

# **Occurrence of normal and anomalous diffusion in polygonal billiard channels**



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- Motivation
- Properties of billiards
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# Introduction

- Definition:
  - Fixed, hard obstacles – scatterers

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- Fixed, hard obstacles – **scatterers**
- Non-interacting point particles collide elastically

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  - electron gas in metal (Lorentz 1905)

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- Motivation: transport processes
  - electron gas in metal (Lorentz 1905)
  - hard-sphere fluid (Sinai 1960s)
  - one of simplest physical systems with macroscopic transport

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- Properties depend on geometry of scatterers:

circular	polygonal
Lorentz gas	Ehrenfest wind-tree model

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Lyapunov exponent $> 0$ ("chaotic")	Lyapunov exponent $= 0$ ("non-chaotic")

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- Necessary microscopic conditions for macroscopic transport?
- Corners separate nearby trajectories: "randomising" effect

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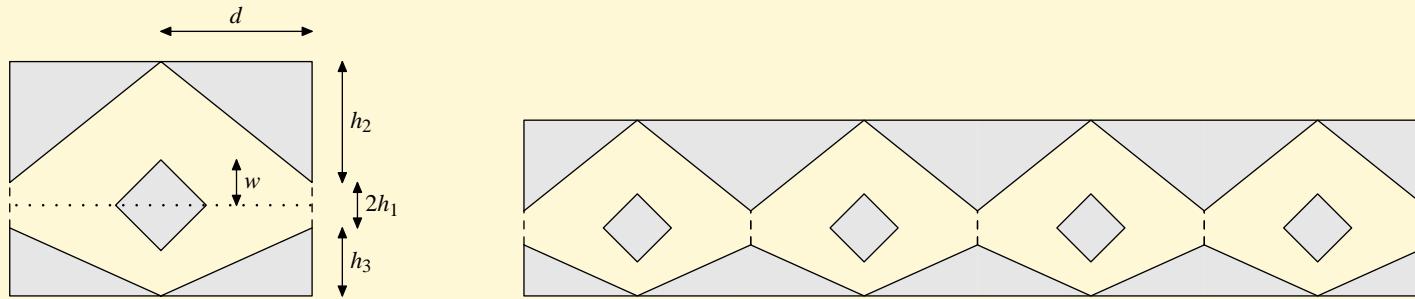
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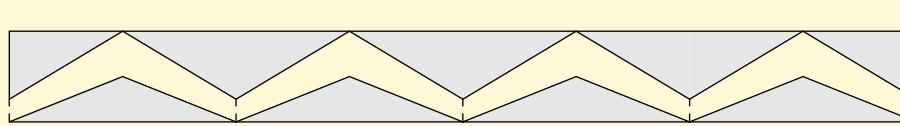
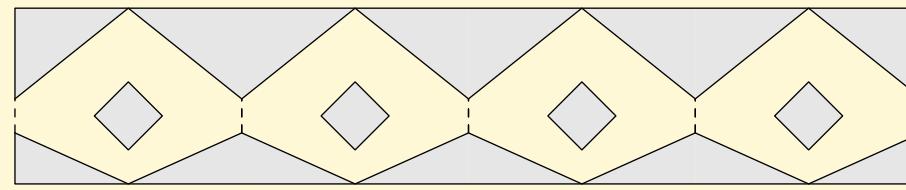
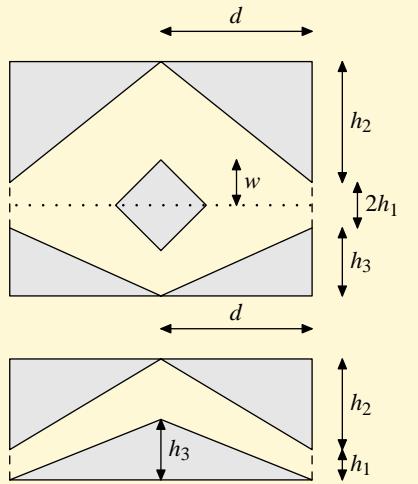
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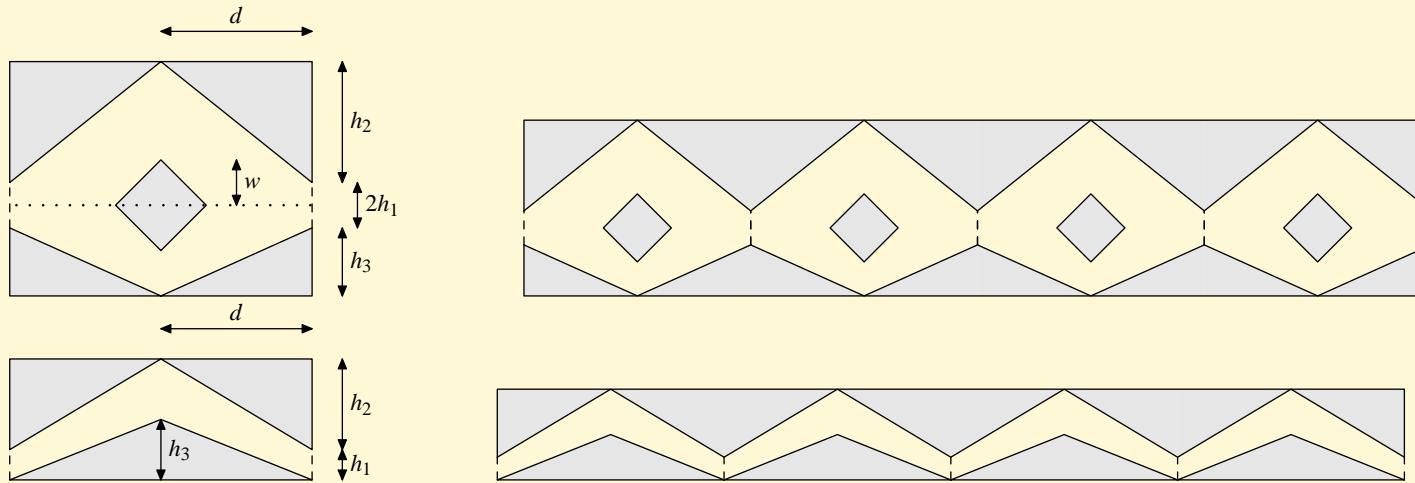
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□ Channels: periodic in  $x$ , bounded in  $y$  (Alonso et al. 2002)

