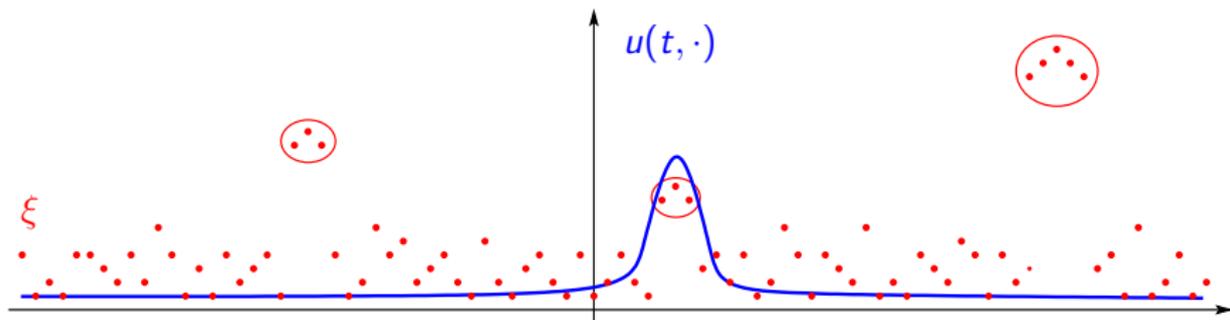


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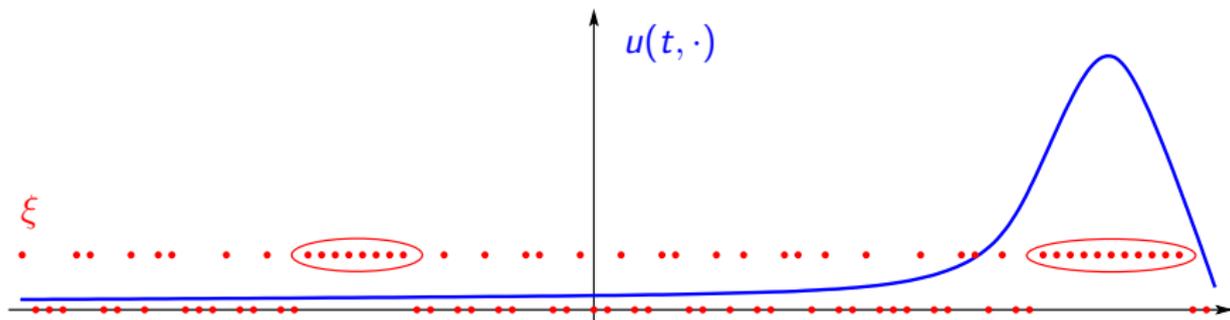


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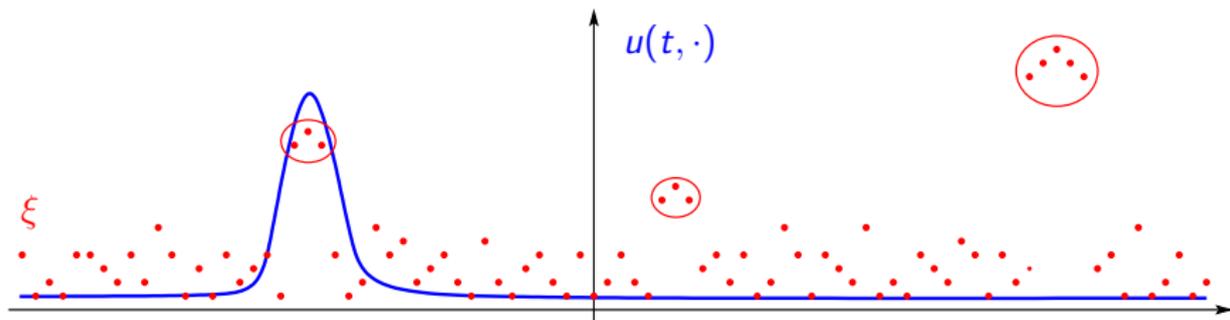


