

Isogeometric analysis for nearly incompressible materials

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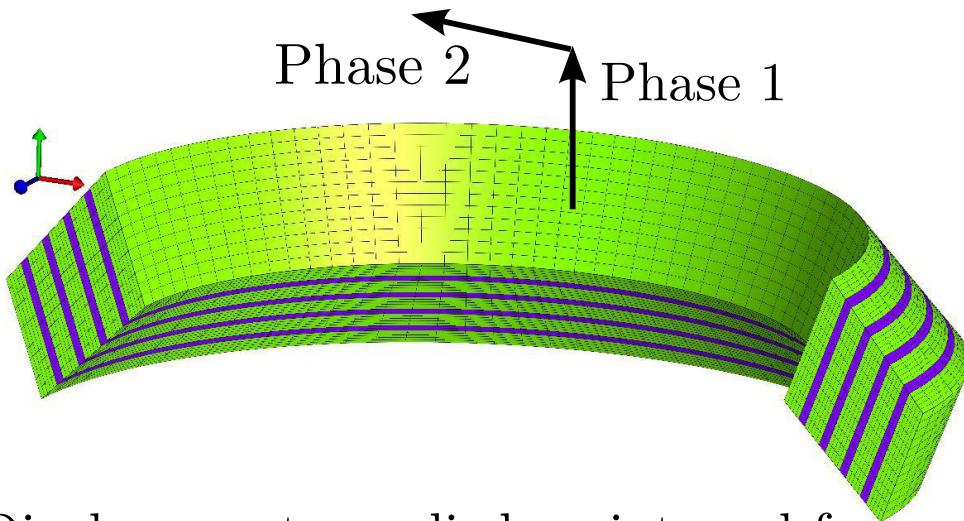
Collaborators: Pablo Antolín, Andrea Bressan,
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In collaboration and with support by: Hutchinson-Total SA

LMS - EPSRC Durham Research Symposia

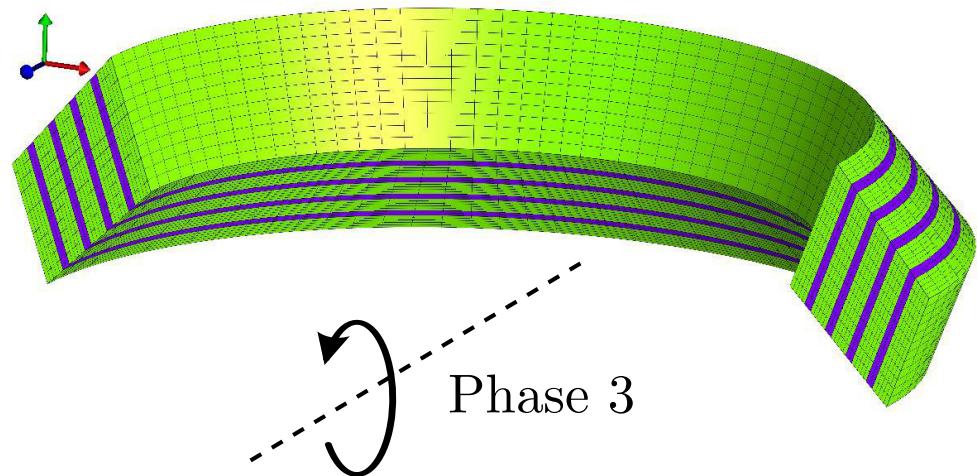
July 9, 2014

Rubber-metal conical bearing



Displacements applied on internal face

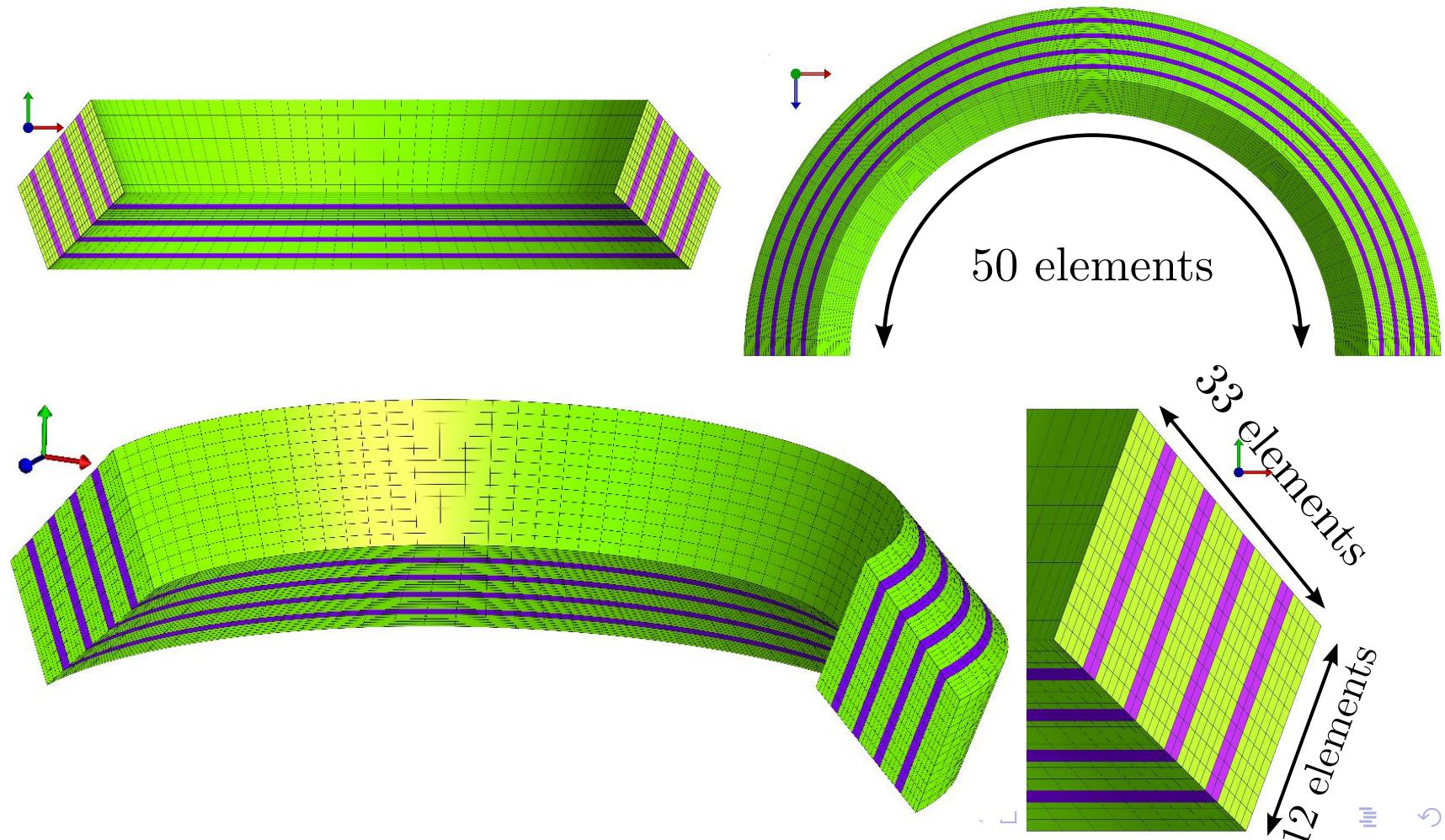
External face is encasted



Rubber-metal conical bearing

19800 elements, 22542 nodes \rightarrow 67626 dofs

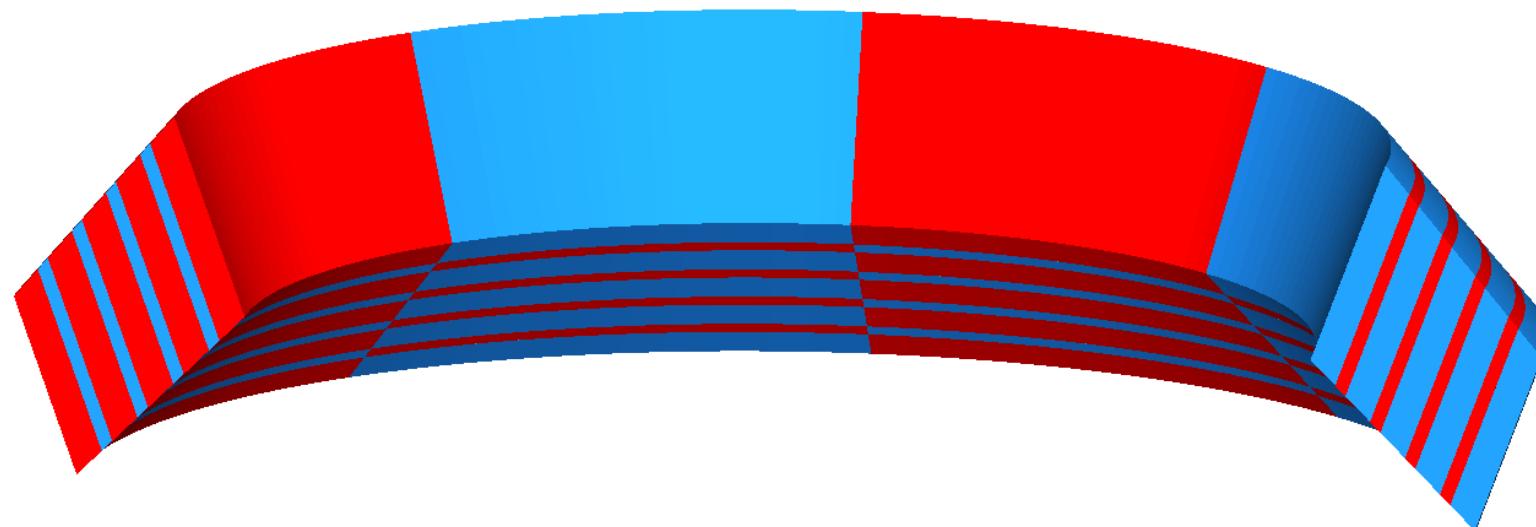
Symmetry conditions imposed



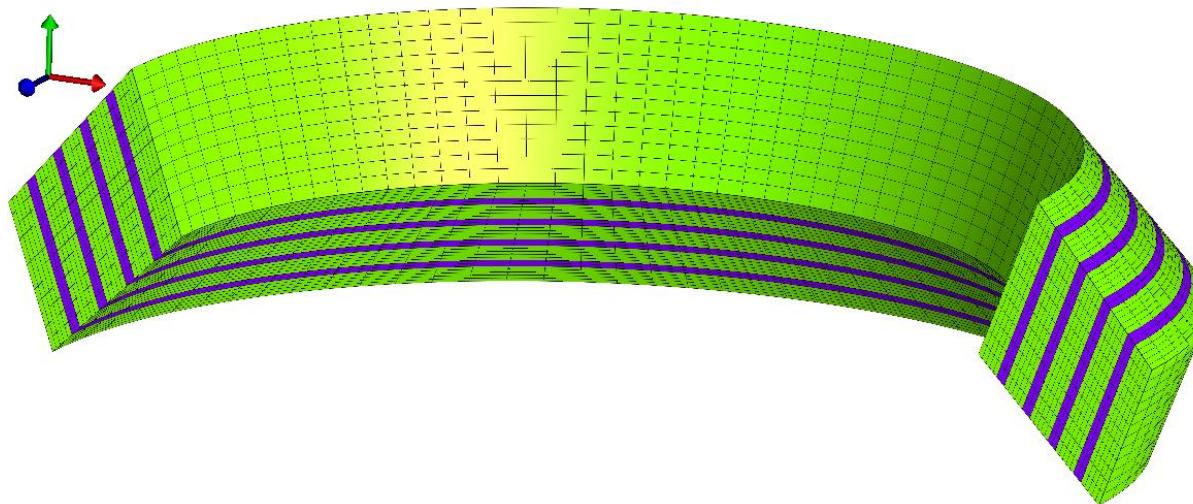
Rubber-metal conical bearing

36 elements ($4 \times 1 \times 9$), $\begin{cases} p = 2 \rightarrow 567 \text{ nodes}, & 1701 \text{ dofs} \\ p = 3 \rightarrow 1296 \text{ nodes}, & 3888 \text{ dofs} \end{cases}$

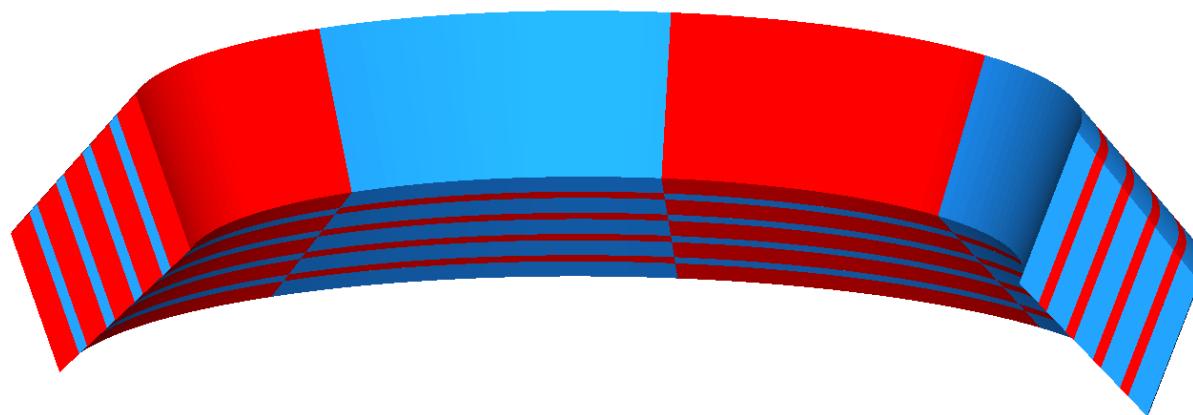
Each layer is a different NURBS patch



Conical bearing: FEM vs IGA mesh



19800 elements
67626 dofs



36 elements
 $p = 2 \rightarrow 1701$ dofs
 $p = 3 \rightarrow 3888$ dofs

Materials and numerical methods

Steel:

$$E = 210 \text{ GPa}, \nu = 0.3$$

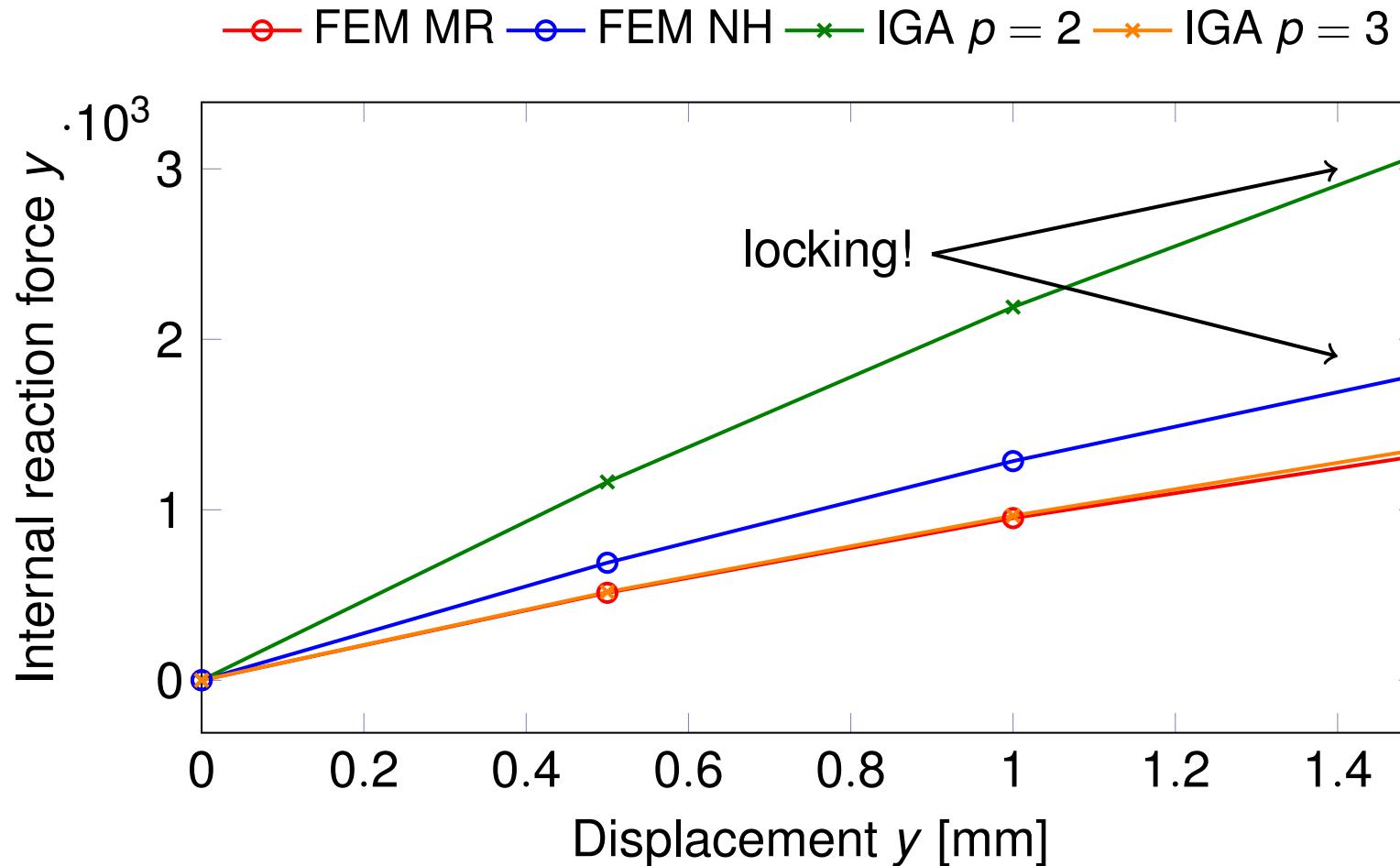
- Isotropic elasticity formulation used.