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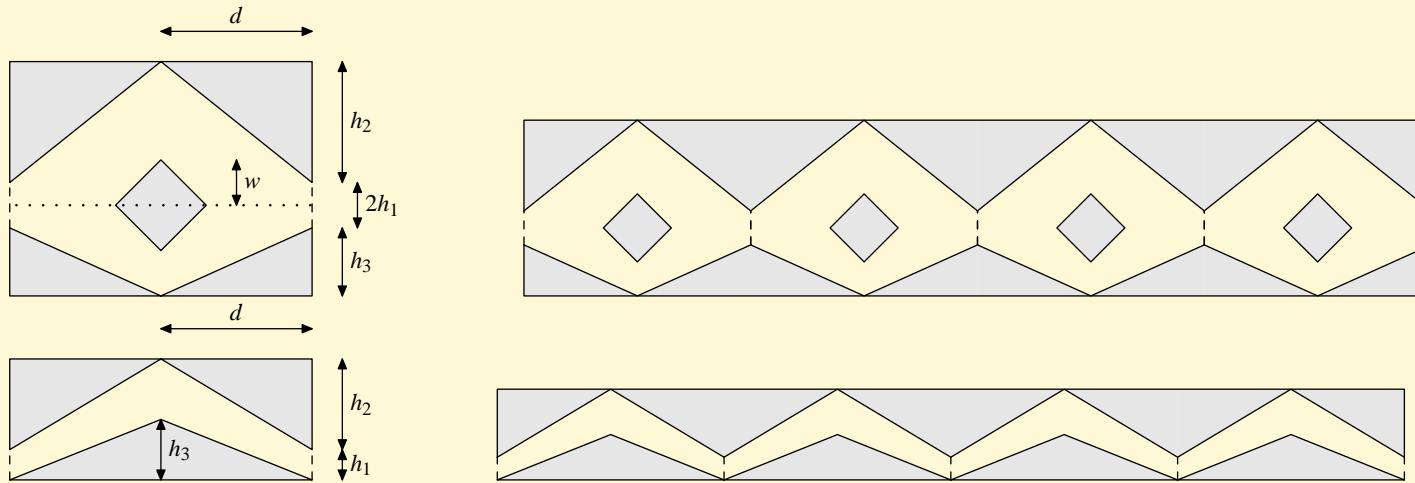
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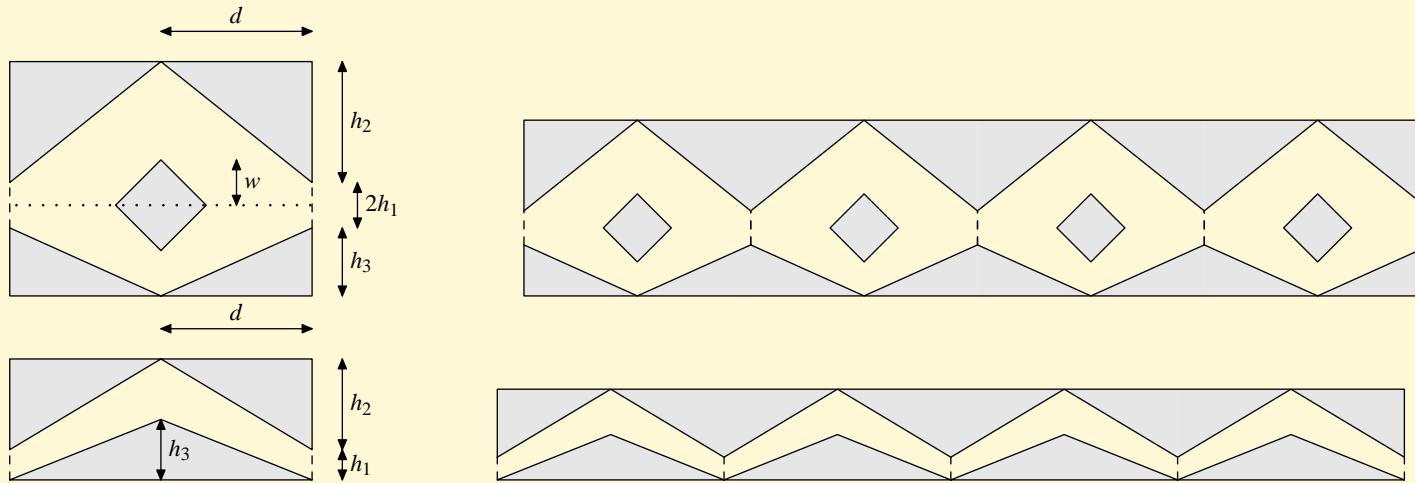
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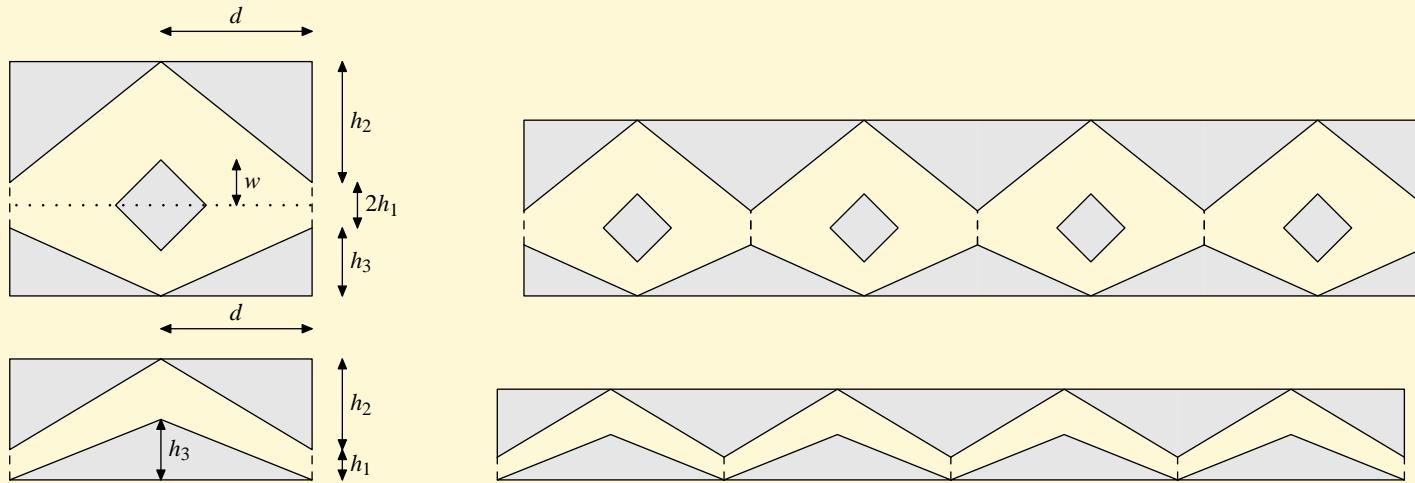
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- Jepps & Rondoni 2006