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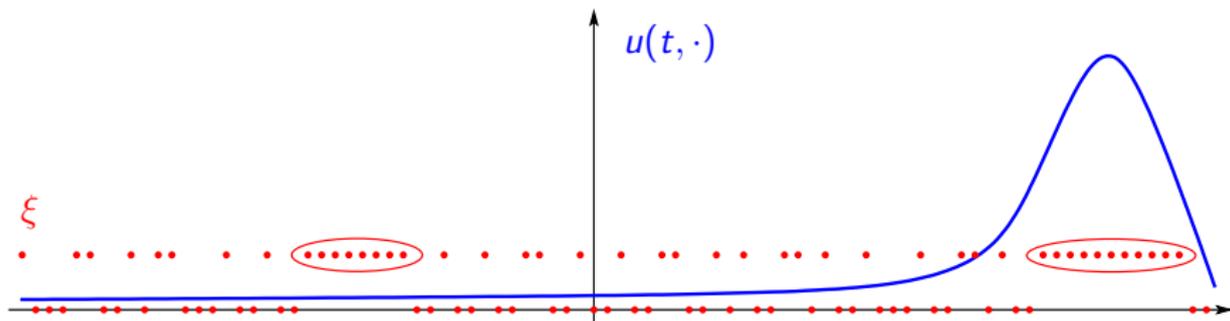


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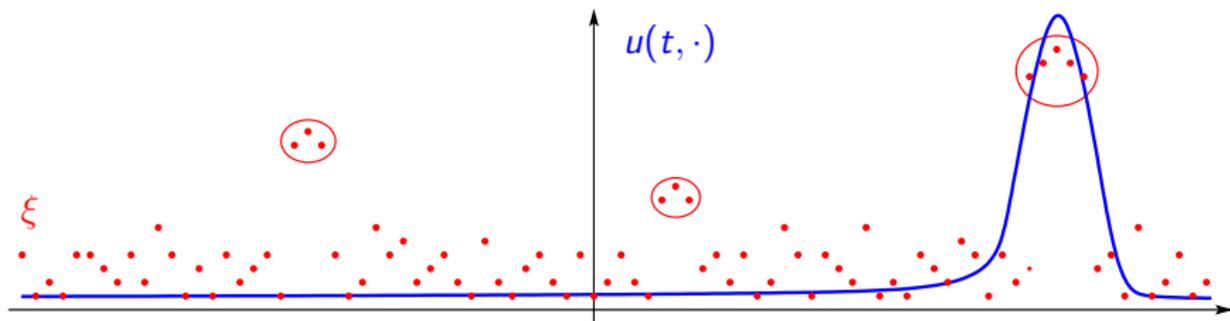


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What is known for **unbounded** potentials?

- **Pareto**: $P(\xi(0) > x) = x^{-\alpha}$, $\alpha > d$
- **Weibull**: $P(\xi(0) > x) = \exp\{-x^\gamma\}$, $\gamma > 0$
- **Double-exponential**: $P(\xi(0) > x) = \exp\{-e^{x/\rho}\}$, $\rho > 0$
- 'Almost bounded' — quite different, not in this talk

What is known

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- [S., Twarowski '12] Weibull with $\gamma < 2$
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Main question of this talk: How can we break this and make the solution spread between two (or more) independent locations?

