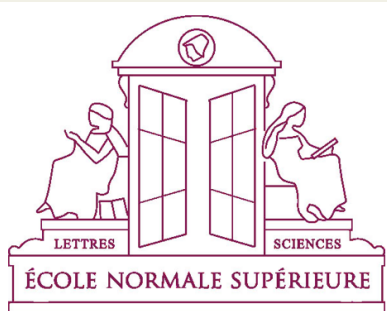


Parameter Estimation and New Application Areas



Michael Ghil

Ecole Normale Supérieure, Paris, and
University of California, Los Angeles



Joint work (recently) with

M. D. Chekroun, D. Kondrashov & Y. Shprits, UCLA; A. Carrassi, IRM, Brussels; L. Roques and S. Soubeyrand, INRA, Avignon; C.-J. Sun, CSIRO, Perth; A. Trevisan, ISAC-CNR, Bologna; and many others: please see

<http://www.atmos.ucla.edu/tcd/> and <http://www.environnement.ens.fr/>

Outline

- Data in meteorology, oceanography and space physics
 - *in situ* & remotely sensed
- Basic ideas, data types, & issues
 - how to combine data with models
 - transfer of information
 - between variables & regions
 - filters & smoothers
 - stability of the forecast-assimilation cycle
- Parameter estimation
 - model parameters
 - noise parameters – at & below grid scale
- Novel areas of application
 - space physics
 - shock waves in solids
 - macroeconomics
 - paleoclimate
- Concluding remarks and bibliography

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Parameter Estimation

a) Dynamical model

$$dx/dt = M(x, \mu) + \eta(t)$$

$$y^o = H(x) + \varepsilon(t)$$

Simple (EKF) idea – augmented state vector

$$d\mu/dt = 0, X = (x^T, \mu^T)^T$$

b) Statistical model

$$L(\rho)\eta = w(t), \quad L - \text{AR(MA) model, } \rho = (\rho_1, \rho_2, \dots, \rho_M)$$

Examples: 1) Dee *et al.* (*IEEE*, 1985) – estimate a few parameters in the covariance matrix $Q = E(\eta, \eta^T)$; also the bias $\langle \eta \rangle = E\eta$;

2) POPs - Hasselmann (1982, *Tellus*); Penland (1989, *MWR*; 1996, *Physica D*); Penland & Ghil (1993, *MWR*)

3) $dx/dt = M(x, \mu) + \eta$: Estimate both M & Q from data (Dee, 1995, *QJ*), Nonlinear approach: Empirical mode reduction (EMR: Kravtsov *et al.*, *J. Clim.*, 2005; Kondrashov *et al.*, *J. Clim.*, 2005; Strounine *et al.*, *Physica D*, 2009)

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Estimating noise – I

$$Q_1 = Q_{slow}, \quad Q_2 = Q_{fast}, \quad Q_3 = 0;$$

$$R_1 = 0, \quad R_2 = 0, \quad R_3 = R;$$

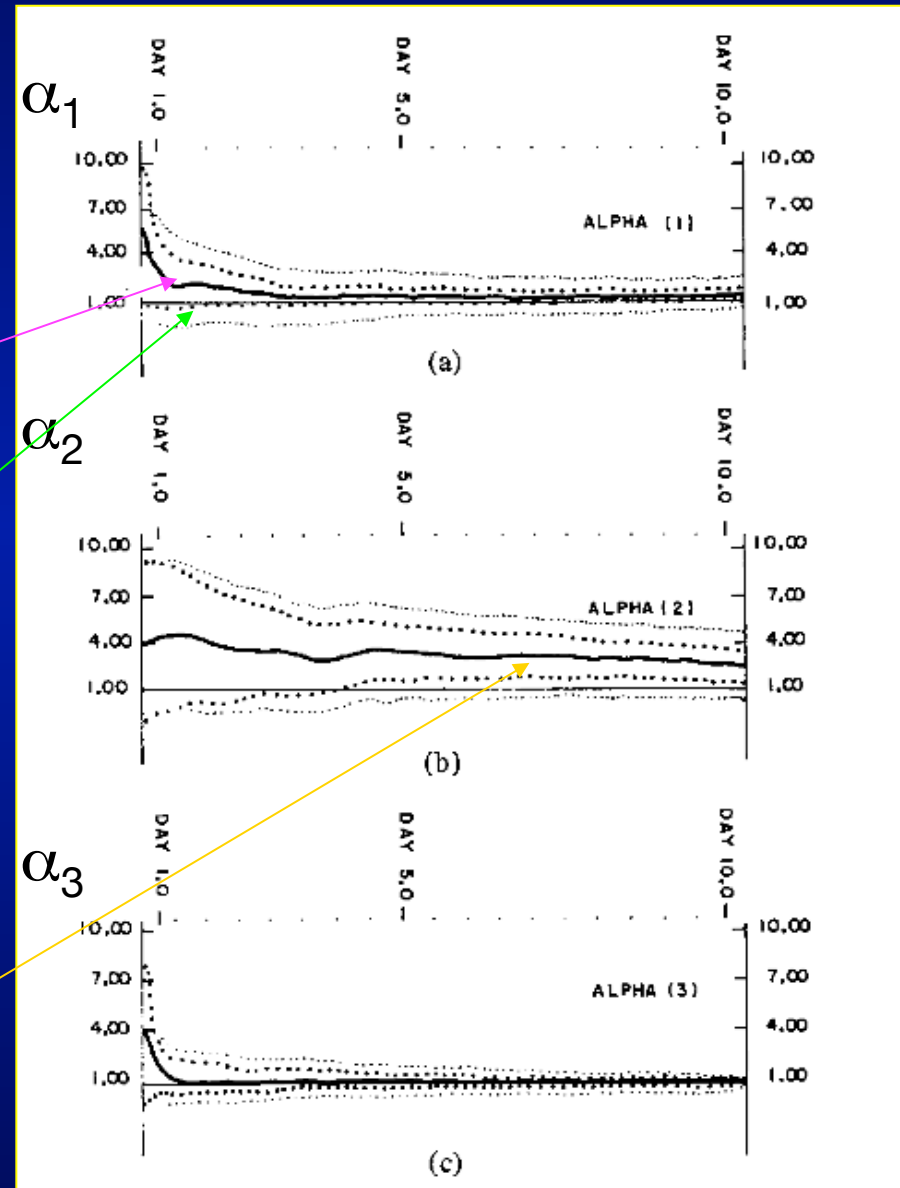
$$Q = \sum \alpha_i Q_i, \quad R = \sum \alpha_i R_i;$$

$$\alpha(0) = (6.0, 4.0, 4.5)^T;$$

$$Q(0) = 25 * I.$$

Dee et al. (1985, *IEEE Trans. Autom. Control*, AC-30)

Poor convergence for Q_{fast} ?