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# Results: polygonal billiards, irrational angles

- Statistical properties: average  $\langle \cdot \rangle$  over initial conditions

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condition	$\sigma^2(t)$ asymptotic	diffusion
generic	$t$	normal
infinite horizon	$t \log t$	marginal anomalous
parallel scatterers	$t^\alpha, \quad \alpha > 1$	anomalous <b>superdiffusion</b>

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# Effect of horizon

- Finite:  $\sigma^2(t) \sim 2Dt$ ; infinite:  $\sigma^2(t) \sim t \log t$

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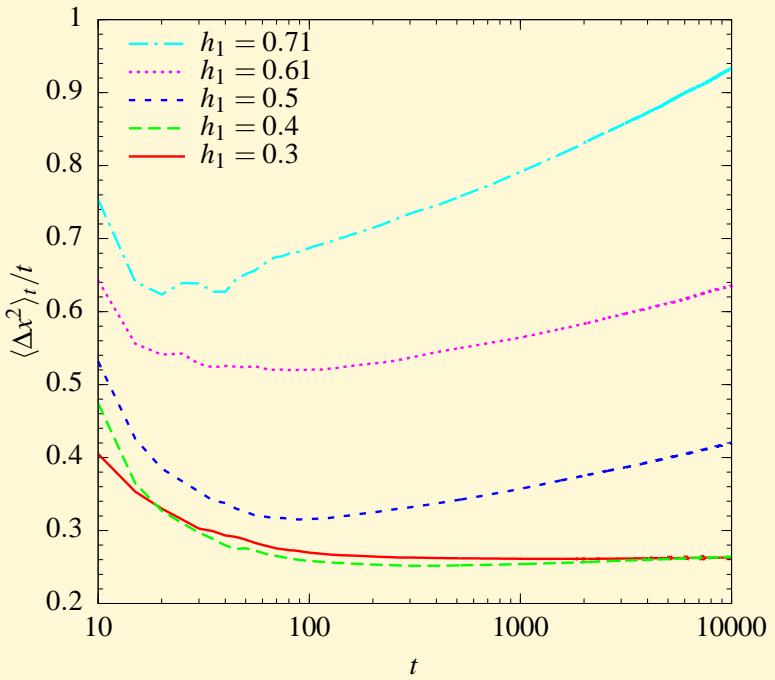
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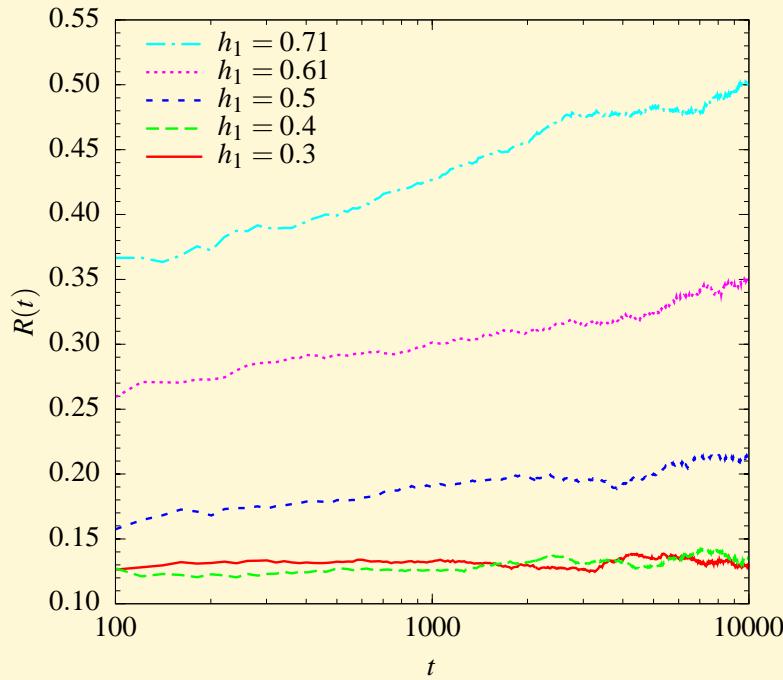
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$$\sigma^2(t)/t$$