Rubber:

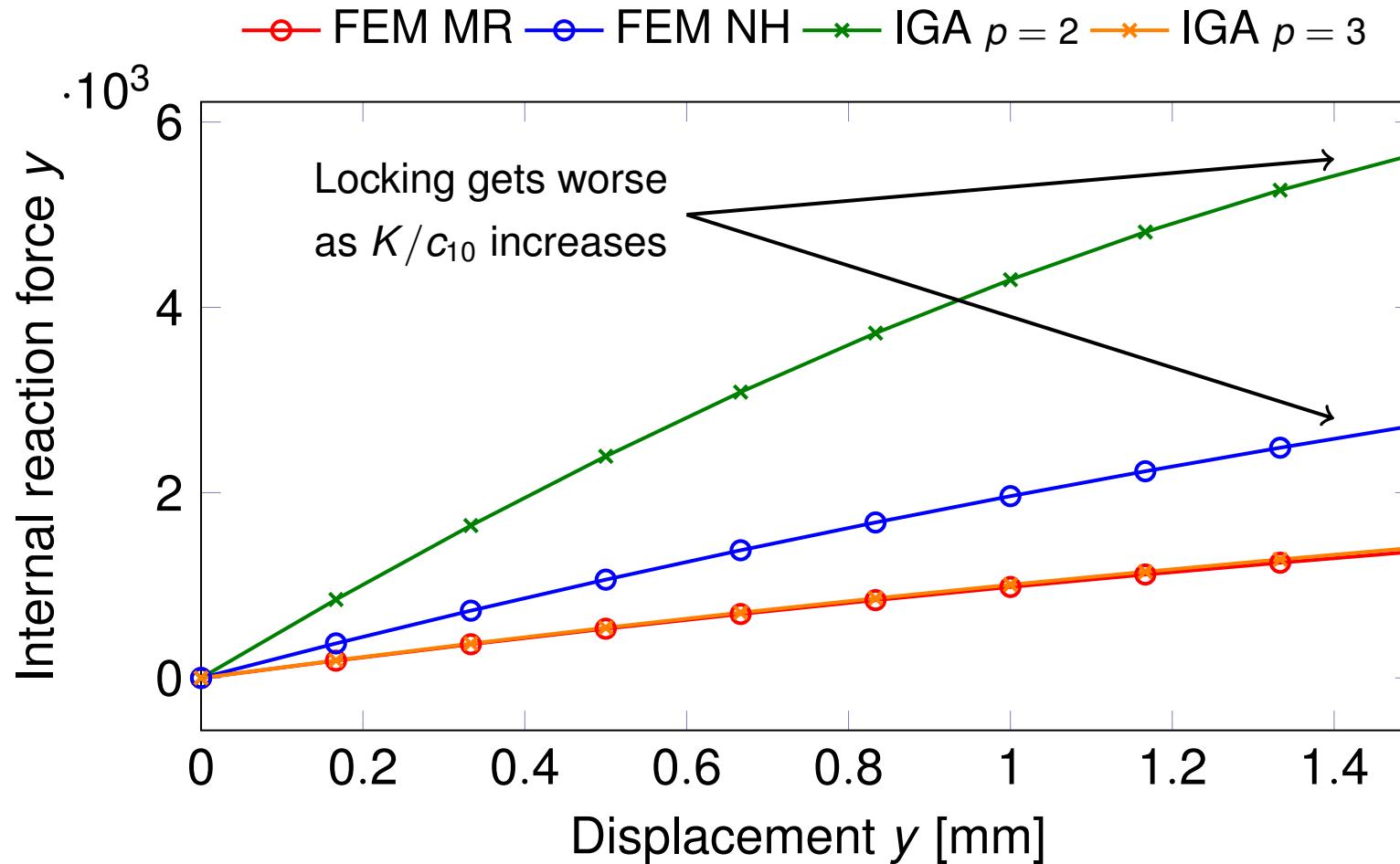
$$C_{10} = 1 \text{ MPa}, K = 1000 \text{ MPa} (E = 5.996 \text{ MPa}, \nu = 0.499)$$

- For FEM:
 - ▶ Neo-Hookean with plain Galerkin formulation → solution locks
 - ▶ Mooney–Rivlin (three field: displacement+pressure+volume ratio) implemented with “selective-reduced integration” → suitable for incompressible materials.
- For IGA: Neo–Hookean with plain Galerkin formulation.

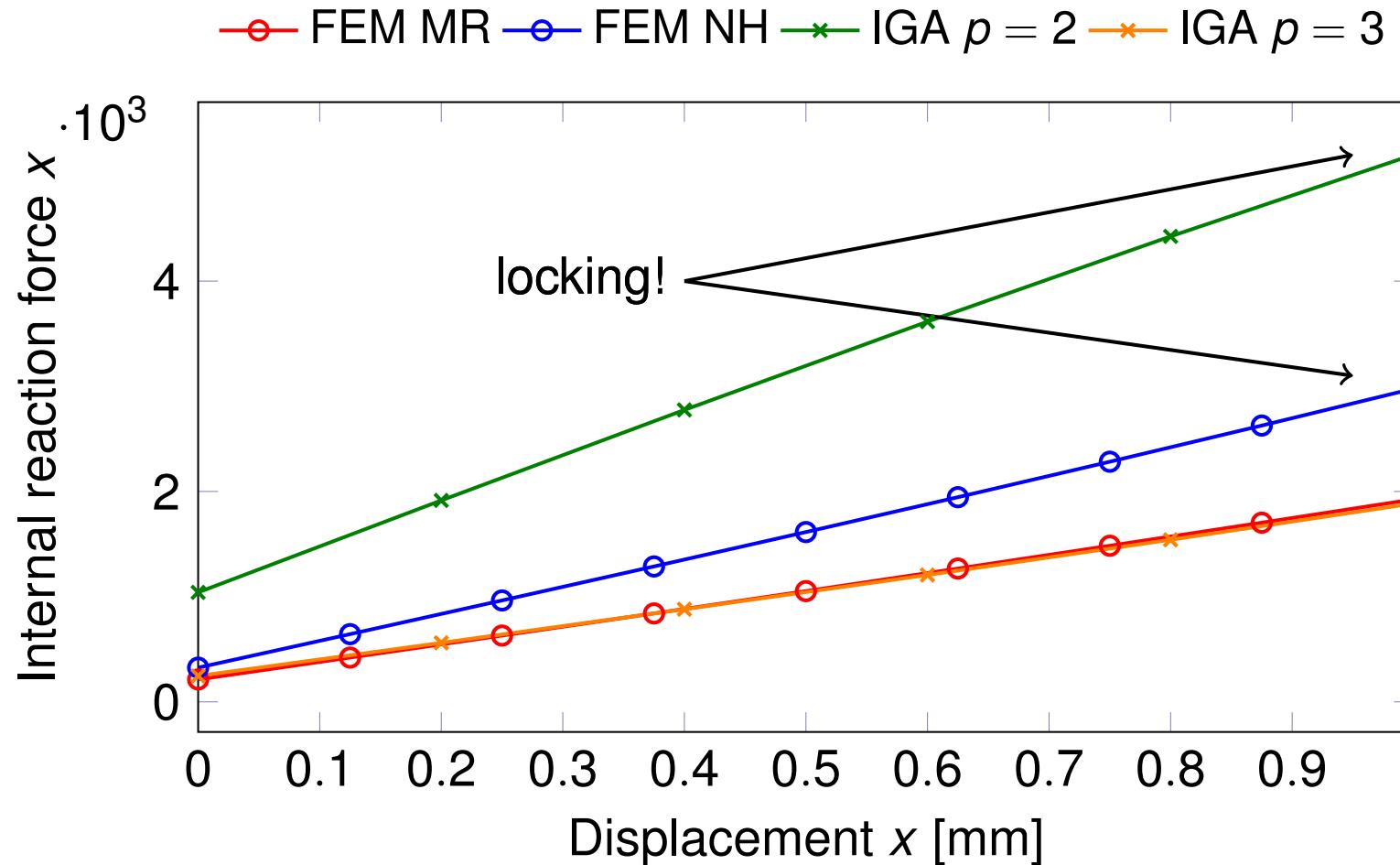
Conical bearing: phase 1



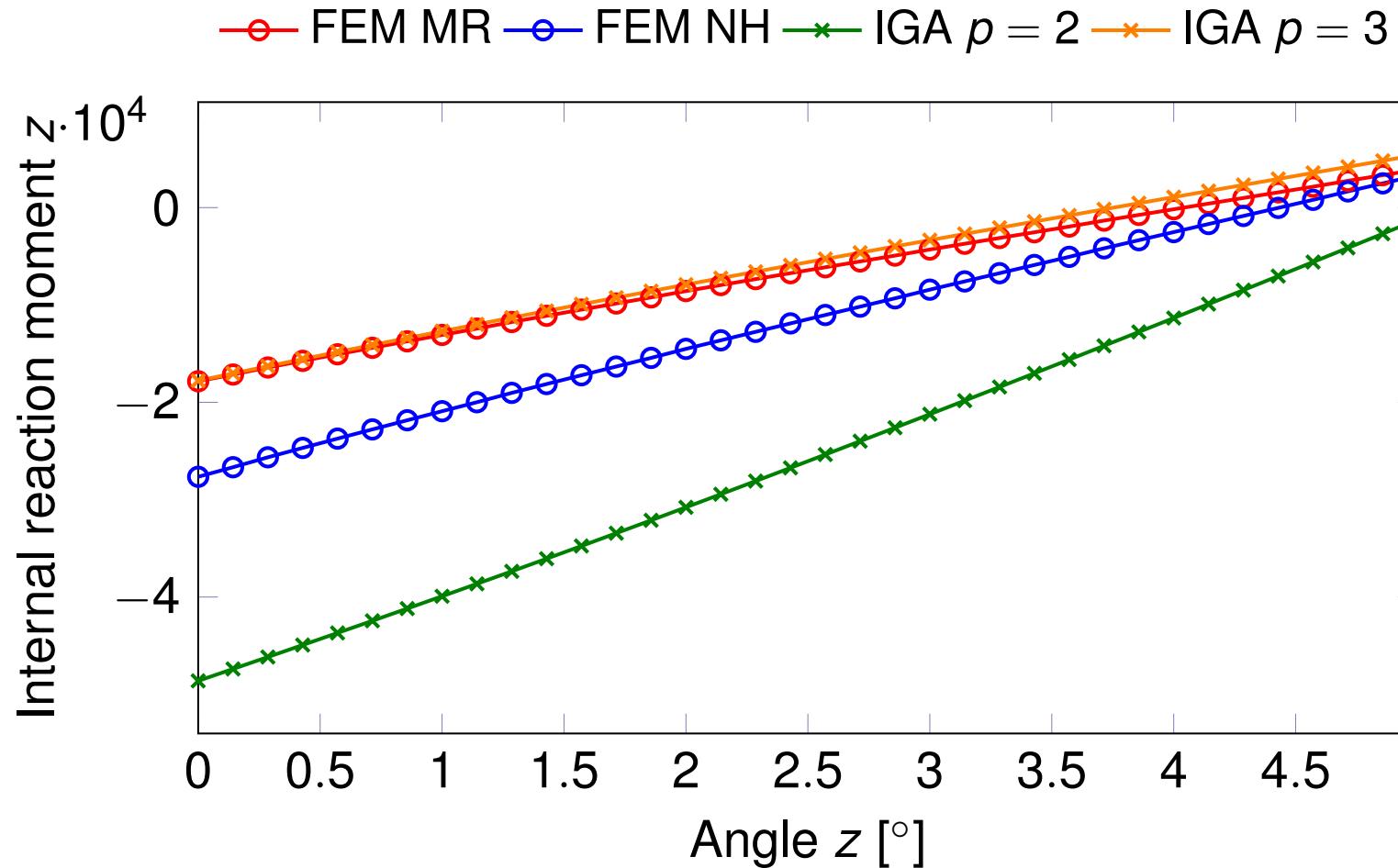
Conical bearing: $K/c_{10} = 3000$. Phase 1