Estimating noise – II

Same choice of $\alpha(0)$, Q_i ,
and R_i but

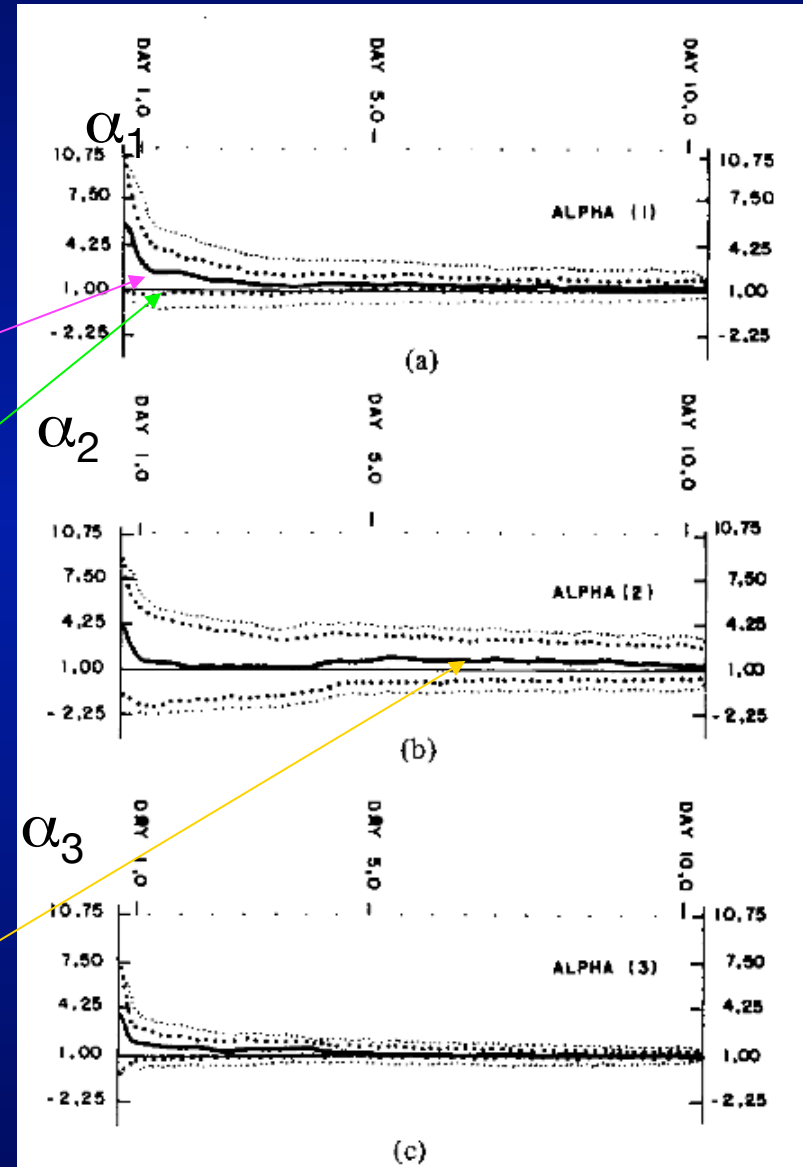
$$\Theta(0) = 25 * \begin{bmatrix} 1 & 0.8 & 0 \\ 0.8 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

estimated

true ($\alpha = 1$)

Dee et al. (1985, *IEEE Trans. Autom. Control*, **AC-30**)

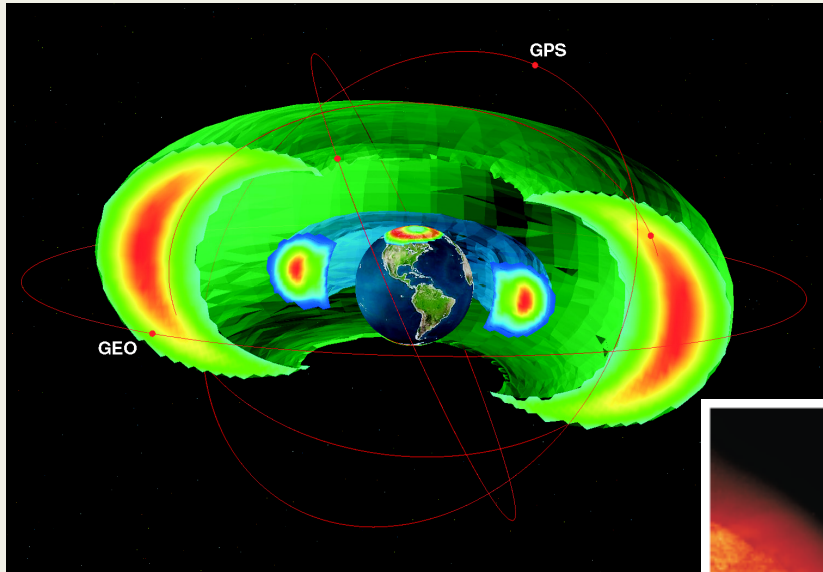
Good convergence for Q_{fast} !



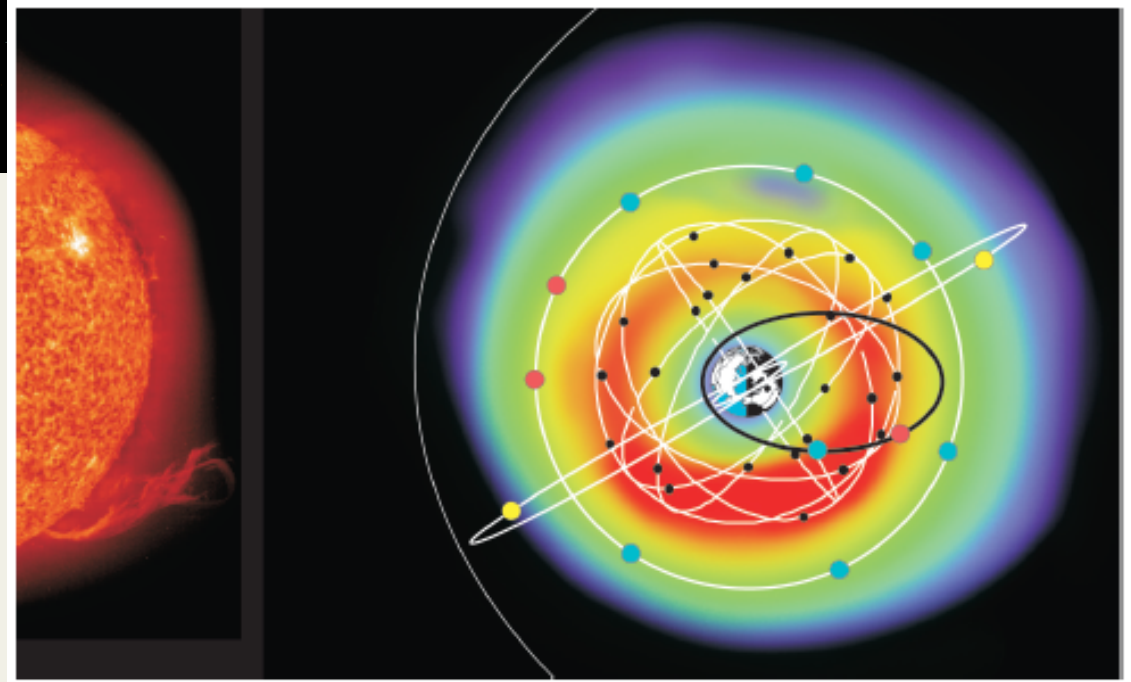
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Space physics data



Two decades ago ...