$$R(t)$$

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# Fine structure of distributions

## □ Diffusion: ‘spreading out’ of distributions

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# Fine structure of distributions

- Diffusion: ‘spreading out’ of distributions
- Probability density  $\rho_t(x)$  of particle positions

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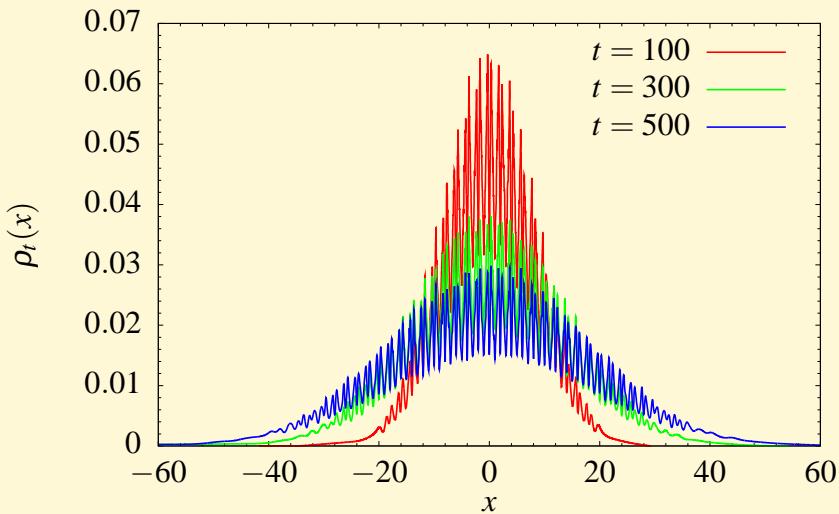