Model: the PAM with duplication

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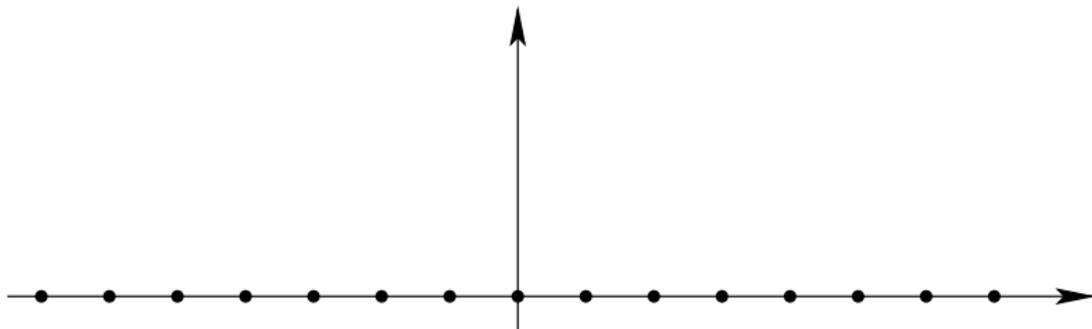
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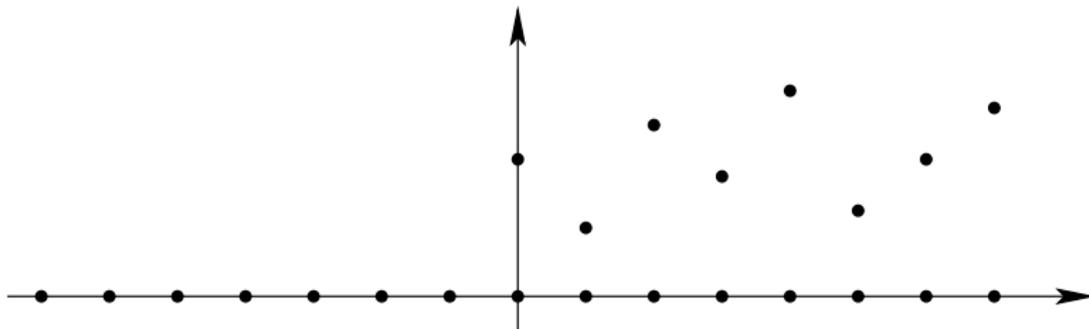
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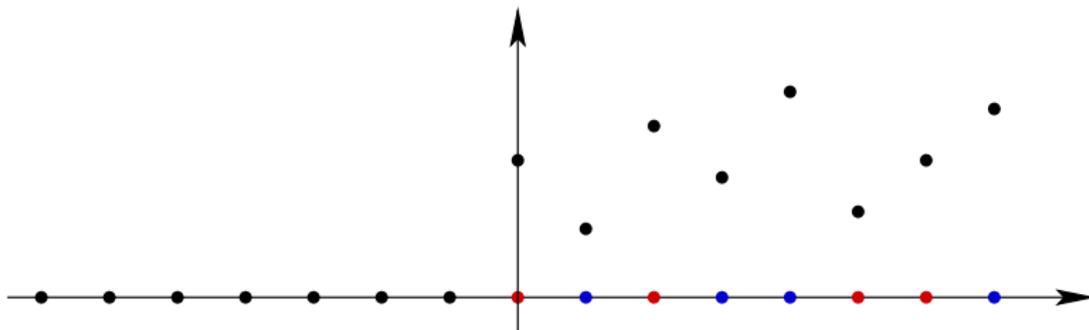
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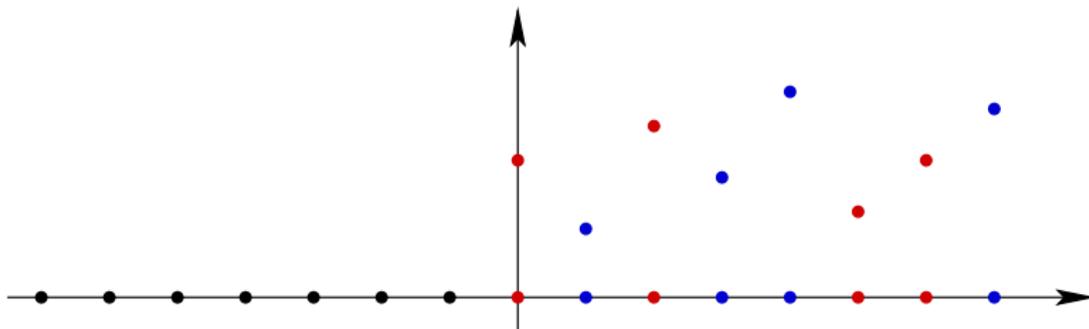