... and now

Space platforms in Earth's magnetosphere

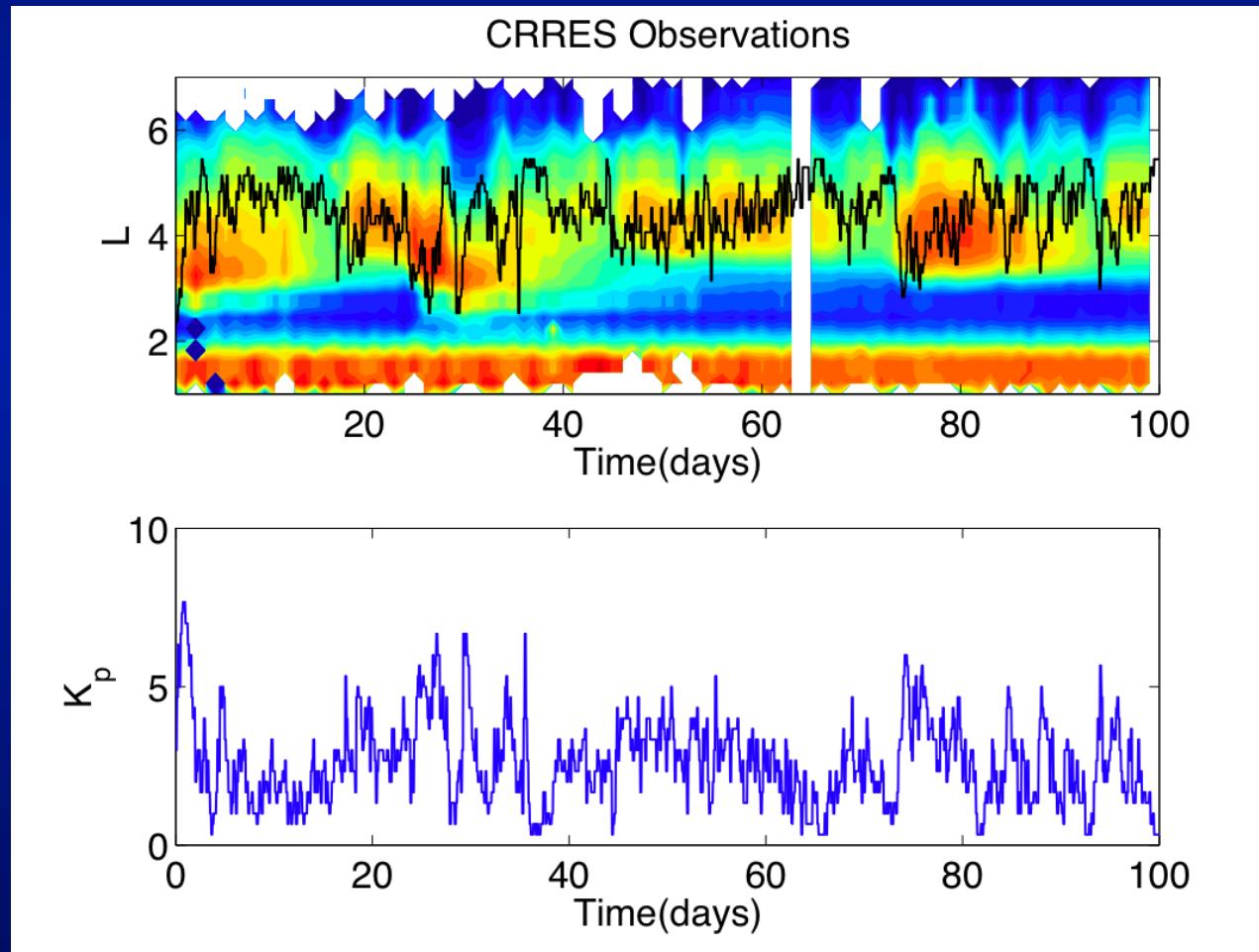
Parameter Estimation for Space Physics – I

Daily fluxes of 1 MeV relativistic electrons in Earth's outer radiation belt
(CRRES observations from 28 August 1990)

K_p - index of solar activity (external forcing) – used to determine the position

of the plasmapause L_{pp}

(black) in the observations



Kondrashov, Shprits,
Ghil & Thorne
(*J. Geophys. Res.*, 2007)

Parameter estimation for space physics – II

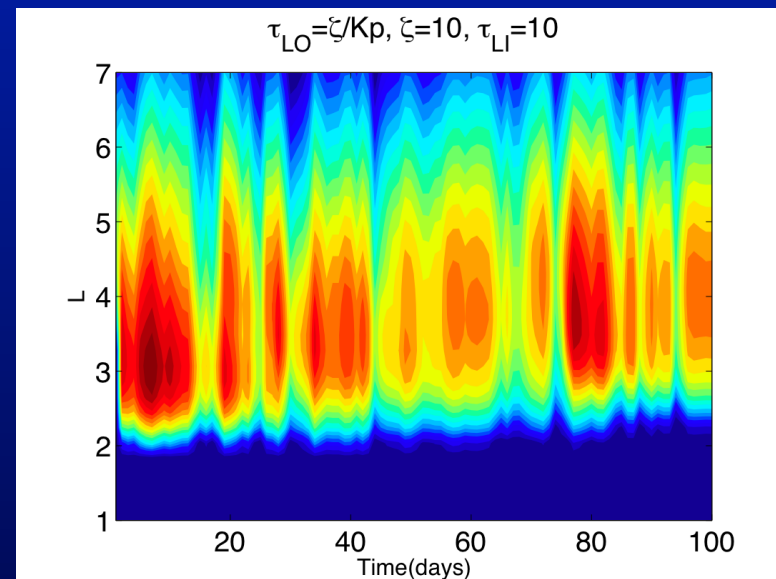
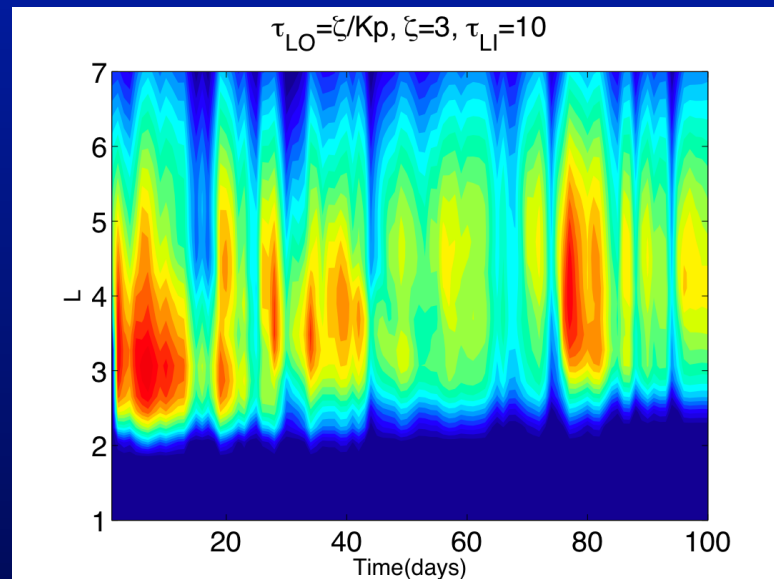
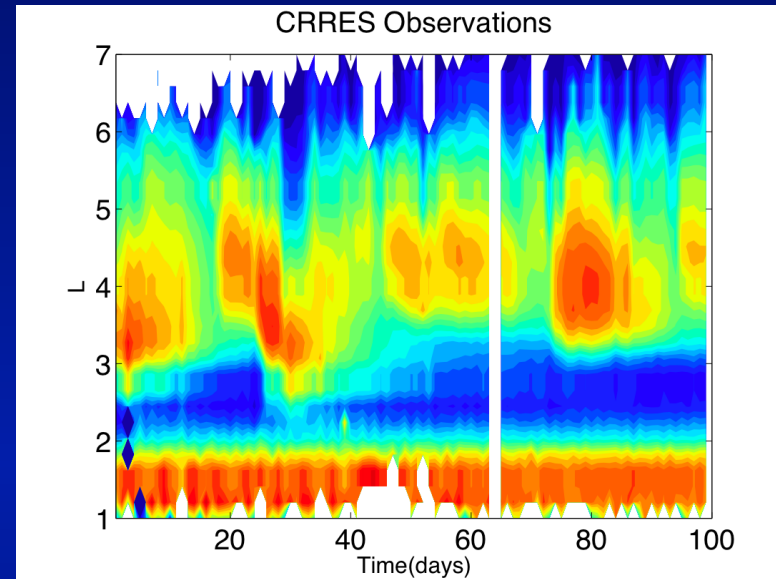
HERRB-1D code (Y. Shprits) –
estimating phase-space density
 f and electron lifetime τ_L :

$$\frac{\partial f}{\partial t} = L^2 \frac{\partial}{\partial L} \left(L^{-2} D_{LL} \frac{\partial f}{\partial L} \right) - \frac{f}{\tau_L}$$