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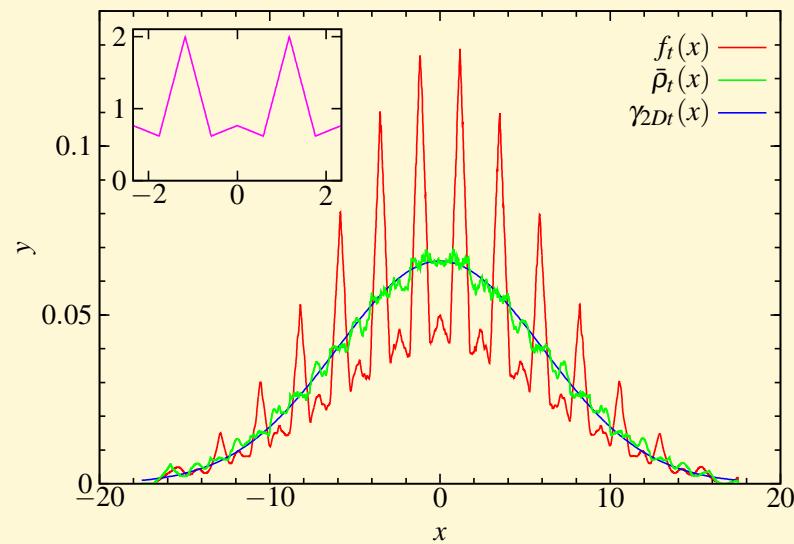
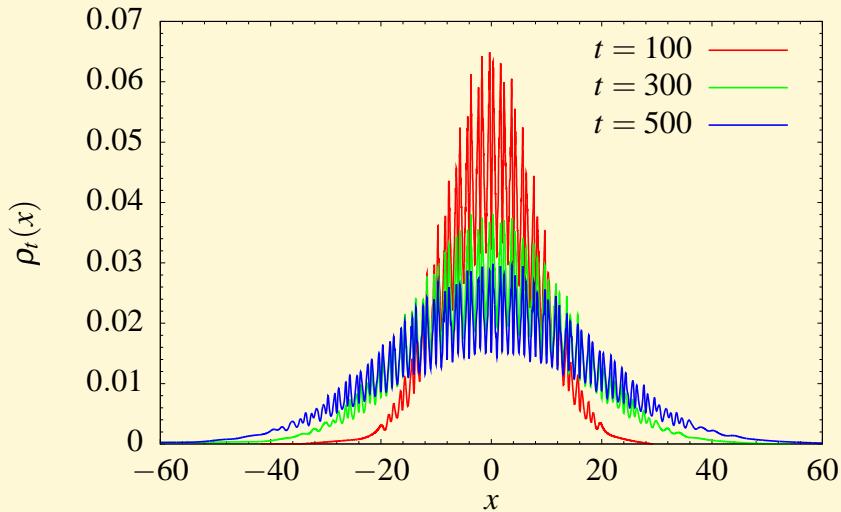
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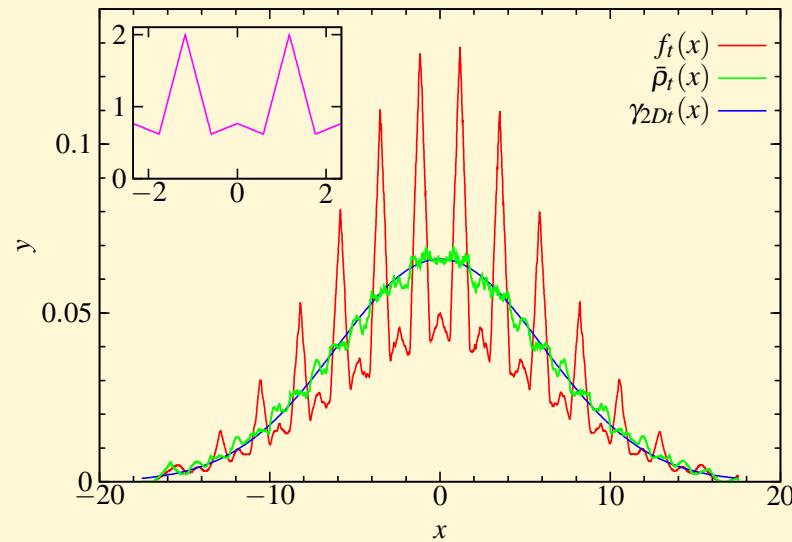
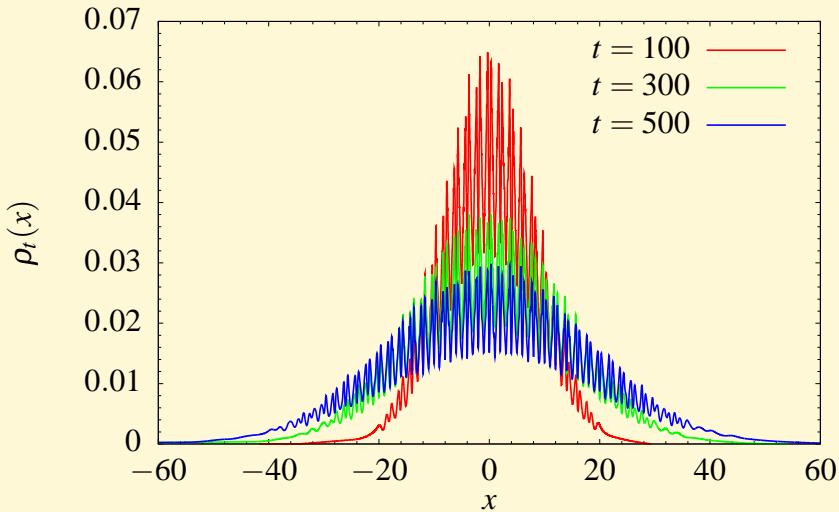
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- Demodulate by available height  $h(x)$ :