Conical bearing: Phase 2

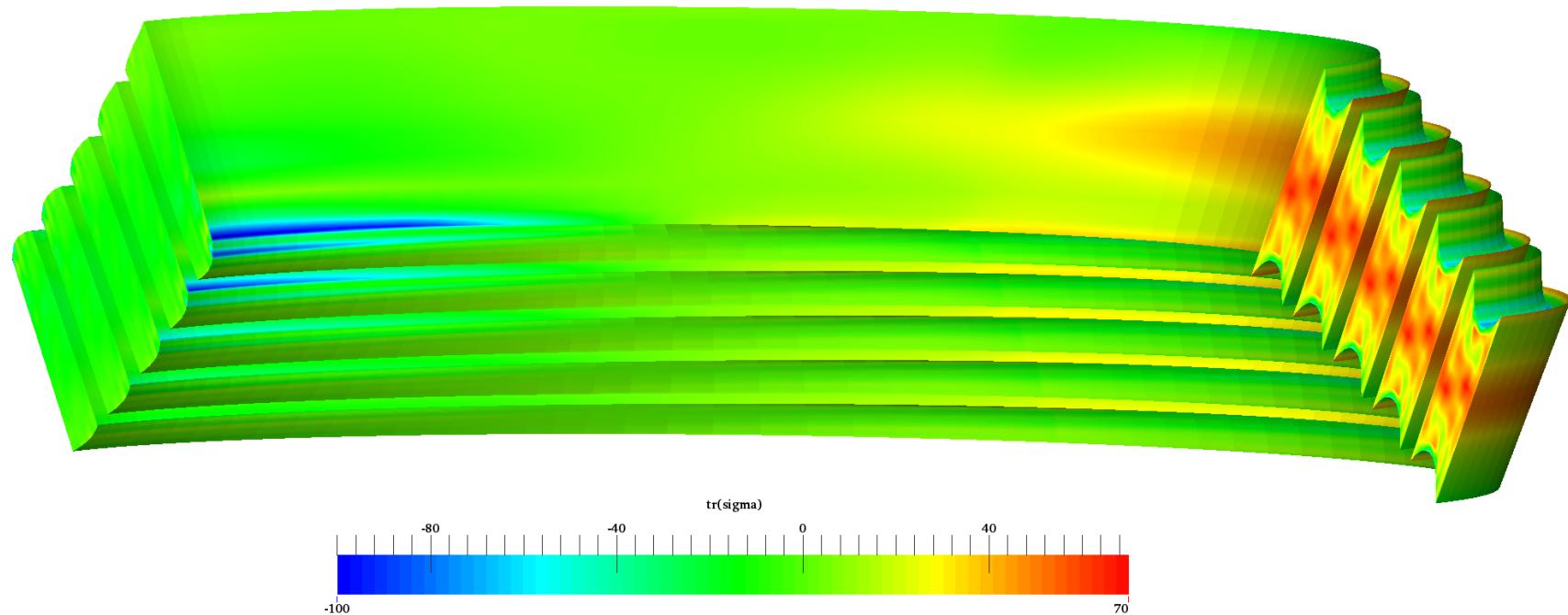


Conical bearing: Phase 3



Conical bearing: stress oscillations for IGA $p = 3$

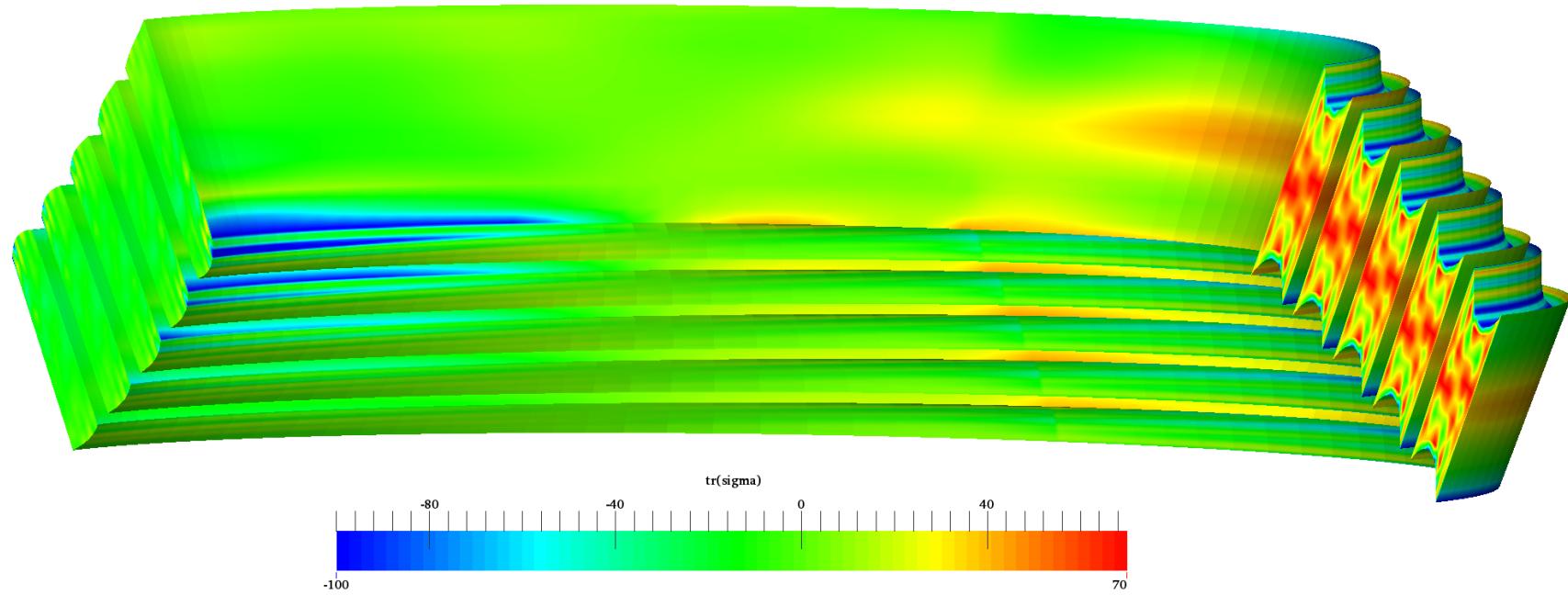
Oscillations appear in $\sigma^{\text{vol}} = \frac{1}{3} \text{tr}(\sigma) \mathbf{1}$



$\frac{1}{3} \text{tr}(\sigma)$ plotted

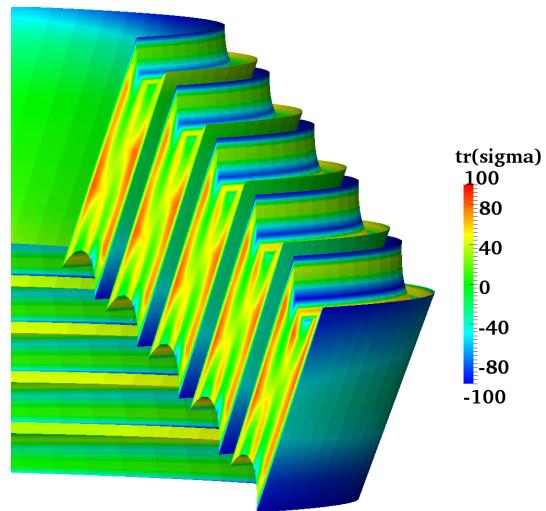
Conical bearing: stress oscillations for IGA $p = 3$

$$K/c_{10} = 3000$$

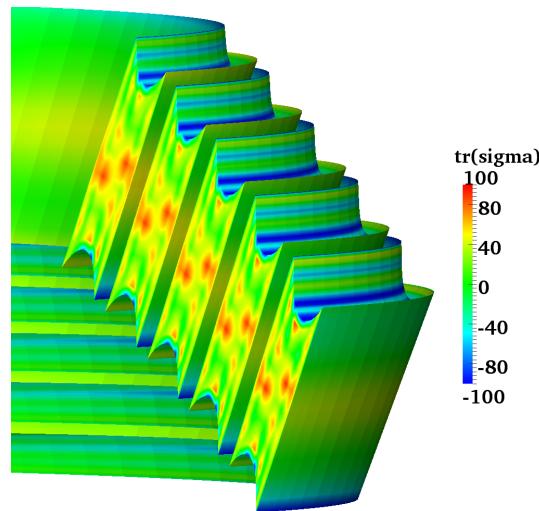


The scale is the same, the oscillations increase with K/c_{10}

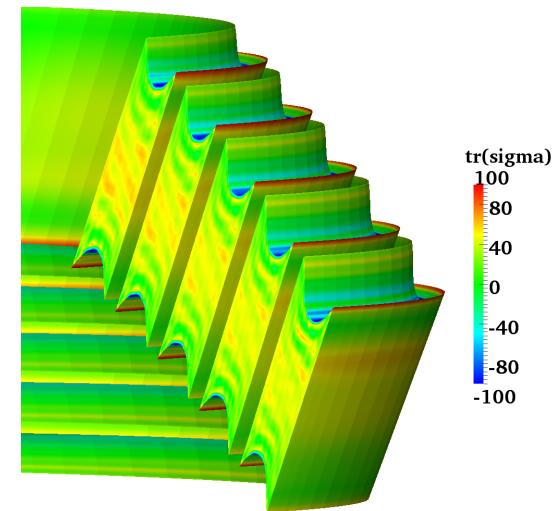
Conical bearing: stress oscillations for IGA $p = 3$



$4 \times 2 \times 2$
elements / layer



$4 \times 4 \times 4$
elements / layer



$4 \times 8 \times 8$
elements / layer

Linear elasticity model problem

Strong form problem

$$\nabla \cdot \boldsymbol{\sigma} + \mathbf{f} = \mathbf{0} \quad \text{in } \Omega$$

$$\mathbf{u} = \bar{\mathbf{u}} \quad \text{on } \Gamma_D$$

$$\boldsymbol{\sigma} \cdot \mathbf{n} = \mathbf{t} \quad \text{on } \Gamma_N$$

Linear elasticity model problem

Isotropic linear elasticity

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$$\boldsymbol{\sigma} = 2\mu \boldsymbol{\varepsilon} + \lambda \nabla \cdot \mathbf{u} \mathbf{1}$$

$$\boldsymbol{\varepsilon} = \nabla^s \mathbf{u}$$

$$\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}$$

$$\mu = \frac{E}{2(1+\nu)}$$

$$\nu \rightarrow 1/2, \quad \lambda \rightarrow \infty$$

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Weak form: find $\mathbf{u} \in (H^1(\Omega))^3$ with $\mathbf{u} = \bar{\mathbf{u}}$ on Γ_D and such that

$$\int_{\Omega} \nabla^s \mathbf{w} : \boldsymbol{\sigma} \, d\Omega = \int_{\Omega} \mathbf{w} \cdot \mathbf{f} \, d\Omega + \int_{\Gamma_N} \mathbf{w} \cdot \mathbf{t} \, d\Gamma, \quad \forall \mathbf{w} \in (H_{\Gamma_D}^1(\Omega))^3$$

Linear elasticity model problem

Isotropic linear elasticity

Strong form problem

$$\begin{aligned}\nabla \cdot \boldsymbol{\sigma} + \mathbf{f} &= \mathbf{0} && \text{in } \Omega \\ \mathbf{u} &= \bar{\mathbf{u}} && \text{on } \Gamma_D \\ \boldsymbol{\sigma} \cdot \mathbf{n} &= \mathbf{t} && \text{on } \Gamma_N\end{aligned}$$

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$$\underbrace{\int_{\Omega} \nabla^s \mathbf{w} : \boldsymbol{\sigma} \, d\Omega}_{a(\mathbf{w}, \mathbf{u})} = \underbrace{\int_{\Omega} \mathbf{w} \cdot \mathbf{f} \, d\Omega + \int_{\Gamma_N} \mathbf{w} \cdot \mathbf{t} \, d\Gamma}_{L(\mathbf{w})}, \quad \forall \mathbf{w} \in (H_{\Gamma_D}^1(\Omega))^3$$

Linear elasticity model problem

Isotropic linear elasticity

Strong form problem

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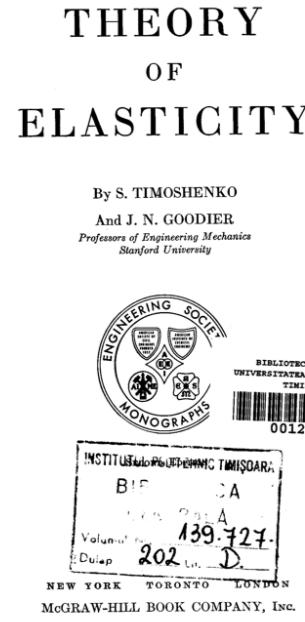
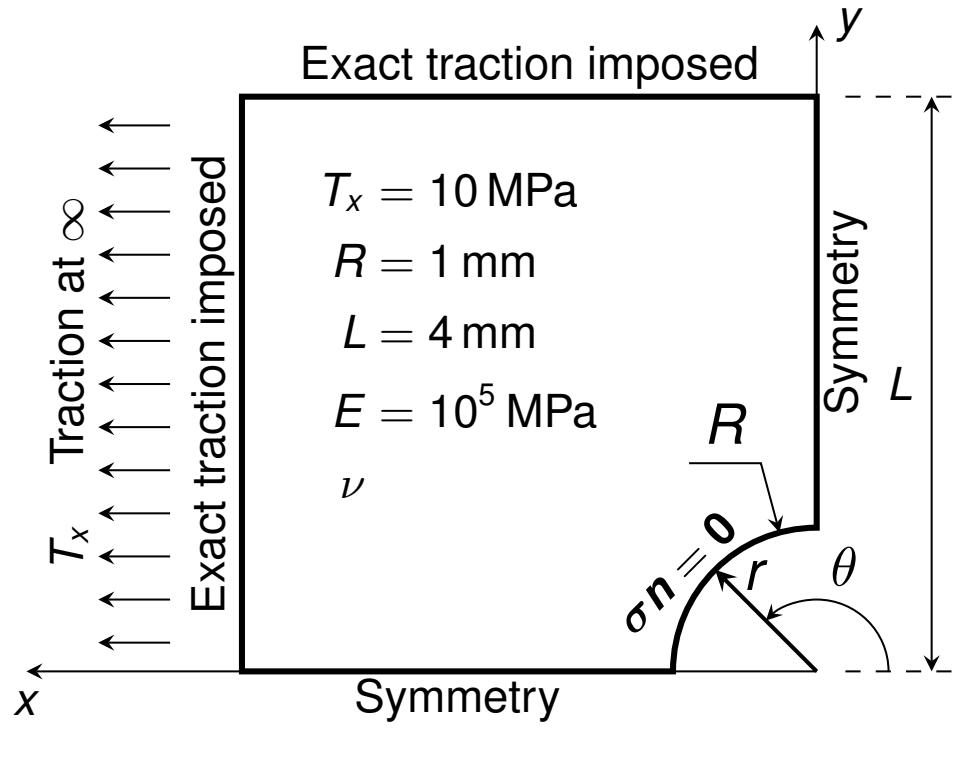
$$\underbrace{\int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \nabla \cdot \mathbf{w} \nabla \cdot \mathbf{u} d\Omega}_{a(\mathbf{w}, \mathbf{u})} = L(\mathbf{w}), \quad \forall \mathbf{w} \in (H_{\Gamma_D}^1(\Omega))^3$$

Target method

Rubber industry is interested in:

- A “locking-free” method: optimal order of convergence for the displacement and no spurious oscillations in the stress, in the range $\nu \in [0.3, 0.4999]$
- Stiffness matrix
 - ▶ Symmetric
 - ▶ Definite positive
- Efficient

Classical benchmark: plate with a hole



$$\sigma_{rr}(r, \theta) = \frac{T_x}{2} \left[1 - \frac{R^2}{r^2} + \left(1 - 4 \frac{R^2}{r^2} + 3 \frac{R^4}{r^4} \right) \cos 2\theta \right],$$

$$\sigma_{\theta\theta}(r, \theta) = \frac{T_x}{2} \left[1 + \frac{R^2}{r^2} - \left(1 + 3 \frac{R^4}{r^4} \right) \cos 2\theta \right],$$

$$\sigma_{r\theta}(r, \theta) = -\frac{T_x}{2} \left(1 + 2\frac{R^2}{r^2} - 3\frac{R^4}{r^4} \right) \sin 2\theta ,$$

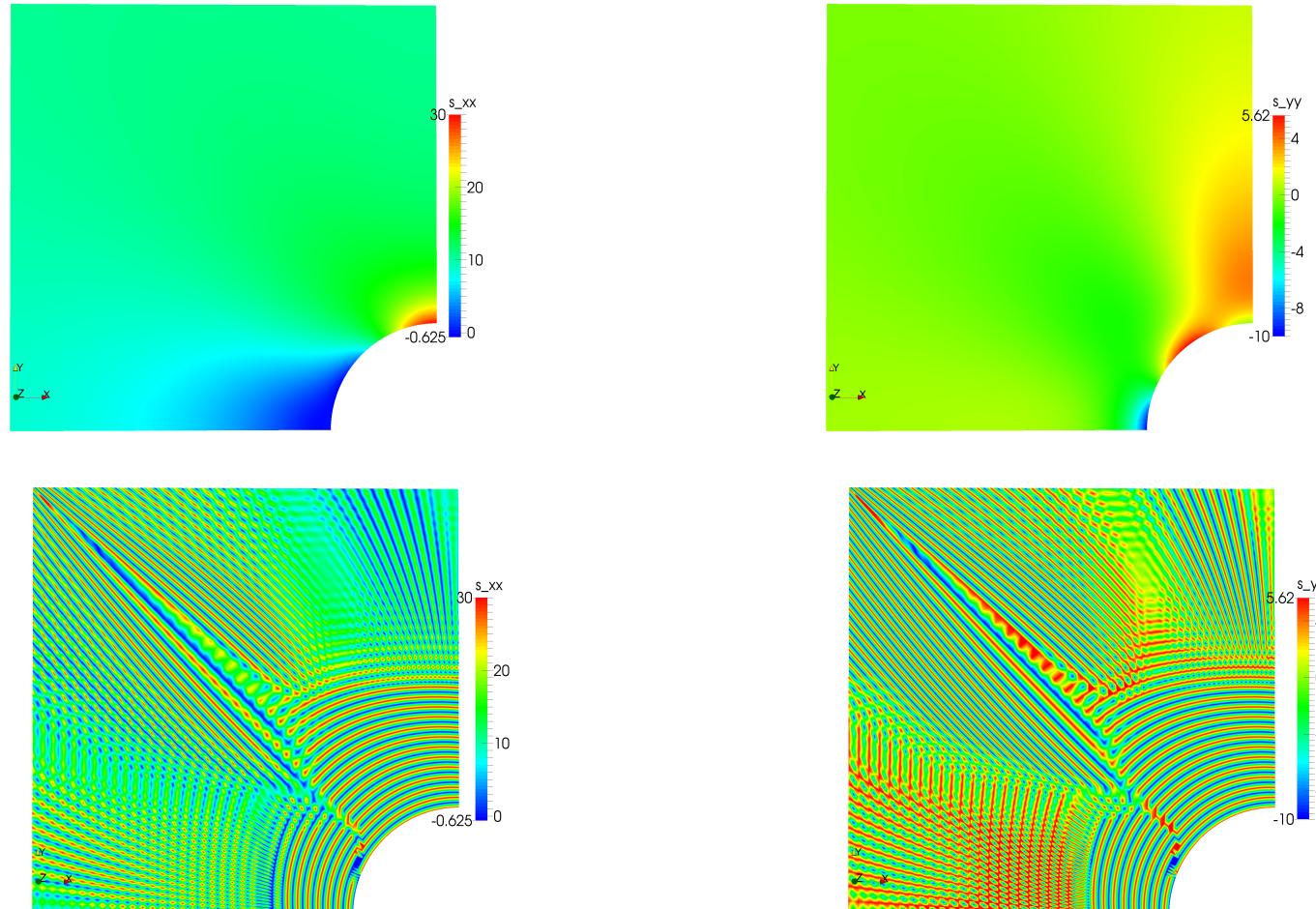
Standard formulation

Plain Galerkin (displacement) formulation

$$a(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \nabla \cdot \mathbf{w} \nabla \cdot \mathbf{u} d\Omega$$