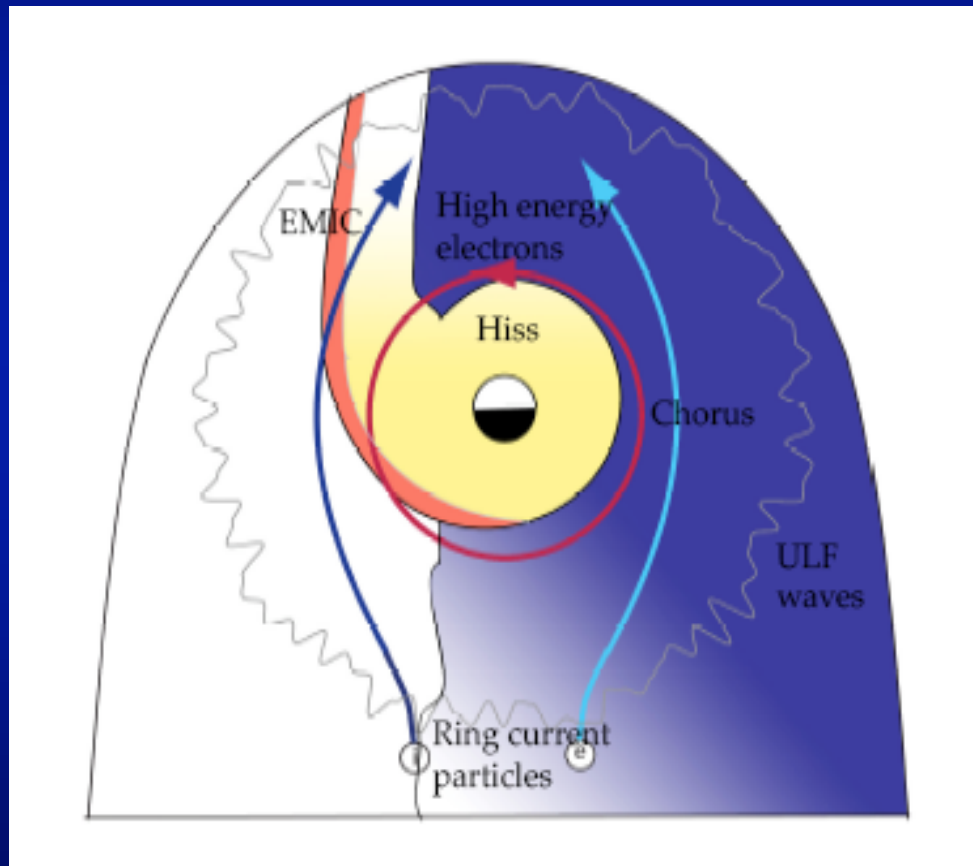
Different lifetime parameterizations for
plasmasphere – out/in:

$$\tau_{Lo} = \zeta / K_p(t); \tau_{Li} = \text{const.}$$

What are the **optimal** lifetimes to match
the observations best?



Dominant loss mechanisms



Pitch angle scattering due to resonance interactions with :

- 1) Plasmaspheric hiss (whistler mode waves) loss time on the scale of 5-10 days (Lyons & Thorne, 1973; Abel & Thorne, 1998; Meredith et al., 2006)
- 2) Chorus waves outside plasmopause provide fast losses on the scale of a day (Horne et al., 2005; Albert et al., 2005; O'Brien, 2004; Thorne et al., 2005)
- 3) EMIC waves mostly in plumes on the dusk side – very fast localized losses (Millan et al., 2002; Summers & Thorne, 2003; Albert, 2003, Bortnik et al., 2006; Shprits et al., 2006a)
- 4) Combined effect of losses to magnetopause and outward radial diffusion (Shprits et al., 2006b).

Parameter estimation for space physics – III

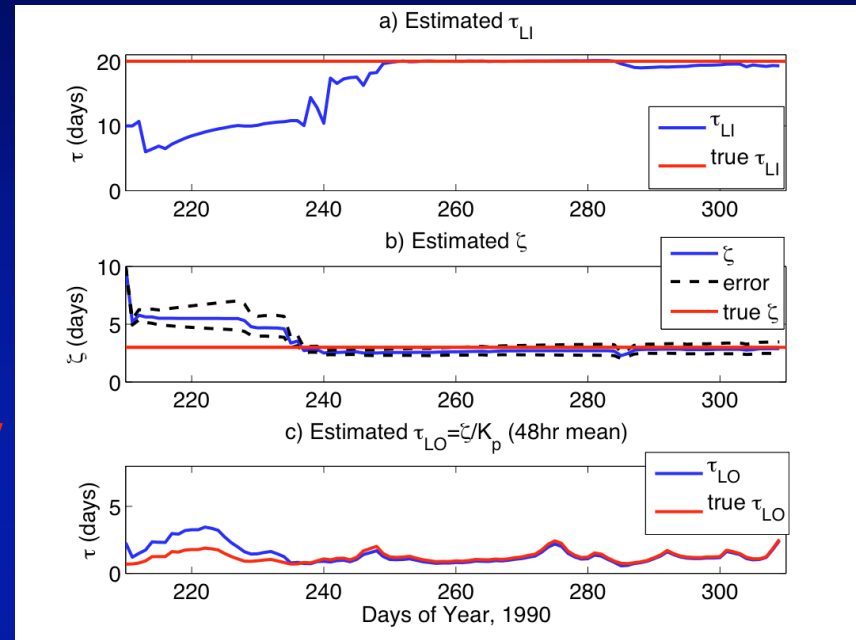
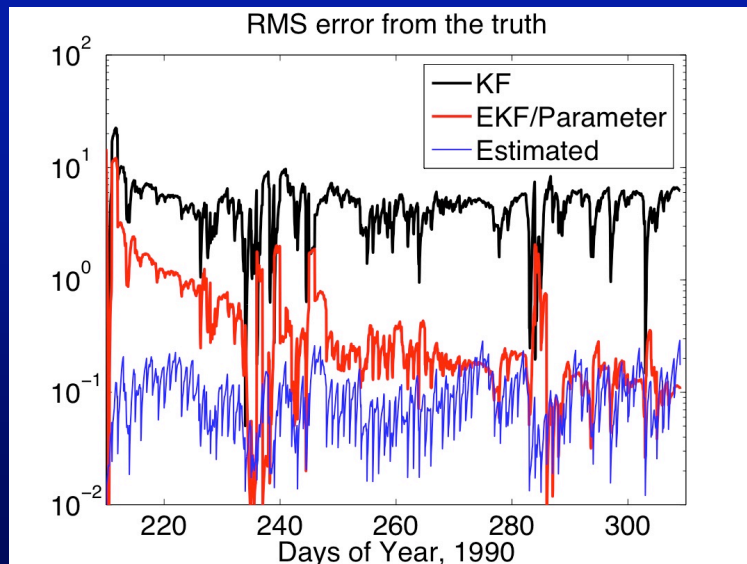
Daily observations from the “truth” —

$$\tau_{Lo} = \zeta/K_p, \quad \zeta = 3, \quad \text{and} \quad \tau_{Li} = 20$$

are used to correct the model’s “wrong” parameters, $\zeta = 10$ and $\tau_{Li} = 10$.

The estimated error $\text{tr}(P^f) \approx$ actual.

When the parameters’ assumed uncertainty is large enough, their EKF estimates converge rapidly to the “truth”.



Black – actual errors for state estimation only

Red – actual errors for state and parameter estimation

Blue – EKF-estimated error ($\text{tr} P_k^f$)