$$f_t(x) := \frac{\rho_t(x)}{h(x)}$$

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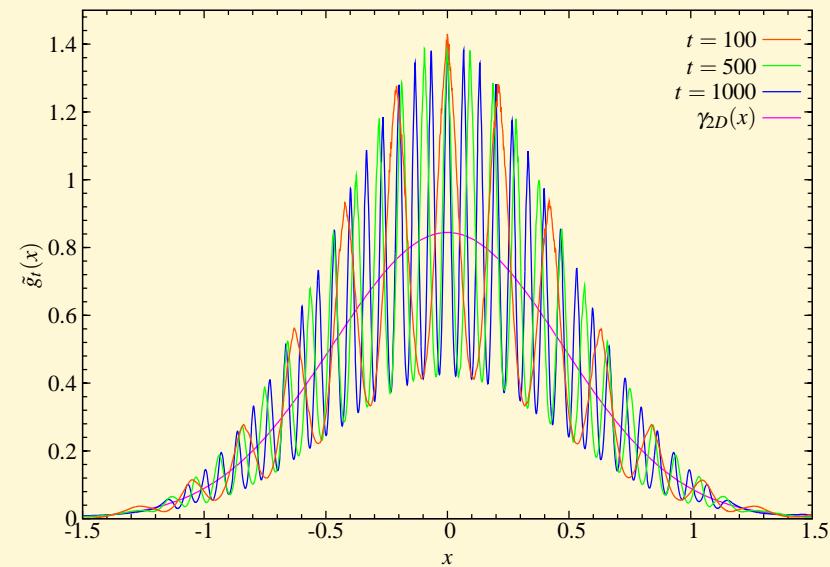
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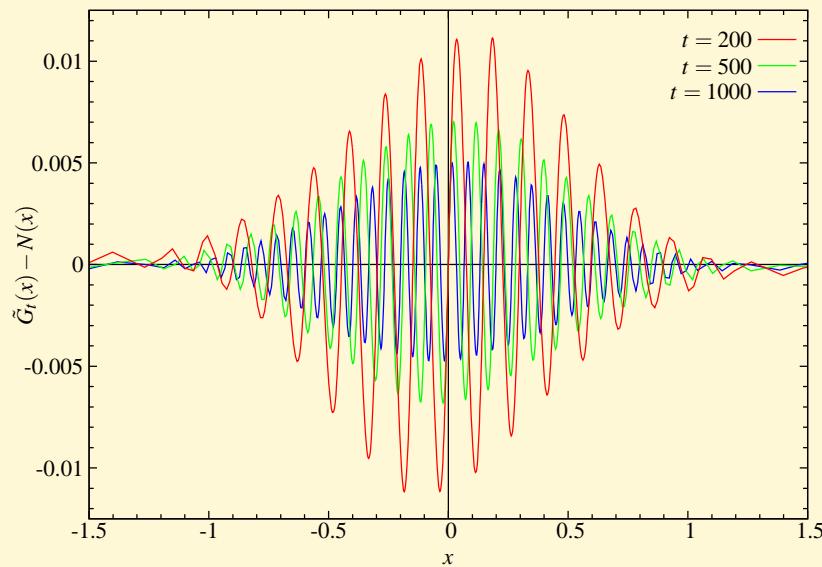
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# Weak convergence (Lorentz gas channel)



rescaled densities  
 $\tilde{f}_t(x) := \sqrt{t} f_t(x\sqrt{t})$



distance of  
cumulative distribution  
from normal

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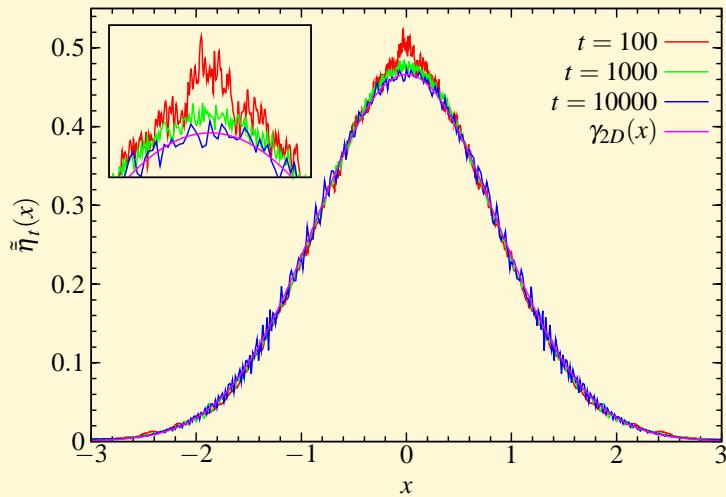
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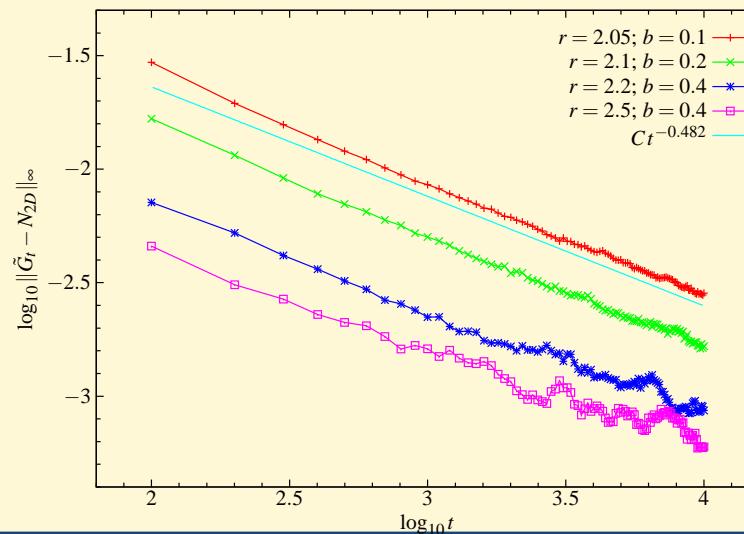
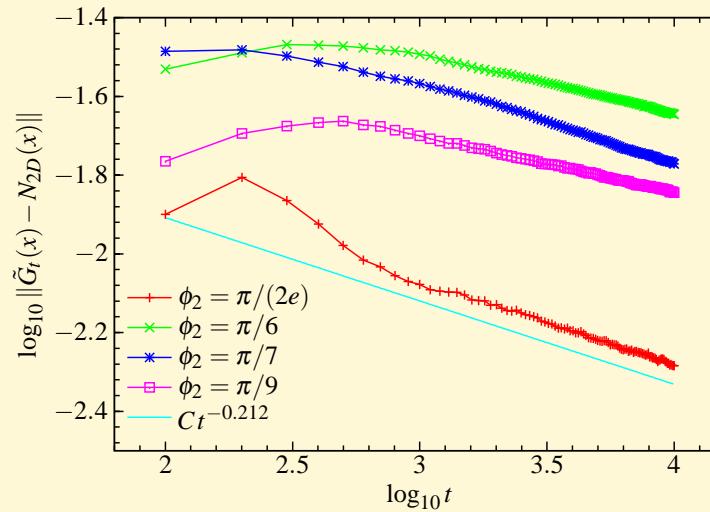
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# Central limit theorem