- Stress oscillations and locking
- Symmetric ✓
- Sparse ✓
- Definite positive ✓

Exact vs plain formulation for $\nu = 0.49999$



Isogeometric \bar{B} projection technique

\bar{B} projection technique [Elguedj, Bazilevs, Calo, and Hughes, 2008]

$$a(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \pi(\nabla \cdot \mathbf{w}) \pi(\nabla \cdot \mathbf{u}) d\Omega$$

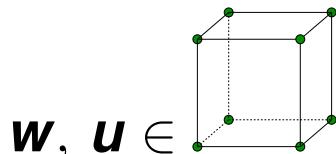
$$\pi(\phi)(\mathbf{x}) = \sum_i \tilde{N}_i(\mathbf{x}) c_i(\phi) \quad L^2\text{-proj. on lower degree splines}$$

Isogeometric \bar{B} projection technique

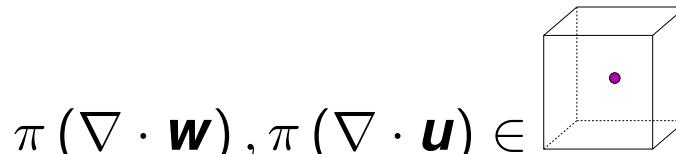
\bar{B} projection technique [Elguedj, Bazilevs, Calo, and Hughes, 2008]

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$\mathbf{w}, \mathbf{u} \in$... then perform degree-elevation and knot-insertion



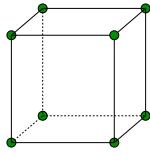
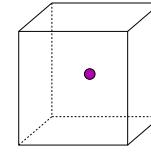
$$\pi(\nabla \cdot \mathbf{w}), \pi(\nabla \cdot \mathbf{u}) \in$$

Isogeometric \bar{B} projection technique

\bar{B} projection technique [Elguedj, Bazilevs, Calo, and Hughes, 2008]

$$a(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \pi(\nabla \cdot \mathbf{w}) \pi(\nabla \cdot \mathbf{u}) d\Omega$$

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$\mathbf{w}, \mathbf{u} \in$ 
 $\pi(\nabla \cdot \mathbf{w}), \pi(\nabla \cdot \mathbf{u}) \in$ 
... then perform degree-elevation and knot-insertion

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$$\pi(\phi)(\mathbf{x}) = \sum_i \tilde{N}_i(\mathbf{x}) c_i(\phi) \quad L^2\text{-proj. on lower degree splines}$$

isogeometric \bar{B} -method

$$\mathbf{u} \in \mathcal{S}_{p-1}^p \times \mathcal{S}_{p-1}^p \times \mathcal{S}_{p-1}^p \text{ and } \pi(\cdot) \in \mathcal{S}_{p-2}^{p-1}$$

Isogeometric \bar{B} projection technique

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$$\pi(\phi)(\mathbf{x}) = \sum_i \tilde{N}_i(\mathbf{x}) c_i(\phi) \quad L^2\text{-proj. on lower degree splines}$$

$$c_i(\phi) = \tilde{\mathbf{M}}_{ij}^{-1} \int_{\Omega} \tilde{N}_j(\mathbf{x}) \phi(\mathbf{x}) d\Omega$$

- Unlocked solution ✓
- Symmetric ✓
- full matrix
- Definite positive ✓

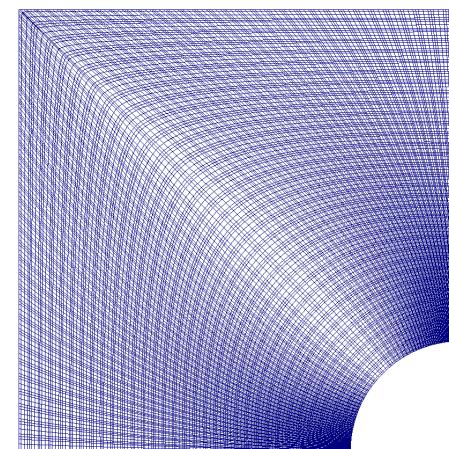
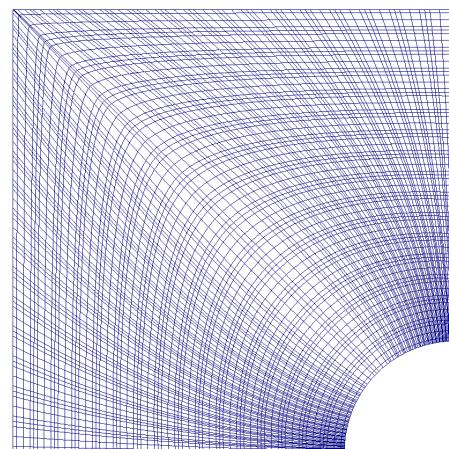
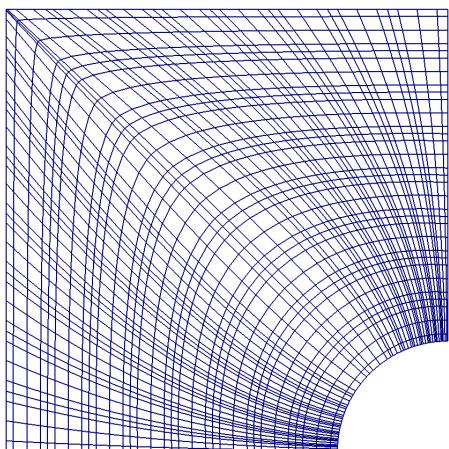
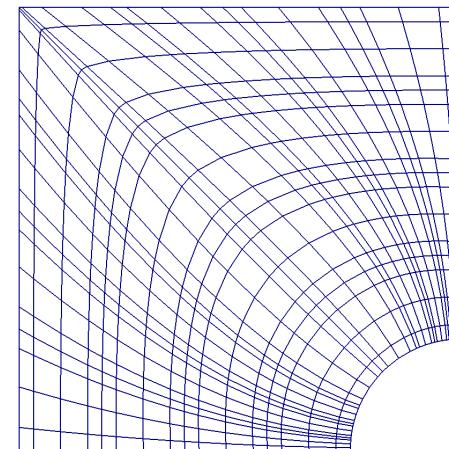
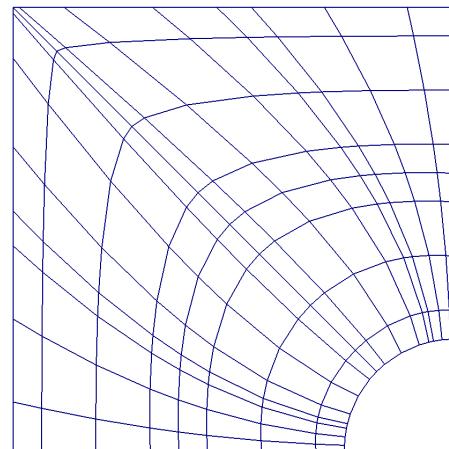
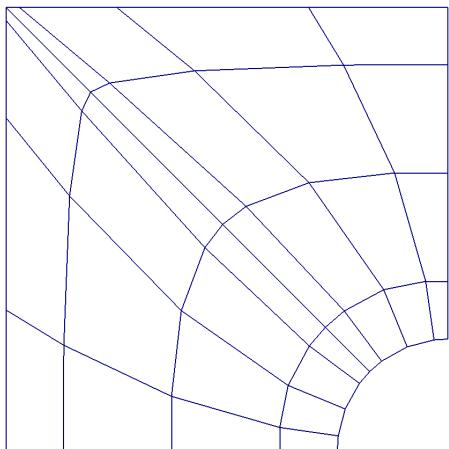
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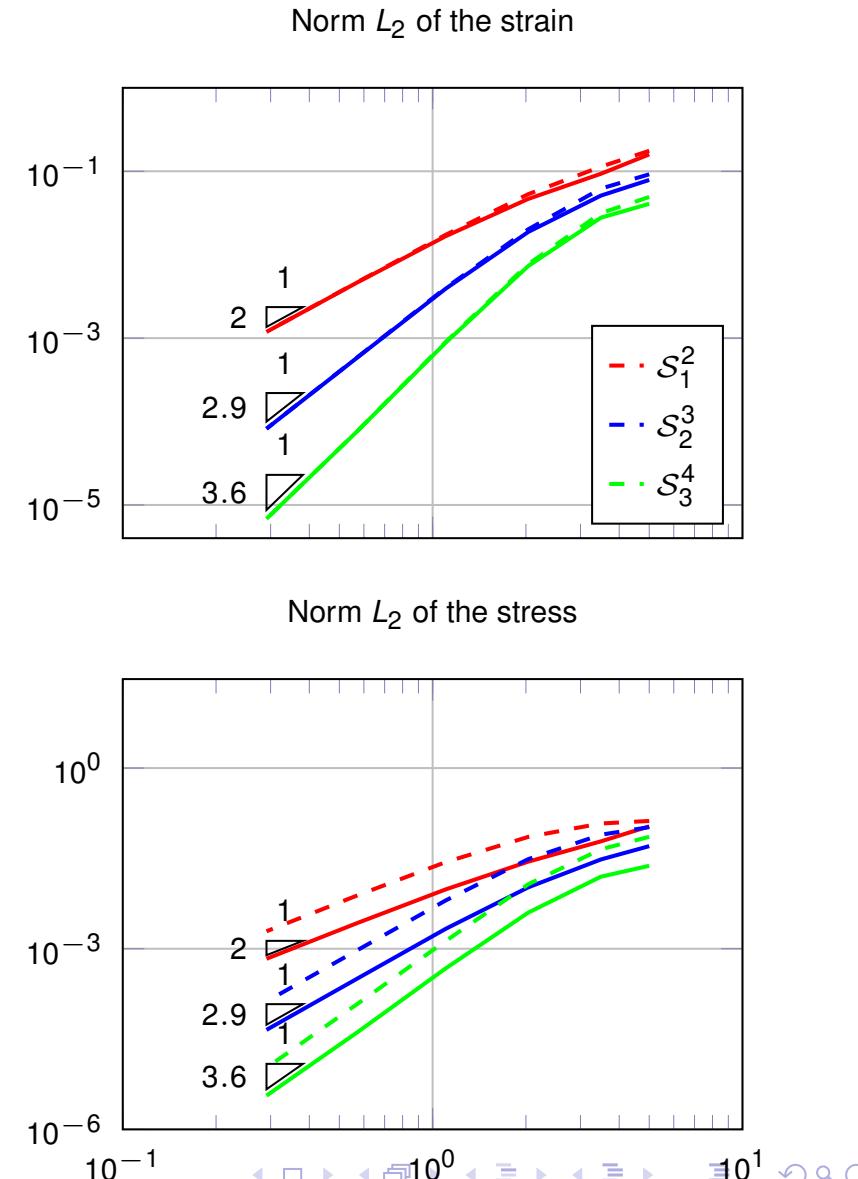
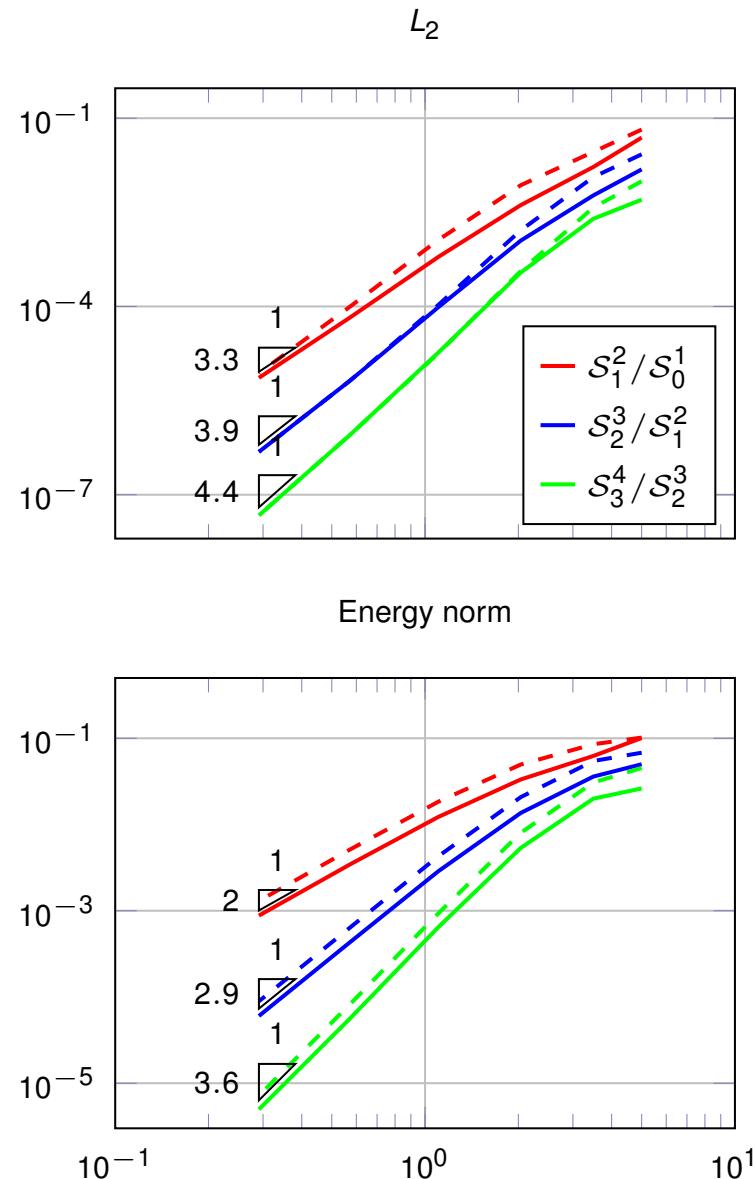
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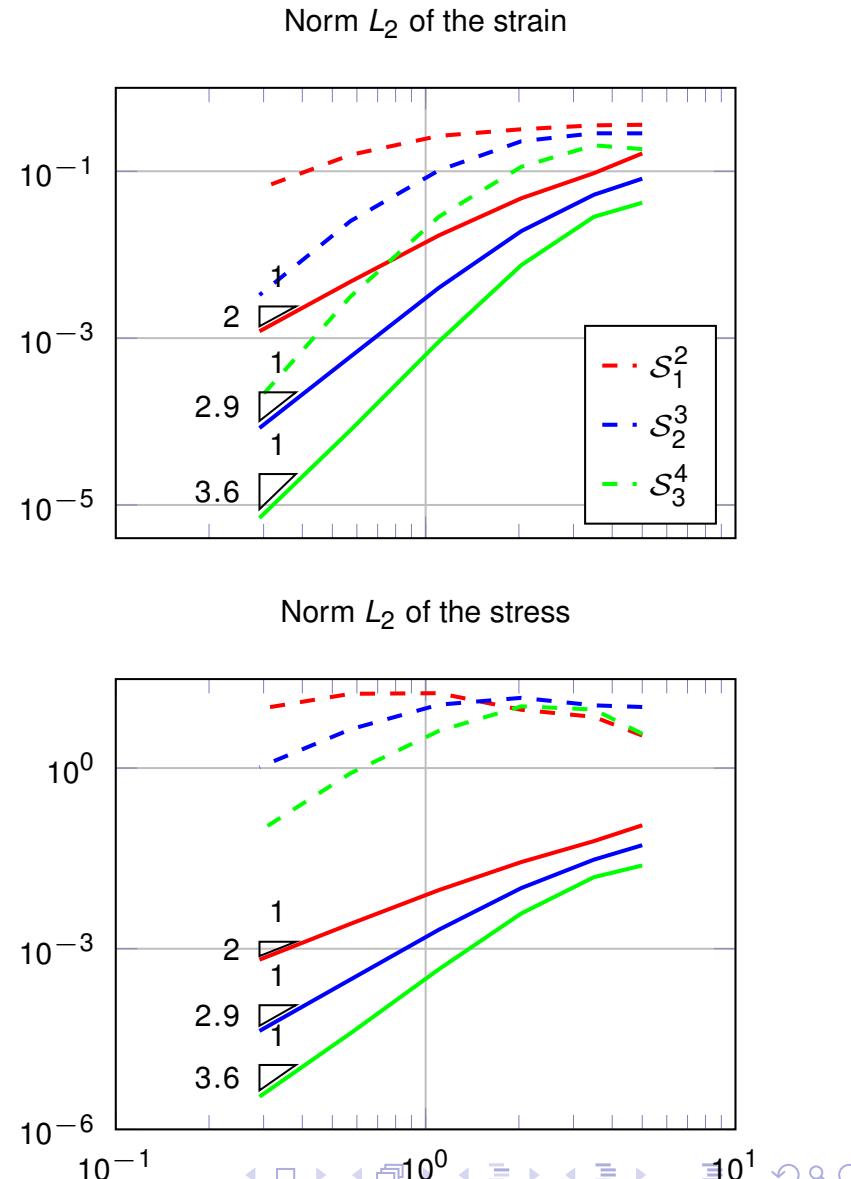
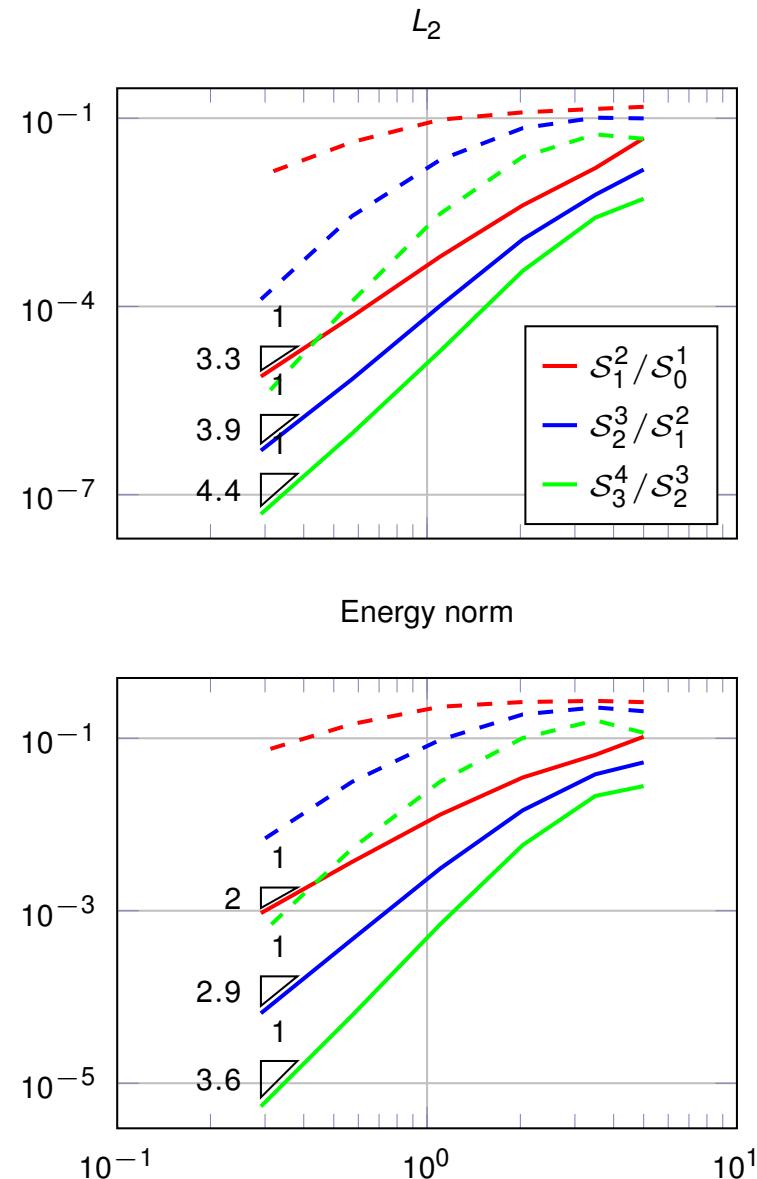
mesh refinement in our next tests