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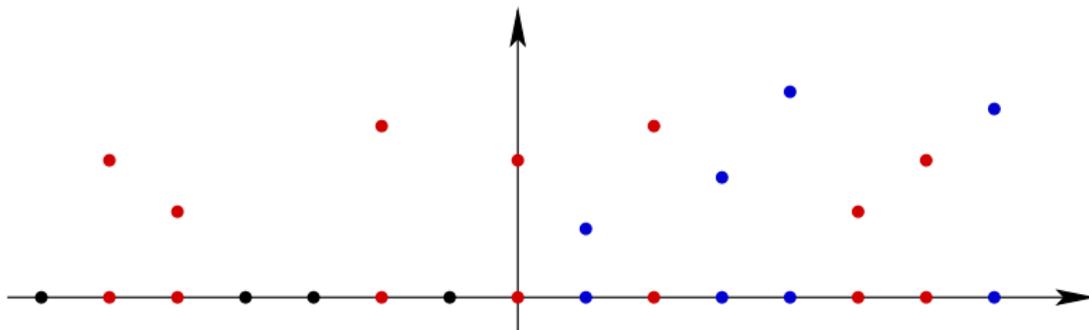
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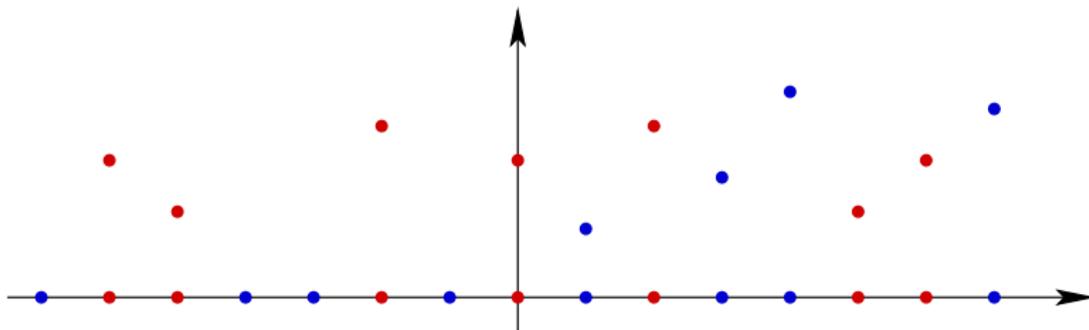
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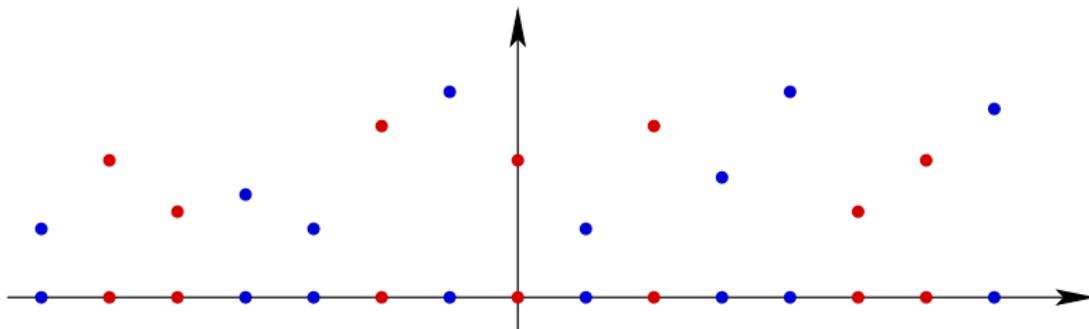
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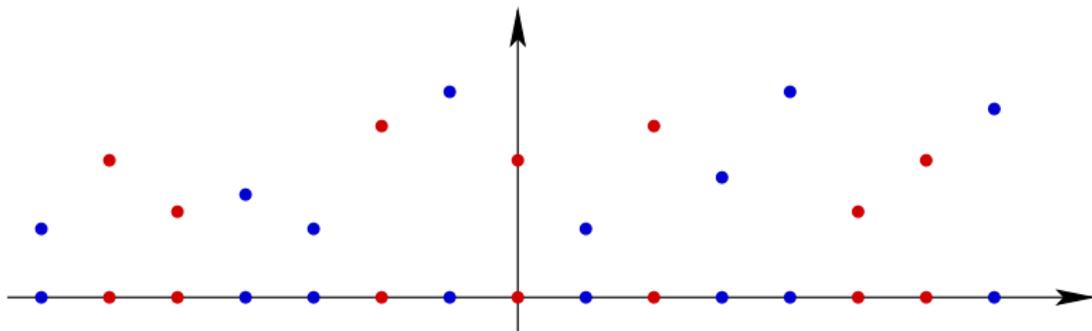
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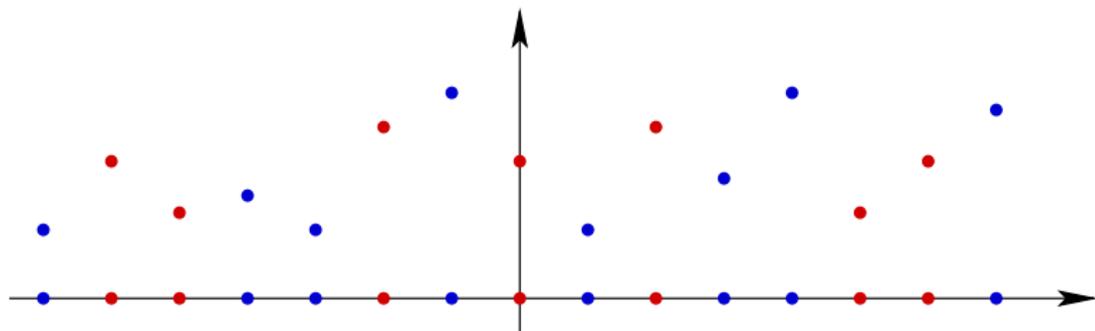