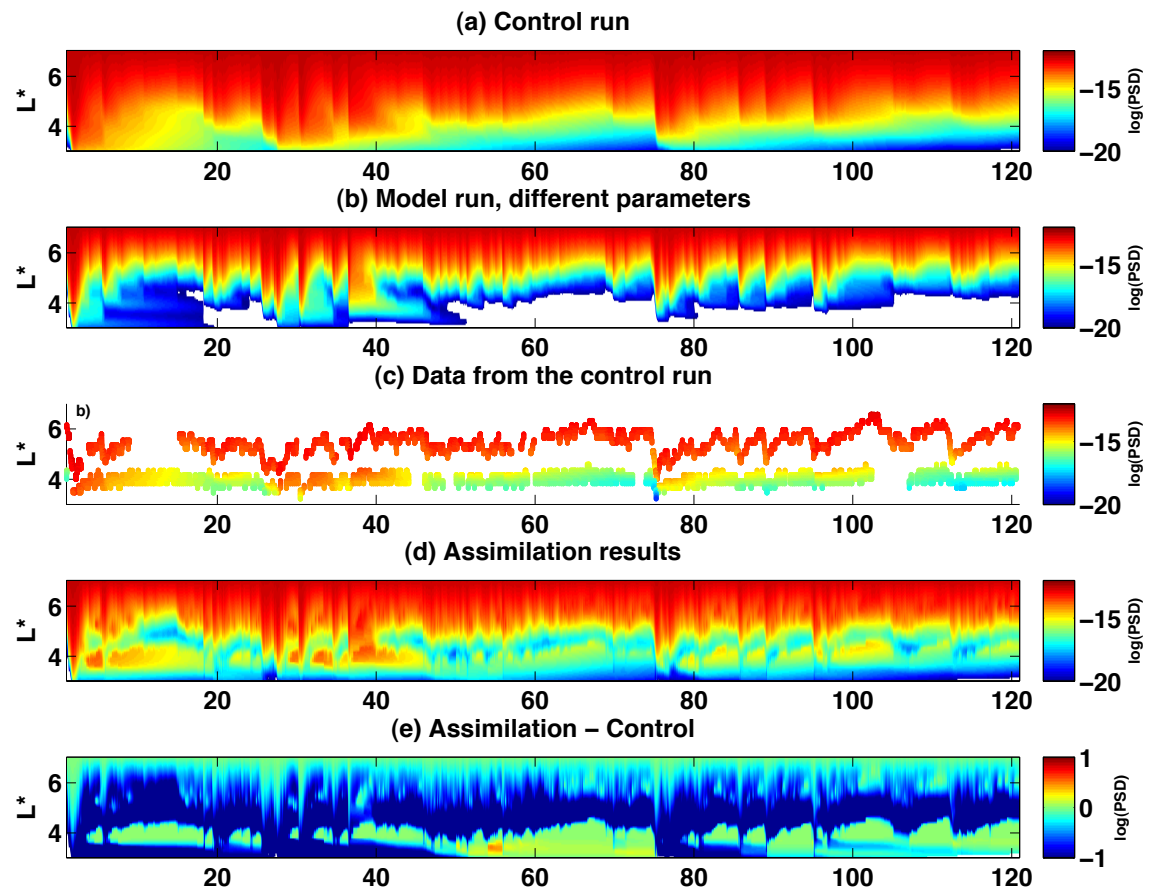
Log-normal EKF for Order-of-Magnitude Changes in Dependent Variables: Space Plasmas – I

Phase-space densities (PSDs) in the Van Allen radiation belts vary by several orders of magnitude over the interval $1 \leq L \leq 6R_E$, where $R_E = \text{Earth's radius}$. This interval includes sharp gradients at the time-varying plasmopause:

$$L_{PP} = 2R_E - 6R_E.$$

Not good for standard sequential (or control) methods that assume normally distributed errors → **Change of variables!**

D. Kondrashov, Y. Shprits & M. Ghil (*Space Weather*, 2011, submitted)



Log-normal EKF for Order-of-Magnitude Changes in Dependent Variables: Space Plasmas – II

Introduce the new variable $S = \log(f)$ to yield the nonlinear PDE in S :

$$\frac{\partial S}{\partial t} = L^2 \frac{\partial}{\partial L} \left(\frac{1}{L^2} D_{LL} \frac{\partial S}{\partial L} \right) - \frac{1}{\tau_L} + D_{LL} \left(\frac{\partial S}{\partial L} \right)^2 .$$

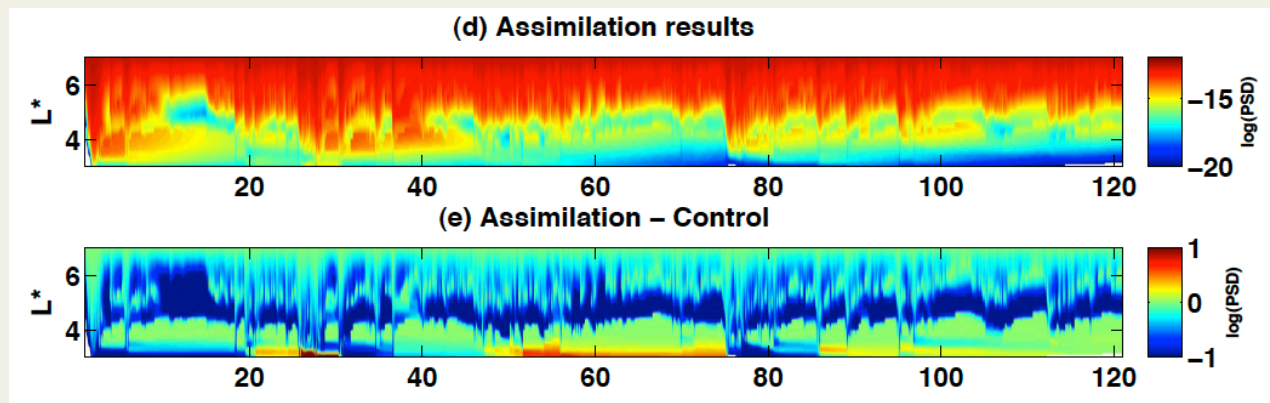
To deal with the nonlinearity and the sharp gradients, we use a total-variation diminishing, second-order scheme (A. Harten, *JCP*, 1983).

The linear Kalman filter for the original PDE in f has to be replaced by an EKF. Results are definitely better with the modified PDE & the log-EKF, as shown by the plot below for “fraternal (dizygotous)-twin” experiments. This is especially so when the observational error covariances R are much larger than the model errors Q .

Another way of evaluating assimilation scheme performance is by considering the variance of the innovation sequence residuals:

$$E \mathbf{z}_k^T \mathbf{z}_k, \text{ where } \mathbf{z}_k \equiv \mathbf{y}_k^o - \mathbf{H} \mathbf{x}_k^f .$$

D. Kondrashov, Y. Shprits & M. Ghil (*Space Weather*, 2011, submitted)



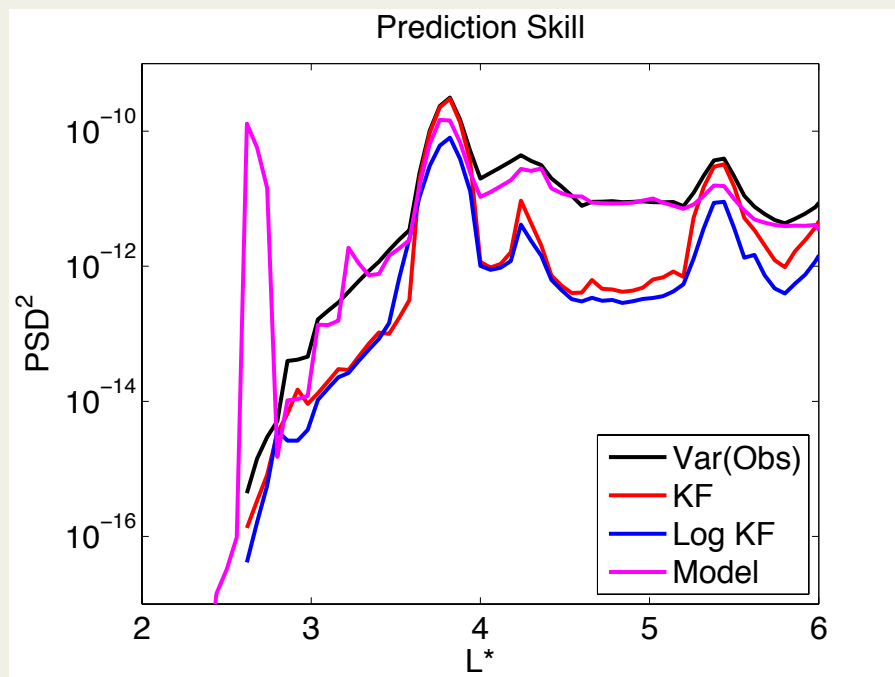
Log-normal EKF for Order-of-Magnitude Changes in Dependent Variables: Space Plasmas – III

We have used real observational data sets from 4 spacecraft missions: the Combined Release and Radiation Effects Satellite (CRRES), GEO-1989 (GEO), GPS NS18 (GPS), and Akebono.