\bar{B} for $\nu = 0.4$



\bar{B} for $\nu = 0.49999$



\bar{B} with lumped mass matrix

\bar{B} with lumped mass matrix: projection technique

$$a^D(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \pi^D(\nabla \cdot \mathbf{w}) \pi^D(\nabla \cdot \mathbf{u}) d\Omega a(\mathbf{w}, \mathbf{u}) :=$$

$$\pi^D(\phi)(\mathbf{x}) = \sum_{i=1}^N \tilde{N}_i(\mathbf{x}) c_i(\phi)$$

$$c_i(\phi) = [\tilde{\mathbf{M}}_{jj}^D]^{-1} \int_{\Omega} \tilde{N}_j(\mathbf{x}) \phi(\mathbf{x}) d\Omega$$

- Only first-order convergence
- Symmetric ✓
- Sparse ✓
- Definite positive ✓

\bar{B} with lumped mass matrix

\bar{B} with lumped mass matrix: projection technique

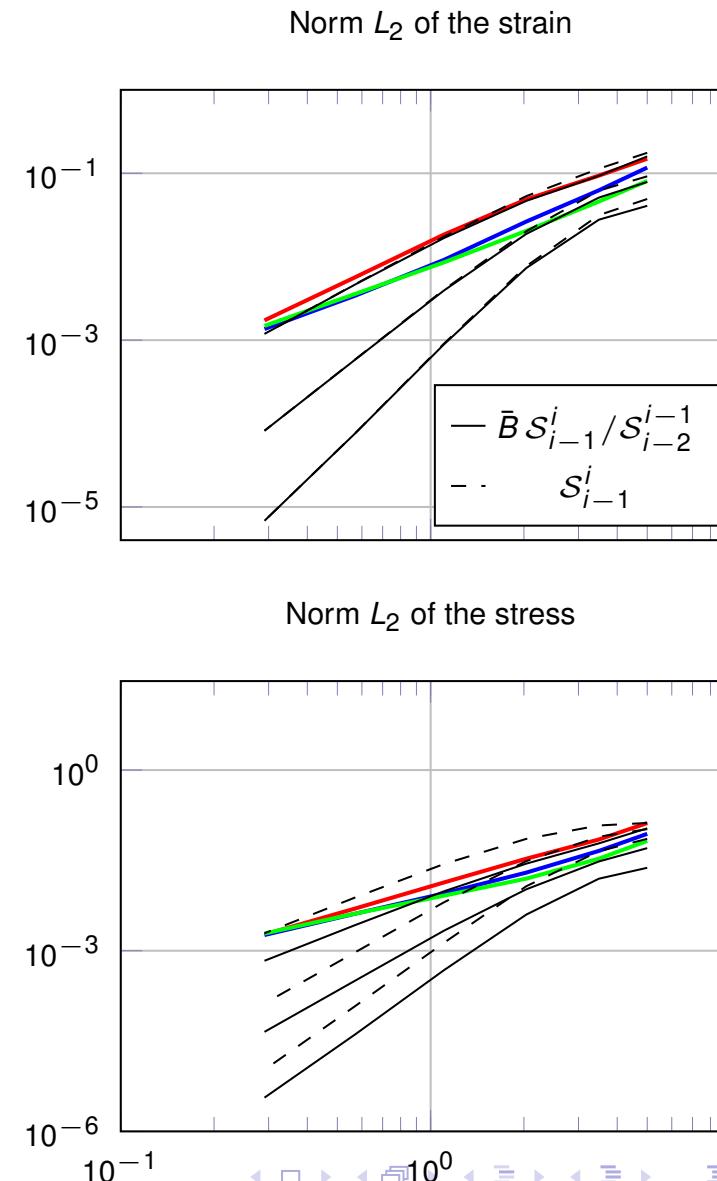
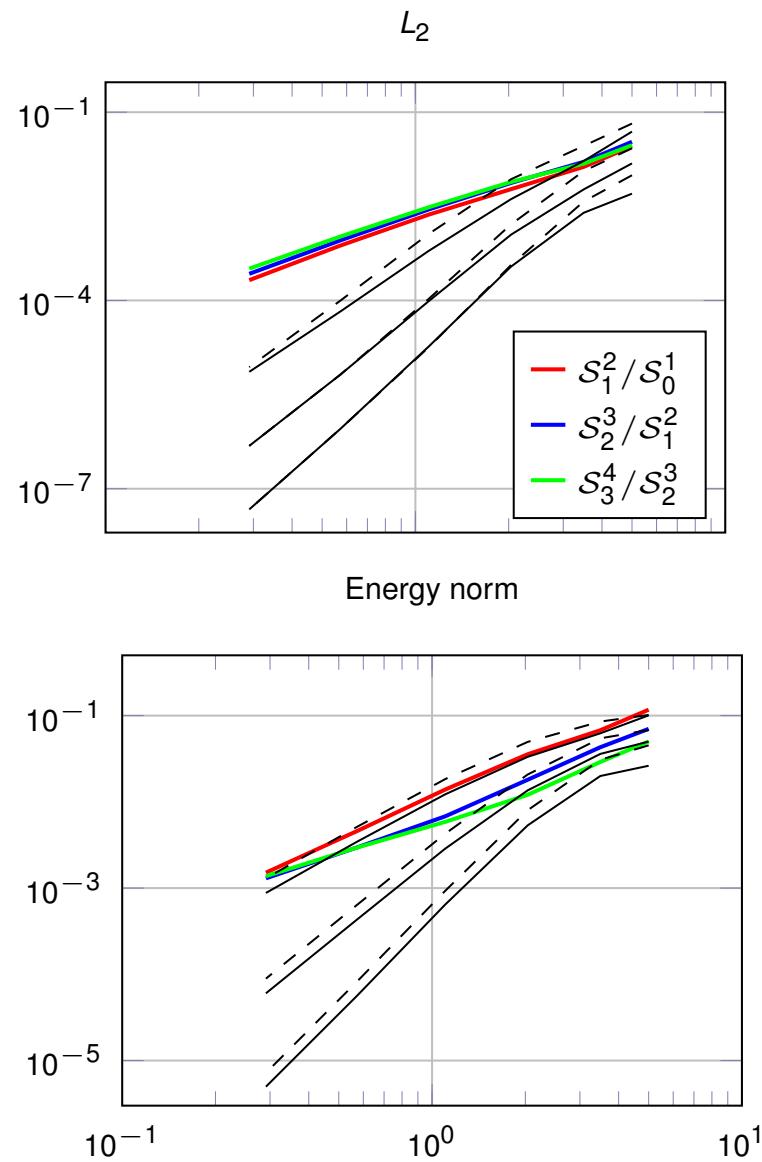
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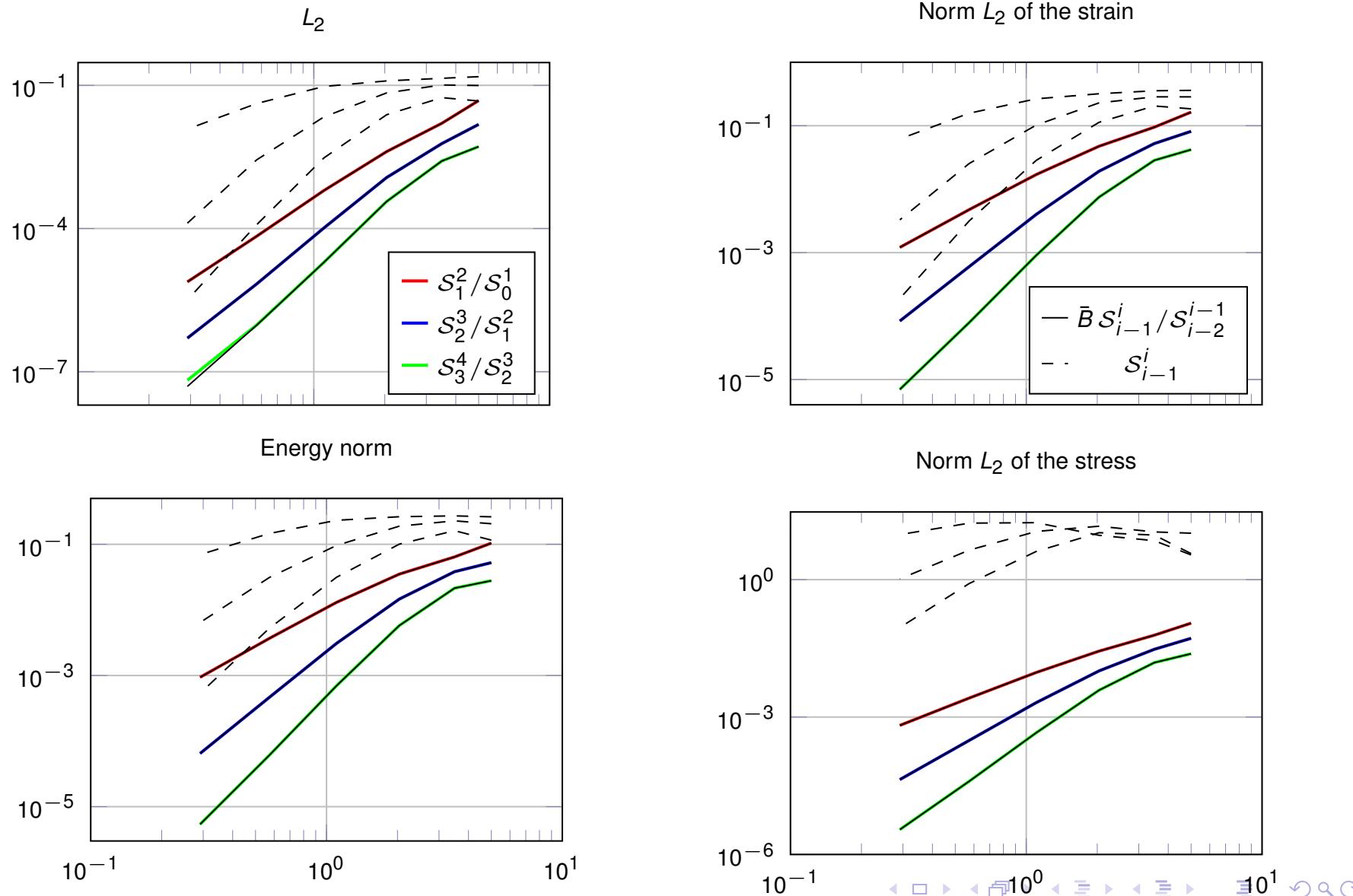
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\bar{B} with lumped mass matrix for $\nu = 0.4$



\bar{B} with lumped mass matrix for $\nu = 0.49999$



\bar{B} with lumped mass matrix and iterative solver

\bar{B} , lumped mass + iterative solver [Elguedj, Bazilevs, Calo, and Hughes, 2008]

$$a(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \pi(\nabla \cdot \mathbf{w}) \pi(\nabla \cdot \mathbf{u}) d\Omega$$