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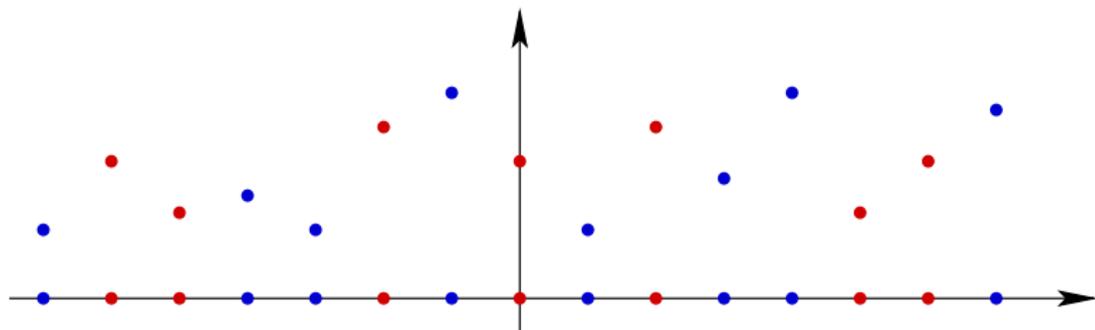
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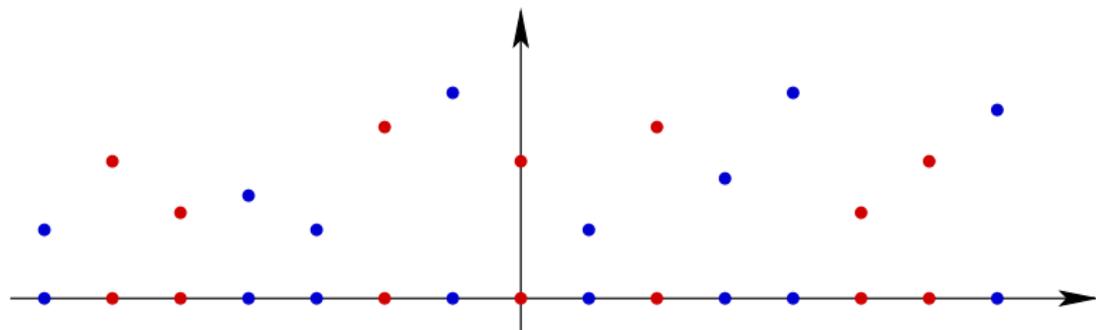
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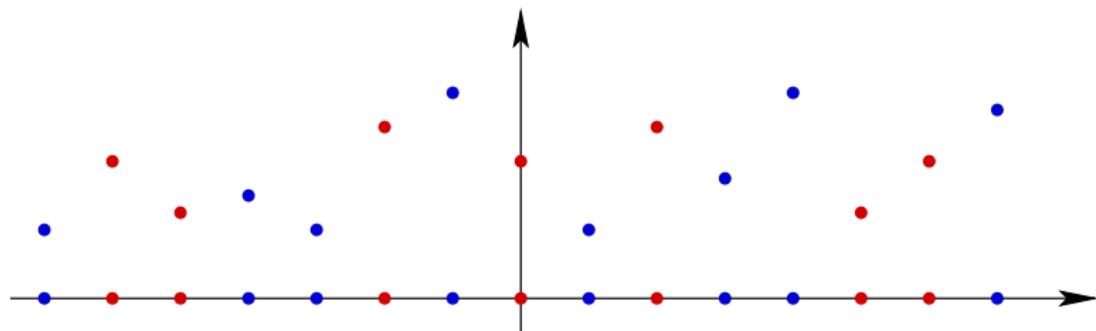
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Denote the total mass of the solution by