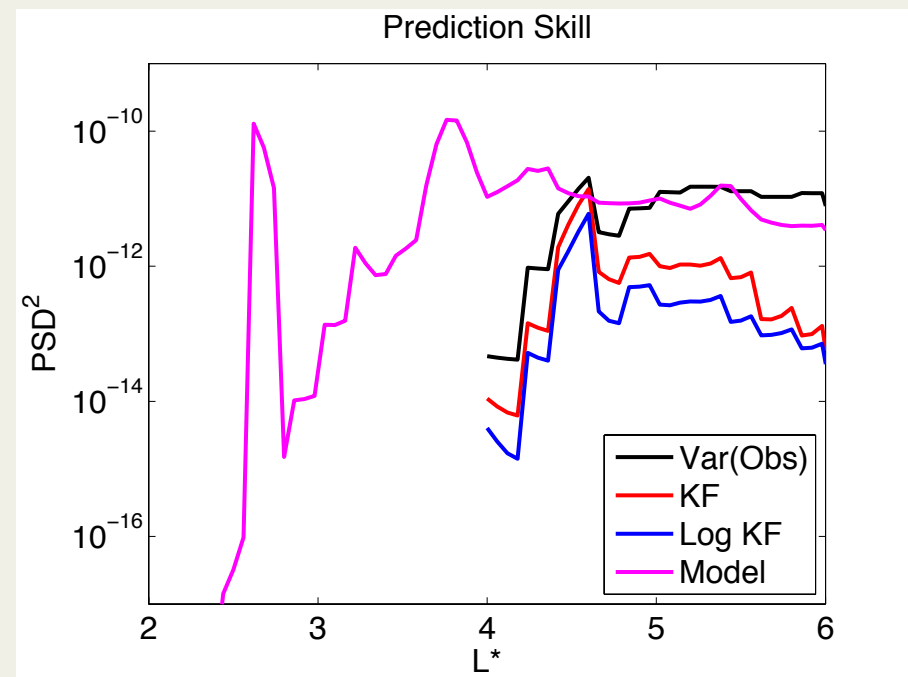
CRESS has the best coverage and accuracy, and was used as a benchmark.

Assimilation was performed with the Akebono and GEO observations, separately.

Plotted are the results for $Ez_k^T z_k$, $z_k \equiv y_k^o - Hx_k^f$.



Akebono



GEO

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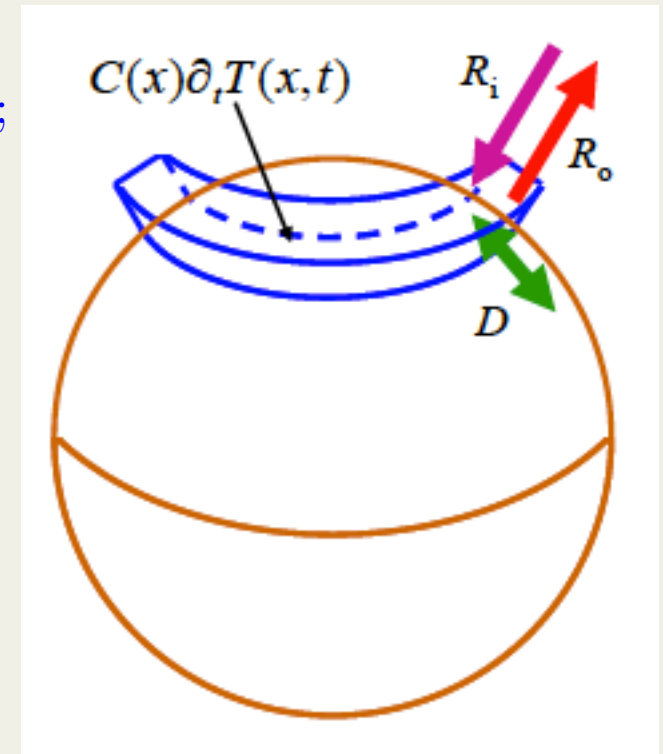
Parameter estimation for energy balance models with memory (EBMMs) – I

One considers a 1-D paleoclimate model governed by an EBM for zonally averaged surface air temperatures $T(t, x)$:

$$c(x) \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left(k(x) \frac{\partial T}{\partial x} \right) + \mu Q(x)[1 - a(x, T)] - g(x, T);$$

here $R_i = \mu Q(x)[1 - a(x, T)]$ is the absorbed solar radiation, with $a = a(x, T)$ the planetary albedo, and $R_o = g(x, T)$ is the terrestrial radiation, modified by the greenhouse effect, while $0 \leq x \leq 1$ is a meridional variable. The albedo depends on past temperatures, because of the long time needed to build up and melt ice sheets.

Ghil (*JAS*, 1976), Bhattacharya, Ghil & Vulis (*JAS*, 1983),
Roques *et al.* (*Phil. Trans.*, 2011, submitted)



Zonal belt with heat capacity $C(x)$ and temperature $T(t, x)$, subject to incoming radiation R_i , outgoing radiation R_o , and meridional diffusion D .

Parameter estimation for energy balance models with memory (EBMMs) – II

The memory effects are represented by a history function $H = H(t, x, T)$,

$$H(t, x, T) = \int_{-\tau}^0 \beta(s, x) T(t + s, x) ds, \quad t > 0, \quad x \in (0, 1),$$

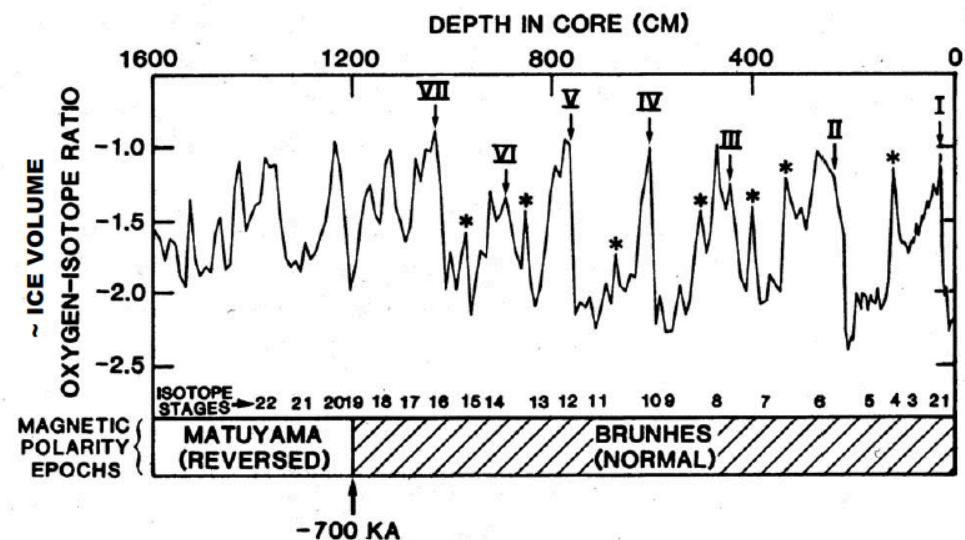
with the non-negative kernel $\beta = \beta(s, x)$ that sums to unity, thus yielding the general EBMM:

$$c(x, H(t, x, T)) \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left(k(x) \frac{\partial T}{\partial x} \right) + f(t, x, T, H(t, x, T));$$

here $f = R_i - R_o$ is the net radiation balance, affected by the past history.

The observational data come from proxy records of past temperatures and ice volume, with errors in both age-dating (**abscissa = time axis**) and “transfer function” (**ordinate = climate variable**).