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$$a^D(\mathbf{w}, \mathbf{u}^{n+1} - \mathbf{u}^n) = L(\mathbf{w}) - a(\mathbf{w}, \mathbf{u}^n)$$

- optimally convergent ✓
- Symmetric ✓
- Sparse in \mathbf{u}^{n+1} ✓
- Definite positive ✓

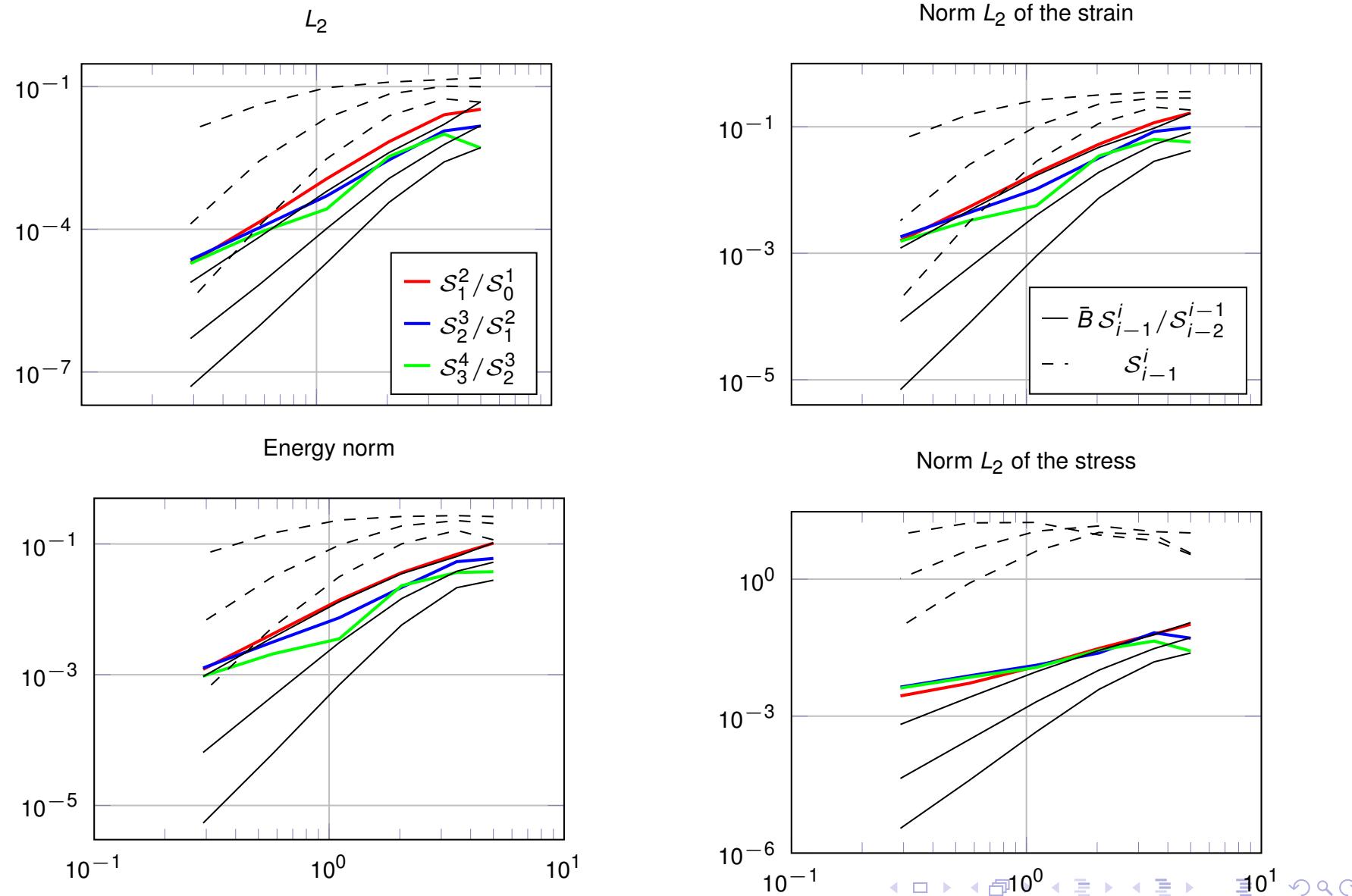
Symmetrized quasi-interpolant π^* [Lee, Lyche, and Mørken, 2001]

Symmetrized quasi-interpolant

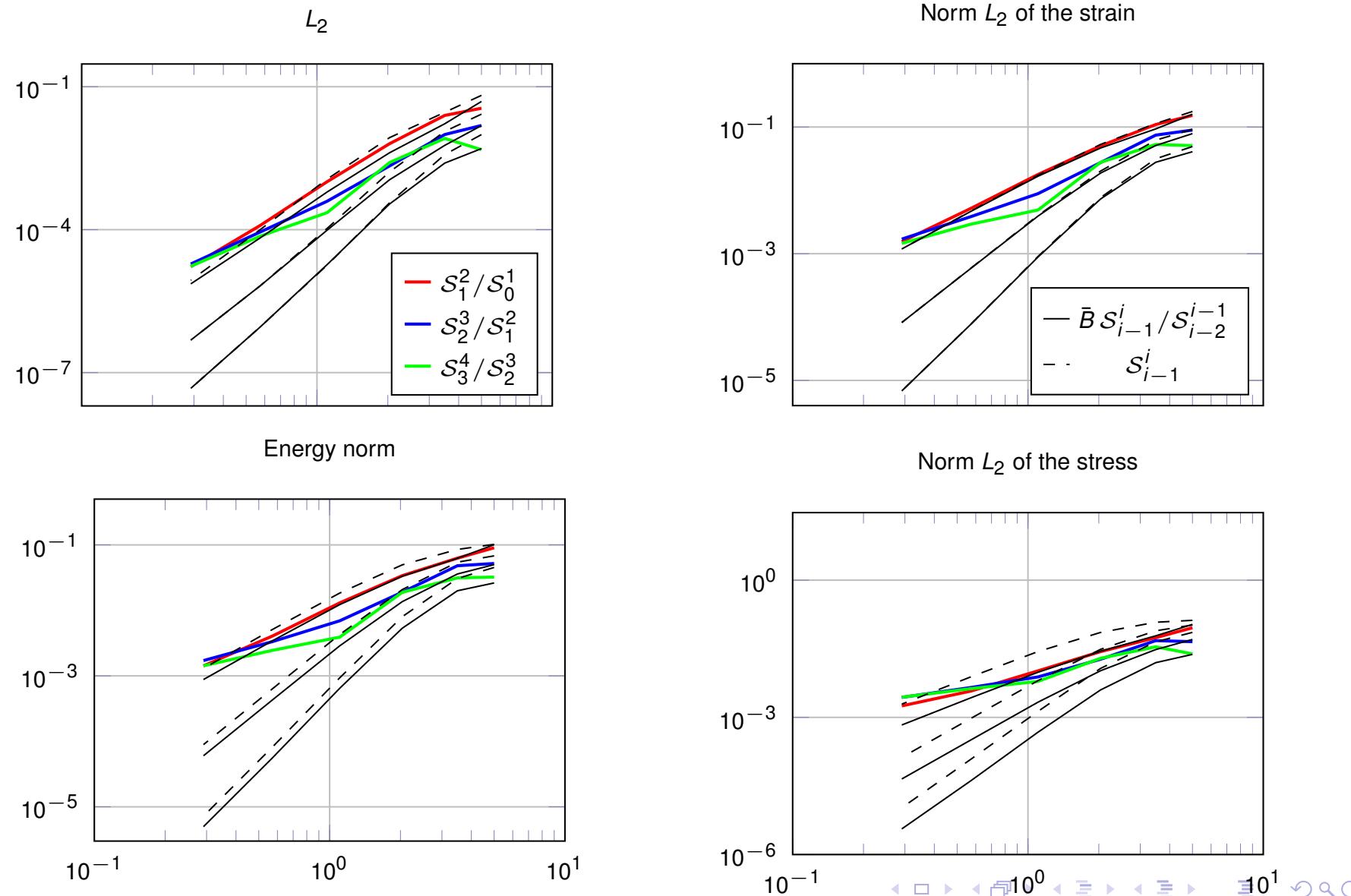
$$a(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \pi^*(\nabla \cdot \mathbf{w}) \pi^*(\nabla \cdot \mathbf{u}) d\Omega$$

- Poor approximation
- Symmetric ✓
- Sparse ✓
- Definite positive ✓

Symmetrized quasi-interpolant for $\nu = 0.49999$



Symmetrized quasi-interpolant for $\nu = 0.4$



Projection formulation as a mixed formulation

projection technique

$$\int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} + \int_{\Omega} \lambda \pi(\nabla \cdot \mathbf{w}) \pi(\nabla \cdot \mathbf{u}) = \int_{\Omega} \mathbf{w} \cdot \mathbf{f}, \quad \forall \mathbf{w}$$

Mixed formulation (for the unknowns \mathbf{u} and p)

$$\int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} + \int_{\Omega} \nabla \cdot \mathbf{w} \wp = \int_{\Omega} \mathbf{w} \cdot \mathbf{f}, \quad \forall \mathbf{w}$$

$$\int_{\Omega} (\nabla \cdot \mathbf{u}) q - \lambda^{-1} \int_{\Omega} \wp q = 0, \quad \forall q$$

where the second equation states $\lambda \pi(\nabla \cdot \mathbf{u}) = \wp$

\bar{B} method: $\mathbf{u}, \mathbf{w} \in \mathcal{S}_{p-1}^p \times \mathcal{S}_{p-1}^p \times \mathcal{S}_{p-1}^p$ and $\wp, q \in \mathcal{S}_{p-2}^{p-1}$.

Projection formulation as a mixed formulation

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Well-posedness of mixed formulation

Babuška-Brezzi inf-sup condition

The discrete displacement space $\mathbf{V}_h \subset (H_{\Gamma_D}^1(\Omega))^2$ and the discrete pressure space $Q_h \subset L^2(\Omega)$ have to fulfill

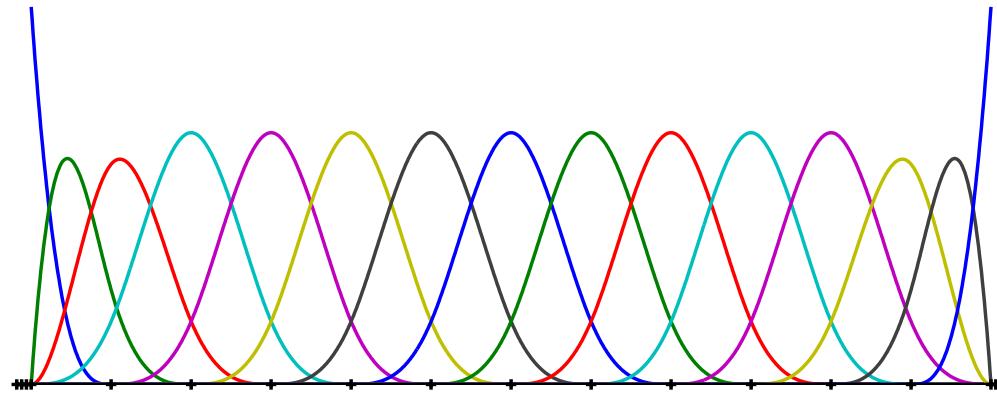
$$\inf_{q \in Q_h} \sup_{\mathbf{v} \in \mathbf{V}_h} \frac{\int_{\Omega} \nabla \cdot \mathbf{v} q \, d\Omega}{\|q\|_{L^2} \|\mathbf{v}\|_{(H^1)^2}} \geq C_{is} > 0 \quad (\text{uniformly w.r.t. } h).$$

The **inf-sup condition** above holds if “locally” and “on average” there are more knot lines of the displacement field than knot lines of the pressure field, roughly speaking... [Bressan and Sangalli, 2013] [▶ quick proof](#)

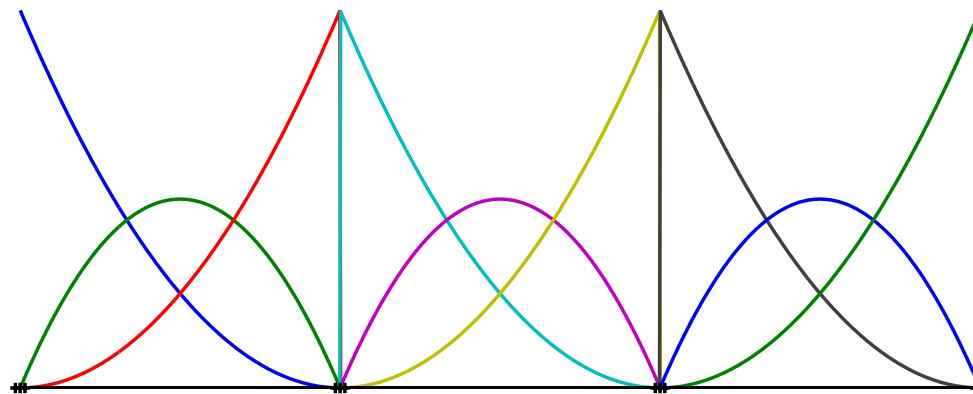
A special case for the mixed formulation

discontinuous pressure and subgrid displacement

cubic
 $\mathbf{u}, \mathbf{v} \in$



quadratic
 $\wp, q \in$



Mixed formulation with discontinuous pressures

[Antolín, Buffa, and Sangalli, 2014]

The mixed formulation with $\mathbf{u} \in (\mathcal{S}_{p-1}^p)^3$ and $\varphi \in \mathcal{S}_{-1,M}^{p-1}$

$$\int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} + \int_{\Omega} \nabla \cdot \mathbf{w} \varphi = \int_{\Omega} \mathbf{w} \cdot \mathbf{f}, \quad \forall \mathbf{w}$$
$$\int_{\Omega} (\nabla \cdot \mathbf{u}) q - \lambda^{-1} \int_{\Omega} \varphi q = 0, \quad \forall q$$

- Well-posed
- optimally convergent
- the 2nd equation is local to macroelements! $\lambda \pi_M(\nabla \cdot \mathbf{u}) = \varphi$

Projection formulation with the local π_M

Mixed formulation with $\mathbf{u} \in (\mathcal{S}_{p-1}^p)^3$ and $\wp \in \mathcal{S}_{-1,M}^{p-1}$

$$\int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} + \int_{\Omega} \nabla \cdot \mathbf{w} \wp = \int_{\Omega} \mathbf{w} \cdot \mathbf{f}, \quad \forall \mathbf{w}$$
$$\int_{\Omega} (\nabla \cdot \mathbf{u}) q = \lambda^{-1} \int_{\Omega} \wp q, \quad \forall q$$



projection technique with macroelement projection π_M

$$\int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} + \int_{\Omega} \lambda \pi_M(\nabla \cdot \mathbf{w}) \pi_M(\nabla \cdot \mathbf{u}) = \int_{\Omega} \mathbf{w} \cdot \mathbf{f}, \quad \forall \mathbf{w}$$

Macroelement projection π_M

Projection technique with macroelement projection π_M

$$a(\mathbf{w}, \mathbf{u}) = \int_{\Omega} \mu \nabla^s \mathbf{w} : \nabla^s \mathbf{u} d\Omega + \int_{\Omega} \lambda \pi_M(\nabla \cdot \mathbf{w}) \pi_M(\nabla \cdot \mathbf{u}) d\Omega$$

$$\tilde{N} \in \mathcal{S}_{-1,M}^{p-1} \quad \pi_M(\phi)(\mathbf{x}) = \sum_i \tilde{N}_i(\mathbf{x}) \left[\sum_j \bar{\mathbf{M}}_{ij}^{-1} \int_{\Omega} \tilde{N}_j(\mathbf{x}) \phi(\mathbf{x}) d\Omega \right]$$