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Let $\alpha \geq 2$. As $t \rightarrow \infty$, one point

$$\frac{u(t, Z_t)}{U(t)} \rightarrow 1 \quad \text{in probability.}$$

Answers

Theorem 1 (Muirhead, Pymar, S. '16)

Let $1 < \alpha < 2$. Conditionally on **no duplication at Z_t** , as $t \rightarrow \infty$, **one point**

$$\frac{u(t, Z_t)}{U(t)} \rightarrow 1 \quad \text{in probability.}$$

Conditionally on the **duplication at Z_t** , as $t \rightarrow \infty$, **two points, each with a random amount of mass**

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where \mathcal{P}_t^\pm are the sets of paths on \mathbb{Z} starting at 0 and ending at $\pm Z_t$.

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$$U(t, y) = \sum_{i=0}^n e^{tc_i - 2t} \prod_{\substack{k=0 \\ k \neq i}}^n \frac{1}{c_i - c_k}$$

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In particular, for $1 < \alpha < 2$ we have

$$u(t, \pm Z_t) \sim e^{t\xi(Z_t) - 2t} \prod_{k=0}^{Z_t} \frac{1}{\xi(Z_t) - \xi(\pm k)}$$

$$1 < \alpha < 2$$

$$\frac{u(t, Z_t)}{u(t, -Z_t)} \sim \prod_{k=1}^{Z_t} \frac{1}{\xi(Z_t) - \xi(k)} : \prod_{k=1}^{Z_t} \frac{1}{\xi(Z_t) - \xi(-k)}$$

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$$= \exp \left\{ - \sum_{k:\text{non-dupl}} \log \left(1 - \frac{\xi(k)}{\xi(Z_t)} \right) + \sum_{k:\text{non-dupl}} \log \left(1 - \frac{\xi(-k)}{\xi(Z_t)} \right) \right\}$$

$\sum_{i=1}^n X_i \approx n\mu + n^{1/\alpha} \mathcal{N}$
\uparrow
stable Pareto(α)

Insightful cheating:

$$- \sum_{k:\text{non-dupl}} \log \left(1 - \frac{\xi(\pm k)}{\xi(Z_t)} \right) \approx \frac{1}{\xi(Z_t)} \sum_{k:\text{non-dupl}} \xi(\pm k) \approx \underbrace{\frac{\mu q |Z_t|}{\xi(Z_t)}}_{\text{LLN}} + \underbrace{\frac{|Z_t|^{1/\alpha}}{\xi(Z_t)} \mathcal{N}^\pm}_{\text{fluctuations}}$$

same for $\pm Z_t$

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Hence everything is determined by

$$\frac{|Z_t|^{1/\alpha}}{\xi(Z_t)} \asymp$$

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The scale of fluctuations remains **finite for all values** of $1 < \alpha < 2$.

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