Roques *et al.* (*Phil. Trans.*, 2011, submitted)



Parameter estimation for energy balance models with memory (EBMMs) – III

The initial data for this functional PDE are

$$T(s, x) = T_0(s, x), \quad s \in [-\tau, 0], \quad x \in [0, 1]$$

and we use Neumann boundary conditions at the 2 poles (or pole and equator, by symmetry). This semi-empirical EBMM requires determining coefficients from the proxy records, e.g., the ratio $\alpha = \alpha(x)$ between R_i and R_o :

$$f = f_\alpha(t, x, T, H) = f_1(t, x, T, H) + \alpha(x) f_2(t, T, H), \quad H = H(t, x, T).$$

Here $f = f[\alpha] = R_i - R_o$ is the reaction function in our reaction-diffusion model.

Under reasonable assumptions on $f_1, f_2, H, c, k, \alpha$ and β , one can prove that — given exact initial data over $-\tau \leq t \leq 0$ and exact data on T and T_x at a single point $0 < x_0 < 1$ (i.e., for a single “core”) over some interval $0 < t < t^*$ — **the coefficient $\alpha(x)$ is determined uniquely!**

But we are interested now in the more realistic situation in which a statistical model of the observation process is needed. We assume that $T(t, x) = T_0$ is the initial data with prior distribution π_1 and that the unknown coefficient α has prior distribution π_2 .

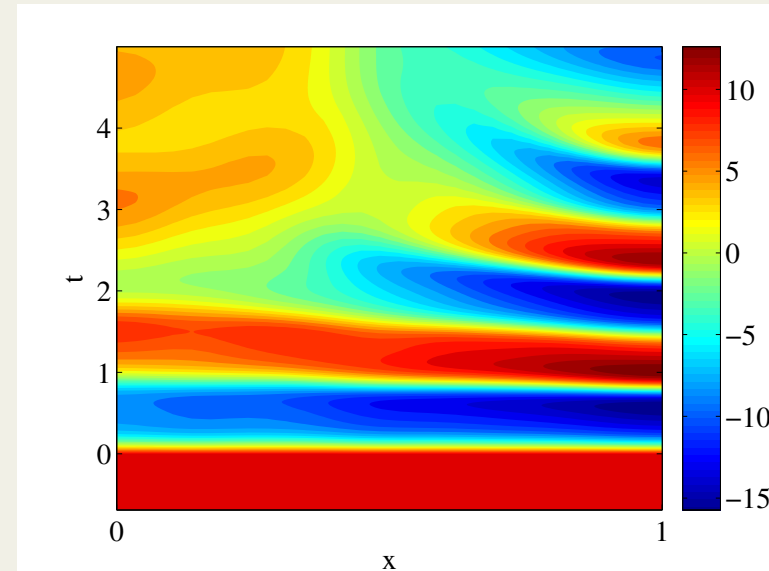
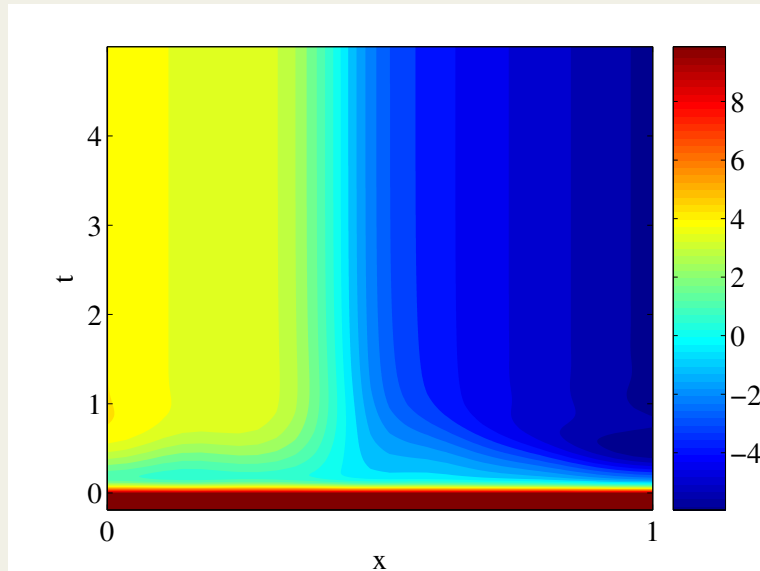
Data will be provided at three sites (cores) $S_k, k=1, 2, 3$, in the interval’s right half.

Parameter estimation for energy balance models with memory (EBMMs) – IV

The mechanistic-statistical model now includes the specific EBMM

$$\frac{\partial T}{\partial t} = \frac{\partial^2 T}{\partial x^2} + \alpha(x) (1 - a(T)) - q_0 - q_1 T - \left(\frac{1}{\tau} \int_{-\tau}^0 T(t+s) ds \right)^3,$$

where the albedo $a(t)$ is a known, piecewise-linear ramp function (cf. Sellers, 1969, and Ghil, 1976), and we study numerically the two cases $\tau = 0.2$ and $\tau = 0.7$ ky.



As expected, the solution tends rapidly to stationarity for small lag and has a longer transient, with large amplitude, for the larger lag.

The proxy records at the three sites S_k , $k=1, 2, 3$ have 2 sources of uncertainty.

Parameter estimation for energy balance models with memory (EBMMs) – V

The statistical model for these uncertainties is as follows:

(i) $Y_k(t_j)$ is the measurement of temperature T at time t_j and location S_k ,

$$Y_k(t_i) | s(t_i) \sim \text{indep. } \mathcal{N} \{T(s(t_i), S_k), \sigma^2\},$$

where σ^2 is the variance of the temperature measurement noise; and

(ii) the date of t_j is in fact $s(t_j)$, with

$$s(t_i) = \theta - \sum_{j=1}^i \eta_j \text{ with } \eta_j \sim \text{indep. } \Gamma \left(\frac{t_{j-1} - t_j}{\kappa^2}, \kappa^2 \right),$$

where Γ is the gamma distribution, $\kappa^2 > 0$ is a shape parameter, and $t_0 = \theta$.

This model is order-preserving, i.e.

$$t_i > t_j \text{ implies } s(t_i) > s(t_j),$$

$E(s(t)) = t$, and its variance increases as we “sink” further into the past.

Sobrino *et al.* (*Boreas*, 2008):
Age-depth models for 3 pollen cores
In NW Iberia.