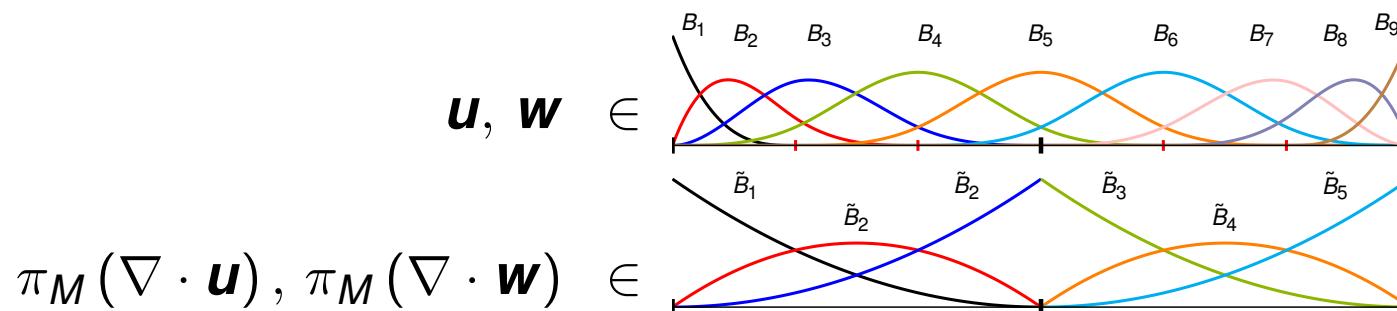
It works with minimal macroelements of p element per direction, but the theory is missing (unless $p = 1$: Pitkäranta and Stenberg [1984])

Macroelement projection π_M

Projection technique with macroelement projection π_M

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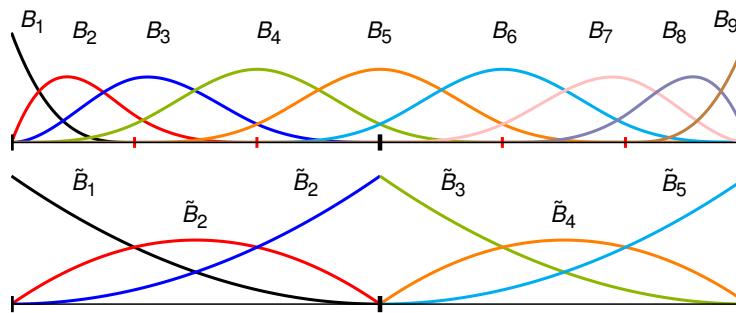
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$$M^{-1} = \begin{pmatrix} p^3 & & \\ & p^3 & \\ & & p^3 \\ p^3 & & \\ & p^3 & \\ & & p^3 \end{pmatrix}^{-1} = \begin{pmatrix} p^3 & & \\ & p^3 & \\ & & p^3 \\ p^3 & & \\ & p^3 & \\ & & p^3 \end{pmatrix}$$



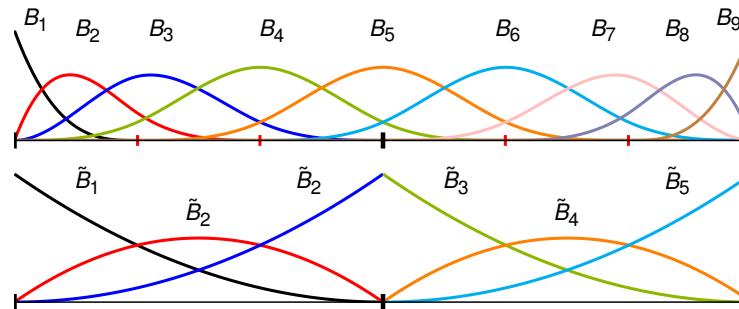
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$$\tilde{N} \in \mathcal{S}_{-1,M}^{p-1} \quad \pi_M(\phi)(\mathbf{x}) = \sum_i \tilde{N}_i(\mathbf{x}) \left[\sum_j \textcolor{blue}{M}_{ij}^{-1} \int_{\Omega} \tilde{N}_j(\mathbf{x}) \phi(\mathbf{x}) d\Omega \right]$$

$$M^{-1} = \left(\begin{array}{cccccc} \textcolor{blue}{\blacksquare} & & & & & \\ & \textcolor{blue}{\blacksquare} & & & & \\ & & \textcolor{blue}{\blacksquare} & & & \\ & & & \textcolor{blue}{\blacksquare} & & \\ & & & & \textcolor{blue}{\blacksquare} & \\ & & & & & \textcolor{blue}{\blacksquare} \end{array} \right)^{-1} = \left(\begin{array}{cccccc} \textcolor{red}{\blacksquare} & & & & & \\ & \textcolor{red}{\blacksquare} & & & & \\ & & \textcolor{red}{\blacksquare} & & & \\ & & & \textcolor{red}{\blacksquare} & & \\ & & & & \textcolor{red}{\blacksquare} & \\ & & & & & \textcolor{red}{\blacksquare} \end{array} \right)$$

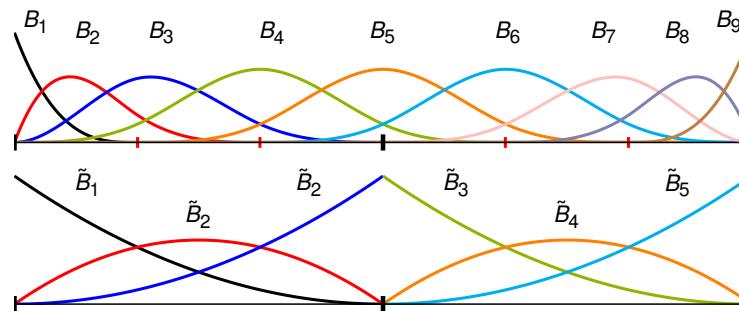


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Macroelement projection π_M

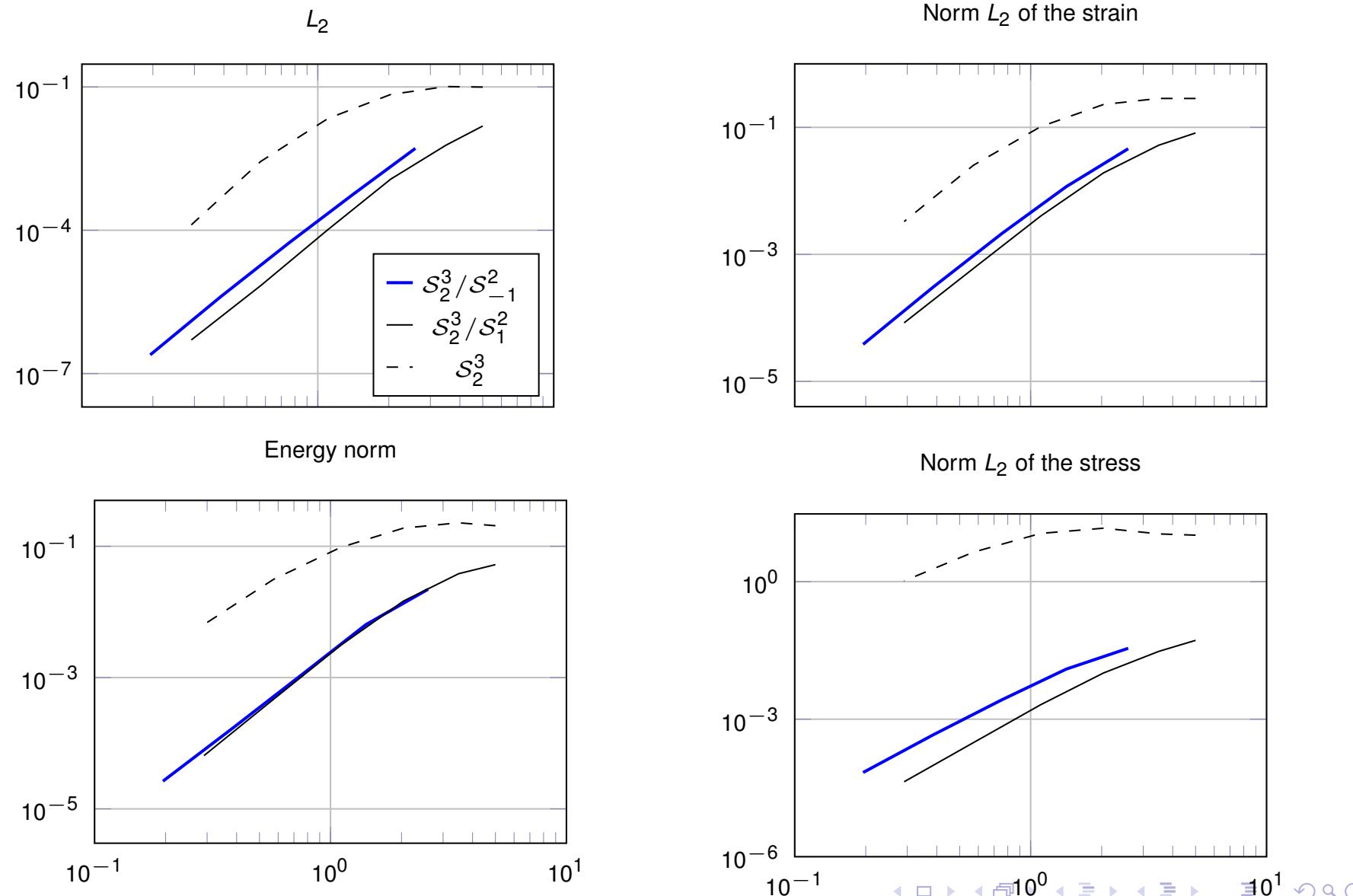
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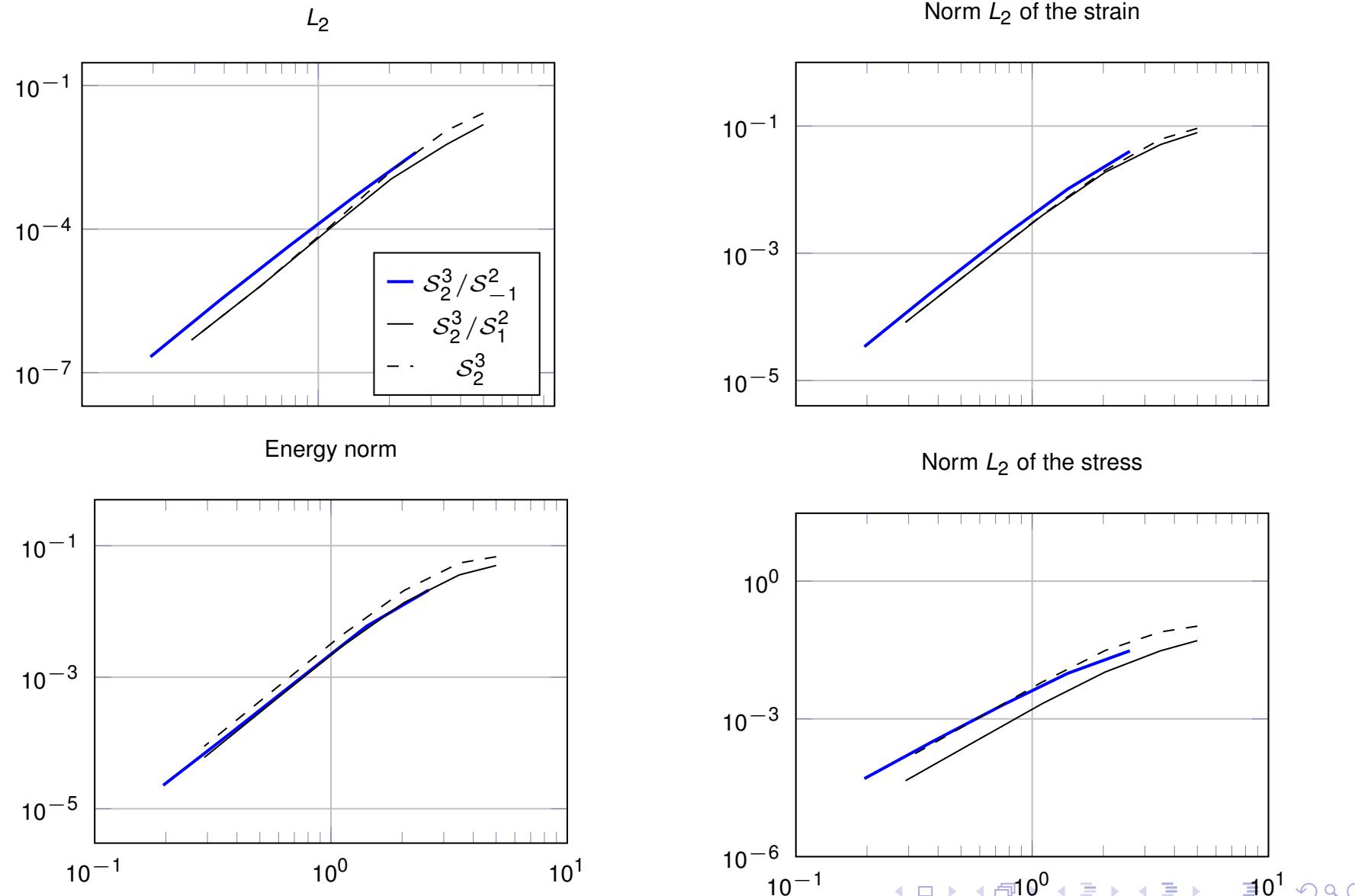
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- Unlocked solution ✓
- Symmetric ✓
- Sparse ✓
- Definite positive ✓

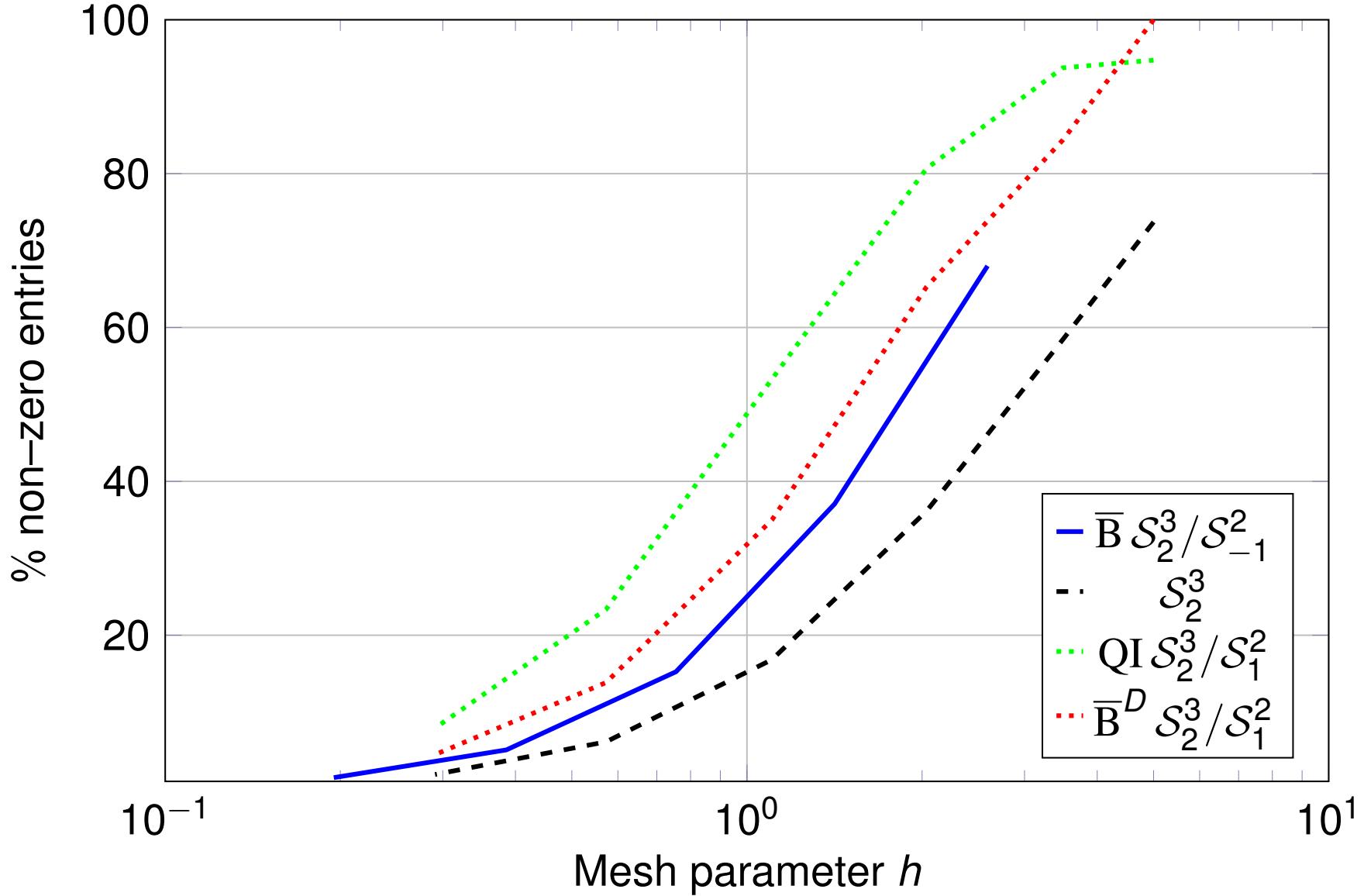
Macroelement projection π_M for $\nu = 0.49999$