polygonal: distance  $\sim t^{-0.21}$   
 Lorentz: distance  $\sim t^{-0.48}$

Pène (2002): faster than  $t^{-1/6}$   
 Heuristic: slower than  $t^{-1/2}$



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# Anomalous super-diffusion

- Anomalous diffusion  $\sigma^2(t) \sim t^\alpha$  when parallel scatterers

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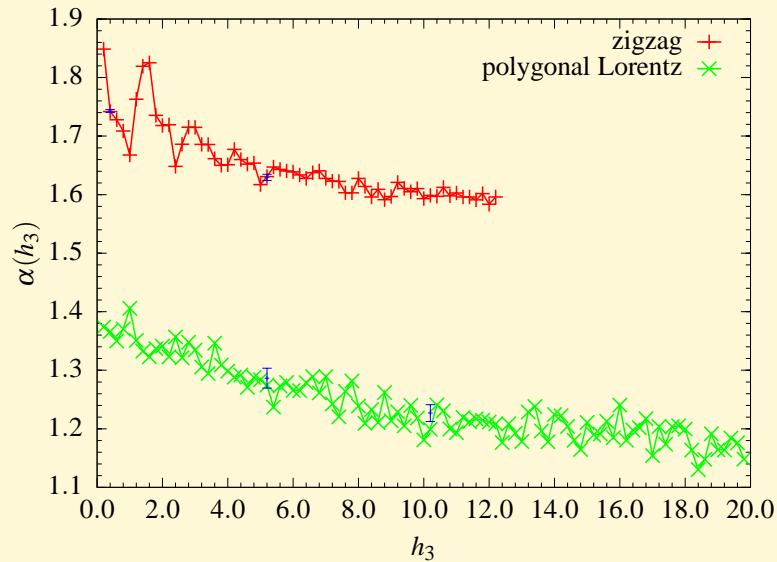
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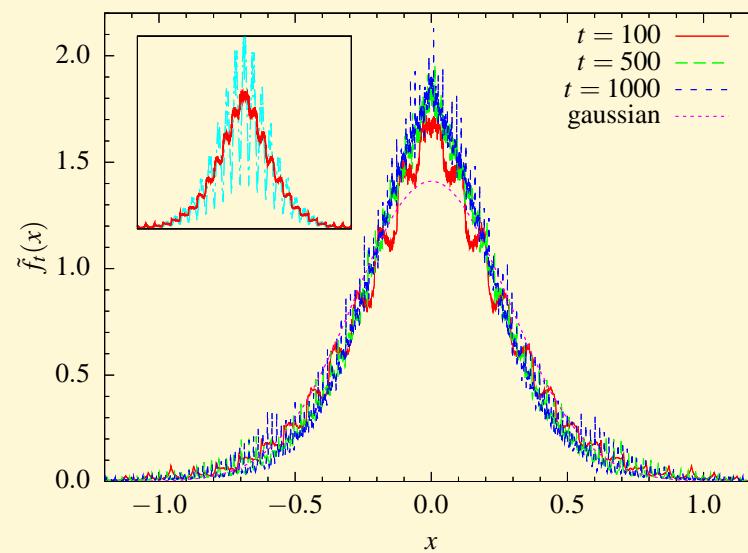
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# Parallel scatterers

□ Anomalous diffusion  $\sigma^2(t) \sim t^\alpha$  when parallel scatterers



exponent  $\alpha$



rescaled densities  $\tilde{f}_t(x) := t^{\alpha/2} f_t(x t^{\alpha/2})$

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## □ What is reason for anomalous diffusion?

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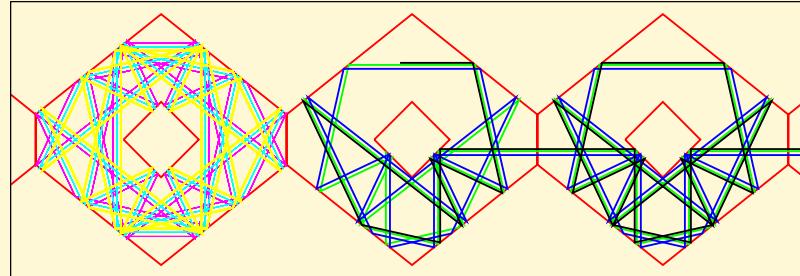
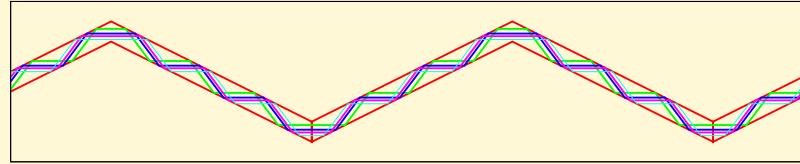
- What is reason for anomalous diffusion?
- Families of **propagating periodic orbits**