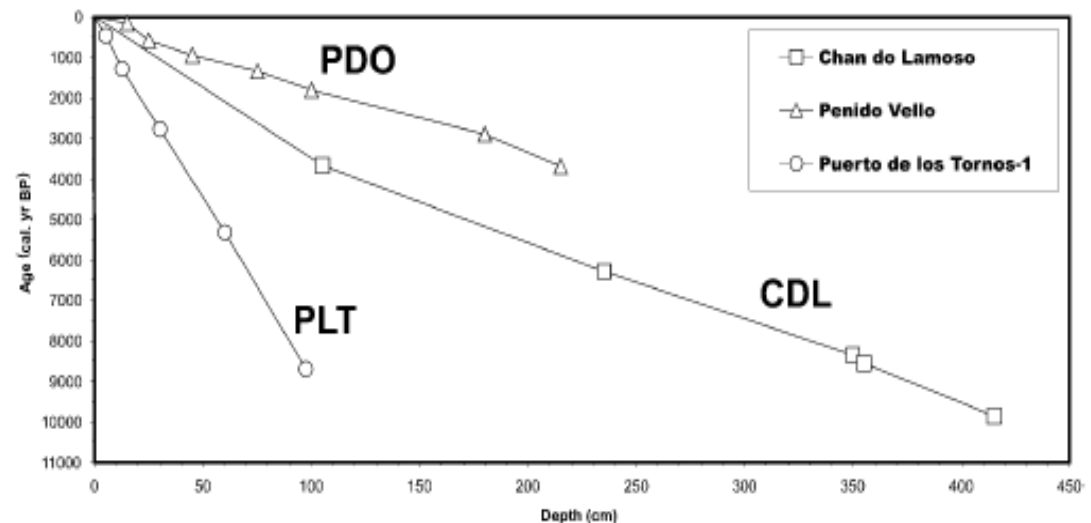
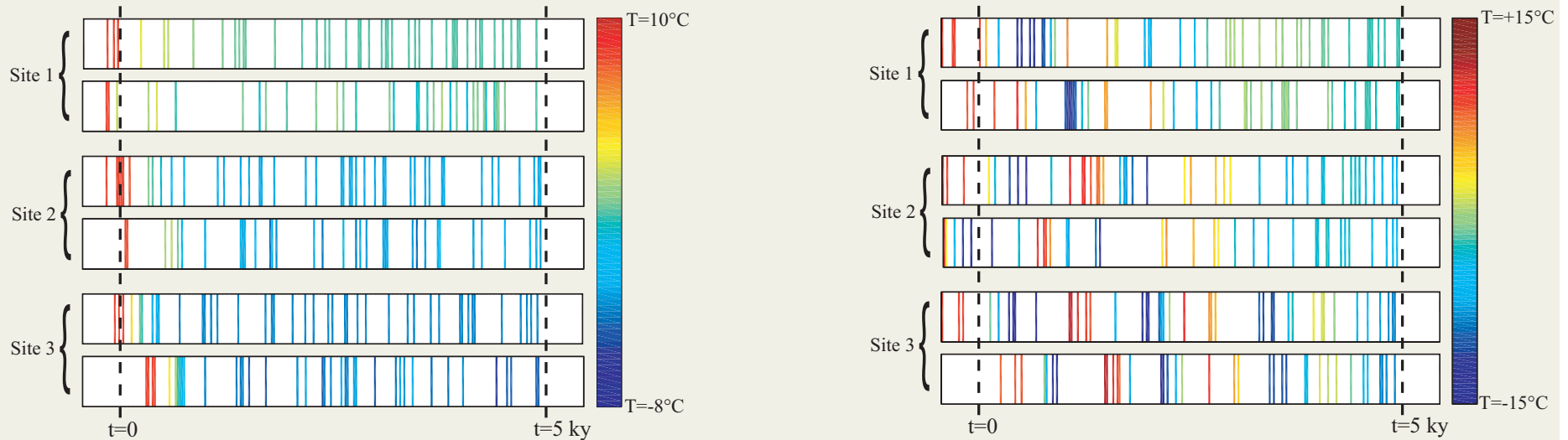


Fig. 6. Age-depth curves for the three cores analysed.

Parameter estimation for energy balance models with memory (EBMMs) – VI



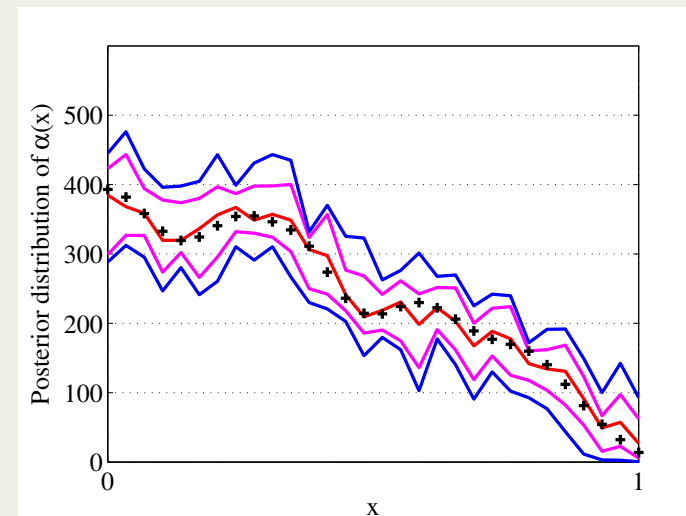
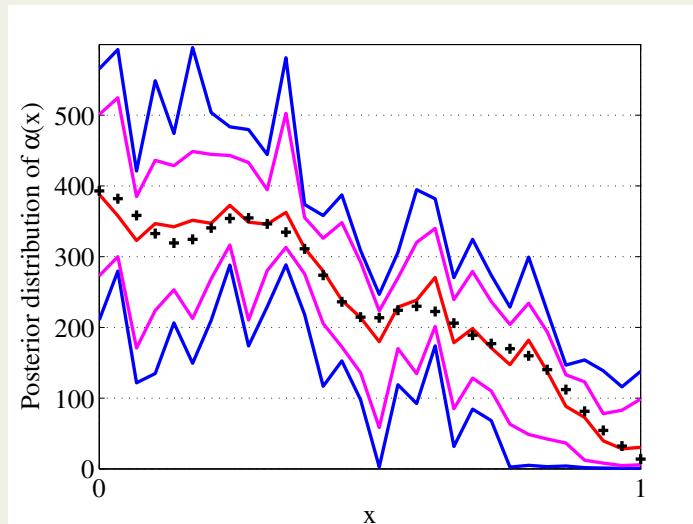
Actual temperatures vs. measured temps, at the 3 sites. At each of them, the upper row corresponds to the actual T s at the actual times, while the lower row corresponds to the measured T s at the estimated times. (a) $\tau = 0.2$ ky, and (b) $\tau = 0.7$ ky. Clearly the errors in both the estimates of T s and times are larger for the larger delay, which resulted in the more irregular solution.

We seek the coefficient $\alpha(x)$ by a Bayesian approach, assume uniform prior distributions for T_0 and for $\alpha(x)$, and draw a sample from the joint posterior distribution of $(T_0, \alpha(x))$ by Markov chain Monte Carlo (MCMC).

Roques *et al.* (2011, submitted)

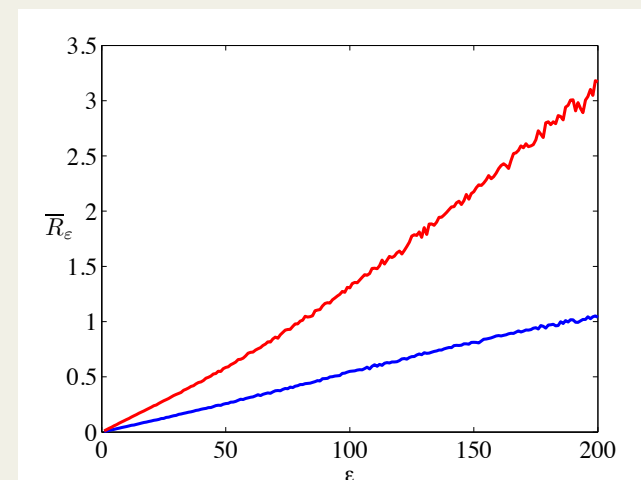
Parameter estimation for energy balance models with memory (EBMMs) – VII

The results are shown below:



Estimates of the coefficient $\alpha(x)$: posterior median (red), first and last deciles (magenta), first and last percentiles (blue); the true values are the + signs. (a) $\tau = 0.2$ ky, and (b) $\tau = 0.7$ ky. Clearly the estimates are better in (b).

The last figure shows the average L_2 -response \bar{R}_ε of our EBMM model to random perturbations $\alpha'(x)$ in $\alpha(x)$ drawn from a random field A with std. dev. ε , over the interval $0 < t < 5$ ky: blue for $\tau = 0.2$ ky, and red for $\tau = 0.7$ ky. Both curves show linear response, but the red one has double the slope.



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Concluding remarks

We've come a long way in 30 years — some advances are laborious and incremental (e.g., sequential vs. control-theoretical methods), but others are fresh and exciting.

The latter include new areas of application

- biology, paleoclimate, space physics, ...;

as well as novel methodological challenges

- multi-scale and multi-model problems

- inverse problems for evolution equations, ...

Technological advances both pose new problems (more data, higher resolution, ...) and help solve them.

Overall, it's a brave new world, in which data and models actively speak to each other, and we do so to both: enjoy!

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General references

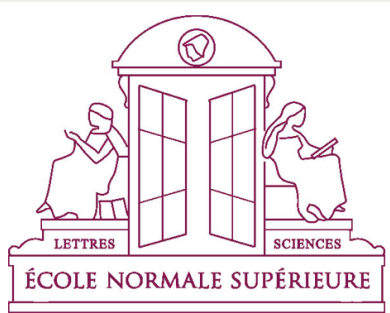
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Empirical Model Reduction and Applications



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