

Spectral Density Calculations for Schrödinger

Operators in 1 Dimension

We consider functions $P = P(x, \lambda)$, $Q = Q(x, \lambda)$,
 $R = R(x, \lambda)$, such that, for each solution
 $y = y(x, \lambda)$ of the differential equation

$$-\frac{d^2y(x, \lambda)}{dx^2} + q(x)y(x, \lambda) = \lambda y(x, \lambda) \quad (x \in [0, \infty), \lambda \in \mathbb{R})$$

the quadratic form

$$Py^2 + Qyy' + Ry'^2 \text{ is independent of } x.$$

We require P, Q, R to satisfy

$$\left. \begin{aligned} \frac{dP}{dx} &= (\lambda - q)Q \\ \frac{dQ}{dx} &= -2P + 2(\lambda - q)R \\ \frac{dR}{dx} &= -Q \end{aligned} \right\} \begin{aligned} B &= \\ &4PR - Q^2 \\ &= \text{const.} \end{aligned}$$

From the linear system for P, Q, R

we have

$$\frac{d^2R}{dx^2} = -\frac{dQ}{dx} = 2P - 2(\lambda - q)R.$$

Hence

$$\frac{d}{dx} \left\{ \frac{d^2R}{dx^2} + 2(\lambda - q)R \right\} = 2 \frac{dP}{dx}$$

$$= 2(\lambda - q)Q = -2(\lambda - q) \frac{dR}{dx}.$$

Hence

$$[R'' + 2(\lambda - q)R]' + 2(\lambda - q)R' \\ = 0$$

Note that if y_1, y_2 both satisfy the Schrödinger equation (with same λ value) then a solution for (P, Q, R) is

$$(P, Q, R)$$

$$= (y_1' y_2', -(y_1 y_2' + y_2 y_1'), y_1 y_2)$$

In particular, defining 2 solutions

$u(x, \lambda), v(x, \lambda)$ by

$$\left. \begin{array}{ll} u(0, \lambda) = 1 & v(0, \lambda) = 0 \\ u'(0, \lambda) = 0 & v'(0, \lambda) = 1 \end{array} \right\}$$

we find the general solution for $R(x, \lambda)$

$$\text{is } R = a u^2 + b u v + c v^2.$$

(a, b, c constants).

Moreover, β

$$\begin{aligned} &= 2R \frac{d^2R}{dx^2} - \left(\frac{dR}{dx} \right)^2 + 4(\lambda - q) R^2 \\ &\quad = 4ac - b^2. \end{aligned}$$

Now consider a solution

$$R = au^2 + buv + cv^2 \text{ for } R(x, \lambda),$$

such that (i) $a > 0$

$$(ii) \beta = 4ac - b^2 = 4.$$

Define A, B by

$$\frac{1}{B} = a, \quad \frac{2A}{B} = b, \quad \frac{A^2 + B^2}{B} = c$$

$$\text{Then } R = \frac{(u + Av)^2 + B^2 v^2}{B}$$

$$= \frac{|u + Mv|^2}{\operatorname{Im} M}, \text{ where}$$

$$M = M(\lambda) = A + iB.$$

IF $M(\lambda) = m_+(\lambda)$ = boundary
value of m -function, then

$$\begin{aligned} \text{Spectral density} &= \frac{1}{\pi} \operatorname{Im} M(\lambda) = \frac{1}{\pi} \overline{R(0, \lambda)} \\ &= \frac{1}{\pi(Pv^2 + Qvv' + Rv'^2)} \end{aligned}$$

Any solution $R(x, \lambda)$ for R , satisfying
 $\beta = 4$, will give rise to an expression

for spectral density (Dirichlet operator)

$$= \frac{1}{\pi (Pv^2 + Qvv' + Rv'^2)} \quad \text{provided either}$$

- (i) $R(x, \lambda) \rightarrow R(\lambda)$ as $x \rightarrow \infty$, or,
more generally,
- (ii) $R(x, \lambda) \sim F(x, \lambda)$ as $x \rightarrow \infty$, with
 $\frac{\partial F(x, \lambda)}{\partial x} \rightarrow 0$, or (less generally)
- (iii) $Q(x, \lambda) \rightarrow 0$ as $x \rightarrow \infty$, or
- (iv) $R(x, \lambda)$ is periodic in x .

Such conditions hold for

- a) potentials $(1+x)^{\beta}$ for any $\beta > 0$.
- b) potential $q(x) = -x^2$
- c) potential $q(x) = \cos x$

Moreover, the same functions P, Q, R
allow one to determine the spectral density
for the Schrödinger operator $-\frac{d^2}{dx^2} + q(x)$
subject to a general boundary condition
 $(\cos \alpha) f(0) + \sin \alpha f'(0) = 0$.