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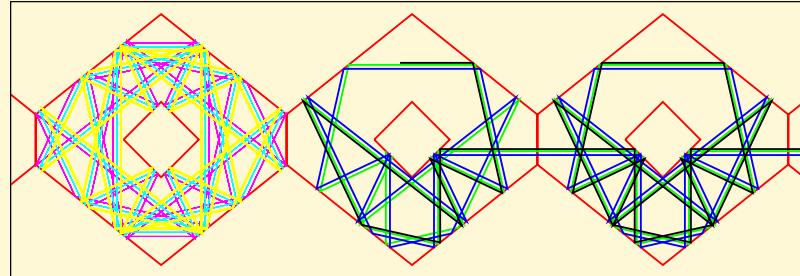
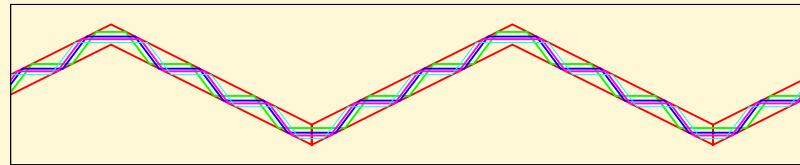
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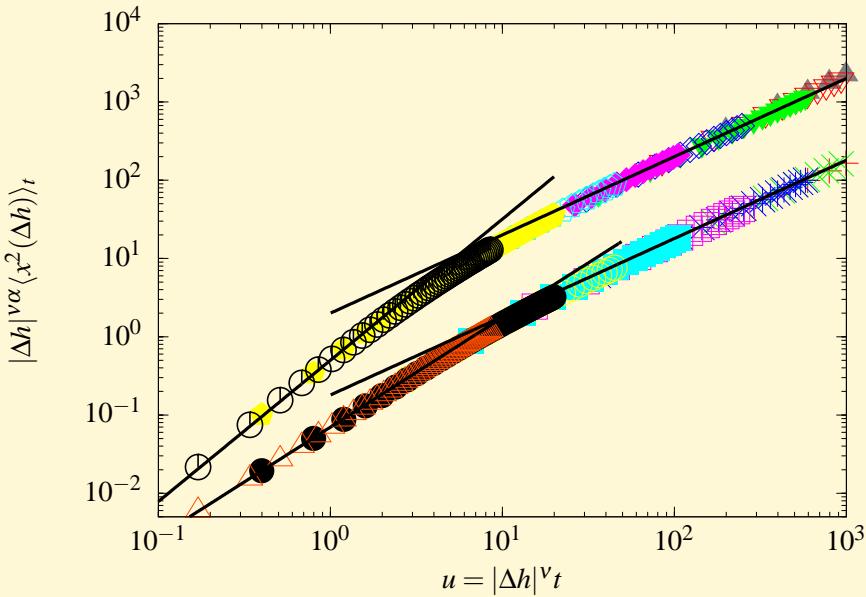
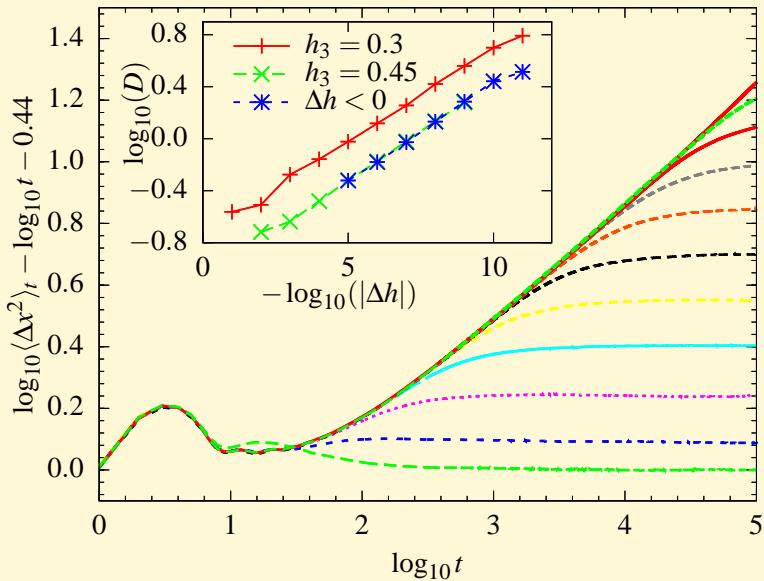
- Much more likely when parallel scatterers
- Model with continuous-time random walks  
(DPS+HL 2006, Schmiedeberg & Stark 2006)

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# Crossover from normal to anomalous



- $\langle x^2(\Delta h) \rangle_t \sim \begin{cases} D(\Delta h)t & \text{for } t > T_c \\ t^\alpha & \text{for } t < T_c \end{cases}; \quad T_c \sim |\Delta h|^{-\nu}$
- Data collapse of  $|\Delta h|^{\nu\alpha} \langle x^2(\Delta h) \rangle_t$  as function of  $u := |\Delta h|^\nu t$ :

$$\phi(u) \sim \begin{cases} u & \text{for } u \gg 1 \\ u^\alpha & \text{for } u \ll 1 \end{cases}$$

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- Conclusions:
  - Numerics give reasonably clear picture

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- Parallel  $\Rightarrow$  propagating orbits  $\Rightarrow$  anomalous diffusion

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