If, in that case, we define solution v_α of
 $-y'' + q y = \lambda y$, with

$$v_\alpha(0, \lambda) = -\sin \alpha \quad \left. \right\}$$

$$v_\alpha'(0, \lambda) = \cos \alpha \quad \left. \right\}$$

then the spectral density is

$$\frac{1}{\pi(P v_\alpha^2 + Q v_\alpha v_\alpha' + R v_\alpha'^2)}$$

Ideas of value distribution may be used to describe the asymptotics of solutions $y(x, \lambda)$ of the Schrödinger equation in the limit as $x \rightarrow \infty$.

One considers the distribution of values of $y(x, \cdot)$ as a function of λ , for large x .

Starting from the identities that

$$u = R^{1/2} \left\{ B^{1/2} \cos \int_0^x \frac{1}{R(t)} dt - A B^{-1/2} \sin \int_0^x \frac{1}{R(t)} dt \right\}$$

$$v = R^{1/2} B^{-1/2} \sin \int_0^x \frac{1}{R(t)} dt,$$

one may show that, asymptotically for large x , $\theta(x, \lambda) = \int_0^x \frac{1}{R(t)} dt$ is uniformly distributed mod π ,

as a function of λ . For the matrix of u_α, v_α

and their derivatives, we have

$$\begin{pmatrix} u_\alpha & v_\alpha \\ u_\alpha' & v_\alpha' \end{pmatrix} = \begin{pmatrix} F_\alpha (\bar{B}_\alpha \cos \theta_\alpha - \bar{A}_\alpha \sin \theta_\alpha) & F_\alpha \sin \theta_\alpha \\ \bar{B}_\alpha G_\alpha \cos(\theta_\alpha + \beta) - \bar{A}_\alpha G_\alpha \sin(\theta_\alpha + \beta) & G_\alpha \sin(\theta_\alpha + \beta) \end{pmatrix}$$

General solution for R is given in terms of particular solution R_0 (with $\beta = 4$) by

$$R = R_0 \left(A + B \cos \int_0^x \frac{2}{R_0(t)} dt + C \sin \int_0^x \frac{2}{R_0(t)} dt \right)$$

where to maintain $\beta = 4$ we require

$$A^2 - B^2 - C^2 = 1.$$

If $R_0 = a_0 u^2 + b_0 uv + c_0 v^2$, we have

$$\int_0^x \frac{1}{R_0(t)} dt = \cot^{-1} \left(\frac{u + A_0 v}{B_0 v} \right) + \pi N(x, \lambda)$$

$(N(x, \lambda) = \text{no. of zeroes of } v(x, \lambda)$
in interval $(0, x)$).

As $x \rightarrow \infty$, $\int_0^x \frac{2}{R_0(t)} dt$, if $M(\lambda) = m_+(\lambda)$, may be regarded as uniformly distributed (modulo 2π) as a function of λ .

Any other solution for $R(x, \lambda)$ has
 $\text{ave} \left\{ \frac{R(x, \lambda)}{R_0(x, \lambda)} \right\} > 1$ in limit as $x \rightarrow \infty$.

Example

Suppose $q' \in L^1(0, \infty)$ with $q \rightarrow 0$ as $x \rightarrow \infty$

Define P_0, Q_0, R_0 by

$$\left. \begin{aligned} P(x, \lambda) &= (\lambda - q(x))^\alpha P_0(x, \lambda) \\ Q &= (\lambda - q)^\beta Q_0 \\ R &= (\lambda - q)^\gamma R_0 \end{aligned} \right\} \begin{aligned} \alpha &= \beta - \gamma \\ &= \frac{1}{2} \end{aligned}$$

Change variable from x to

$$s = \int^x (\lambda - q(t))^{1/2} dt$$

Then

$$\frac{d}{ds} \begin{pmatrix} P_0 \\ Q_0 \\ R_0 \end{pmatrix} = A \begin{pmatrix} P_0 \\ Q_0 \\ R_0 \end{pmatrix} + \frac{q'}{(\lambda - q)^{3/2}} \begin{pmatrix} \alpha P_0 \\ \beta Q_0 \\ \gamma R_0 \end{pmatrix}$$

$$\text{Here } A = \begin{pmatrix} 0 & 1 & 0 \\ -2 & 0 & 2 \\ 0 & -1 & 0 \end{pmatrix} \quad \begin{array}{l} \text{(eigenvalues:} \\ 0, \pm 2i \end{array}$$

Converts to integral equation

$$\underline{P}_0(s) = \left\{ \begin{array}{l} \text{eigenvector with} \\ \text{eigenvalue 0} \end{array} \right\} - \int_s^\infty e^{A(s-t)} \Phi(t) \underline{P}_0(t) dt$$

$$\text{Take eigenvector } = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \alpha = \frac{1}{2}, \beta = 0, \gamma = -\frac{1}{2}$$