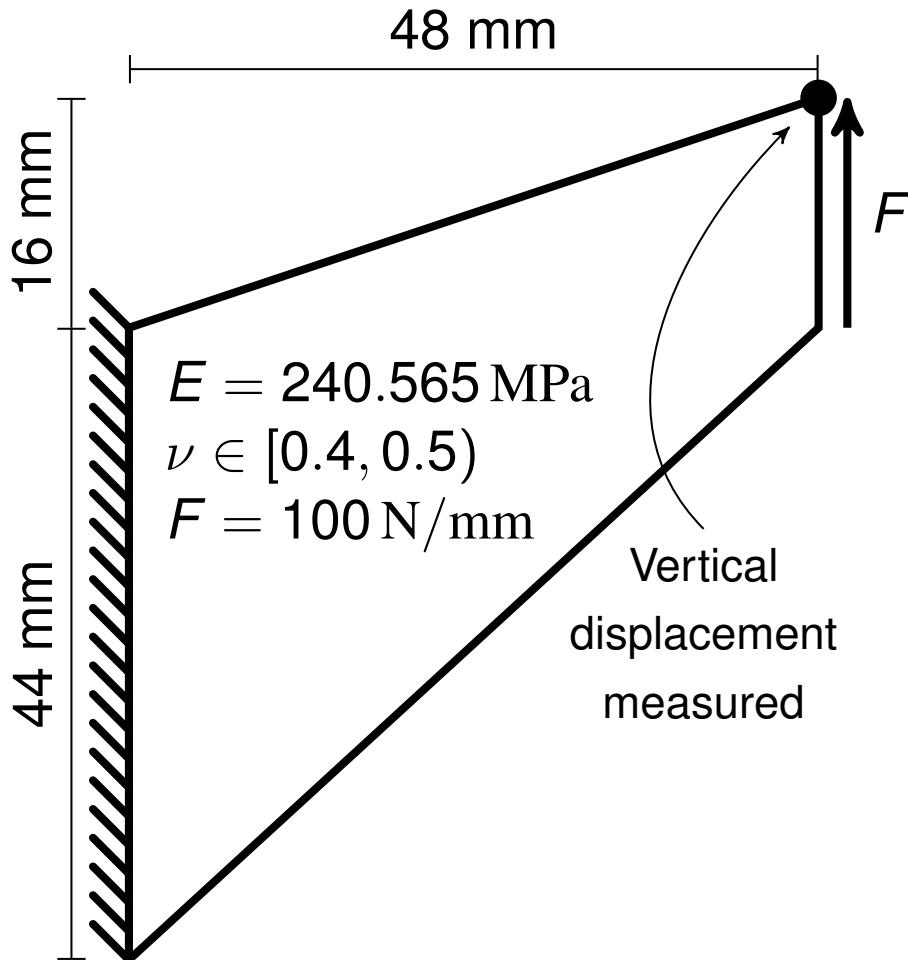
Macroelement projection π_M for $\nu = 0.4$



Macroelement projection π_M : sparsity

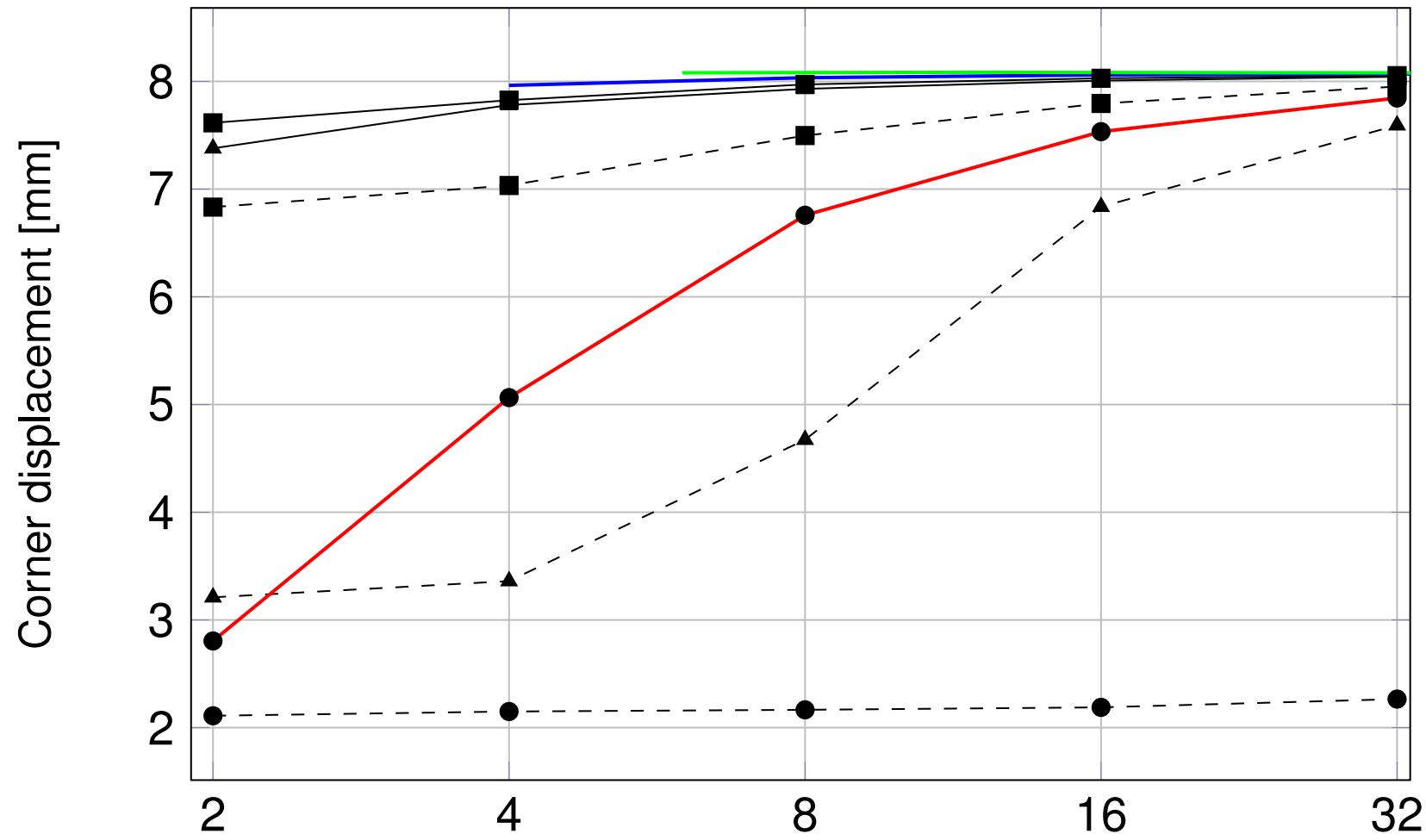


Cook Membrane



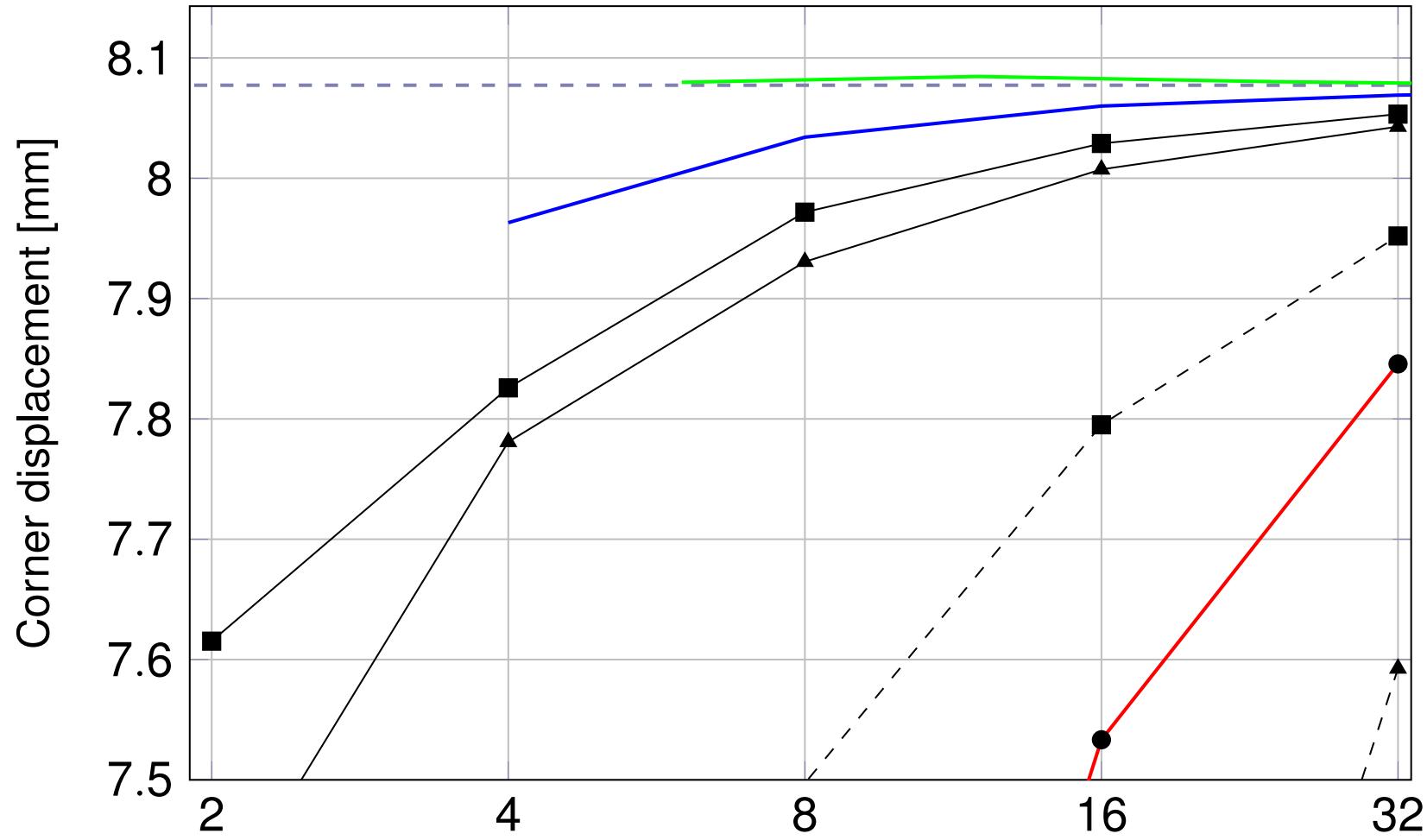
Cook Membrane $\nu = 0.49999$

- ● - \mathcal{S}_0^1 - ▲ - \mathcal{S}_1^2 - ■ - \mathcal{S}_2^3 - ● - $\bar{\mathcal{B}}\mathcal{S}_0^1/\mathcal{S}_{-1}^0$ - ▲ - $\bar{\mathcal{B}}\mathcal{S}_1^2/\mathcal{S}_0^1$
- ■ - $\bar{\mathcal{B}}\mathcal{S}_2^3/\mathcal{S}_1^2$ - — $\bar{\mathcal{B}}\mathcal{S}_0^1/\mathcal{S}_{-1}^0$ - — $\bar{\mathcal{B}}\mathcal{S}_1^2/\mathcal{S}_{-1}^1$ - — $\bar{\mathcal{B}}\mathcal{S}_2^3/\mathcal{S}_{-1}^2$



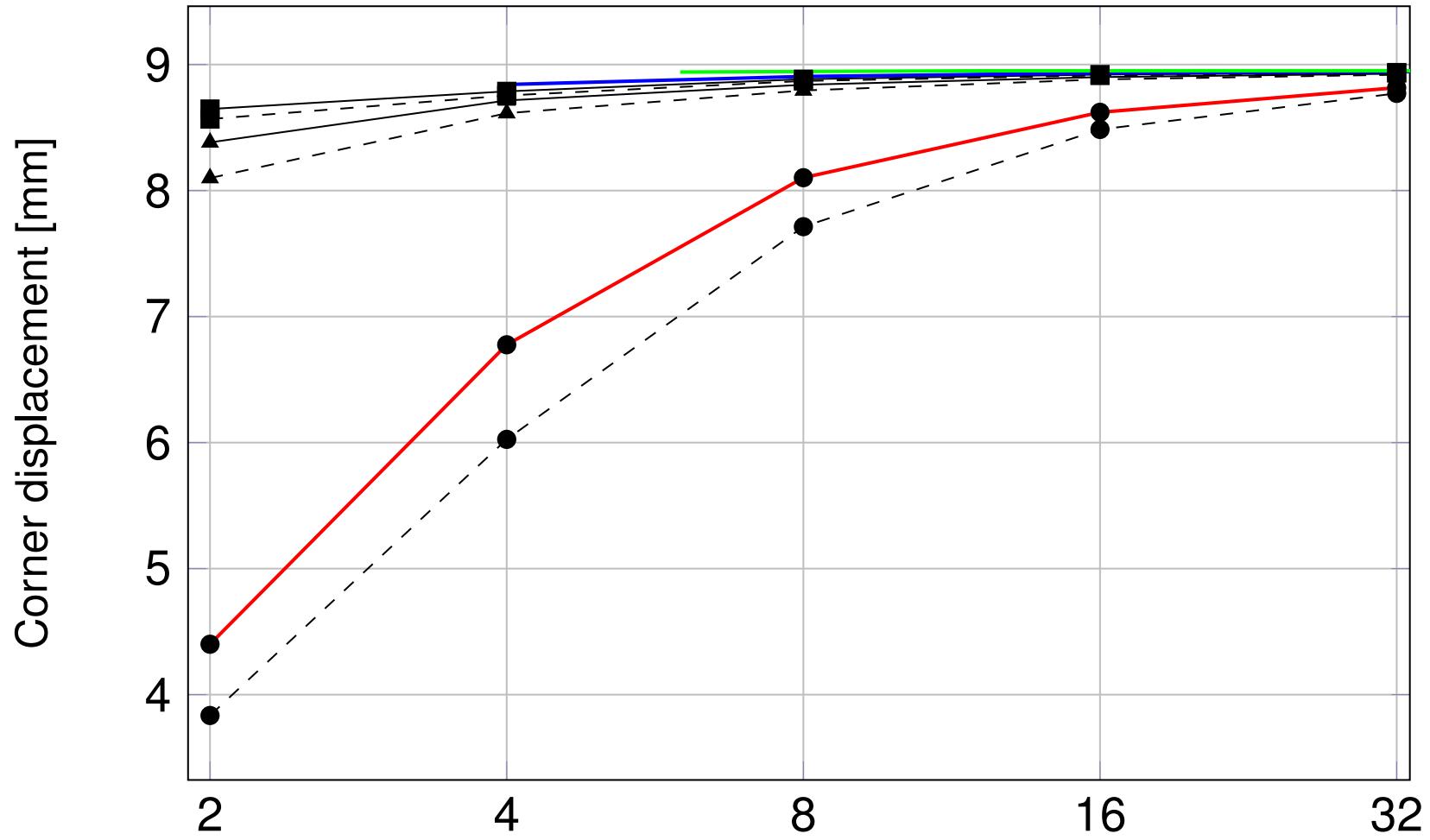
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- ■ - $\bar{\mathcal{B}} \mathcal{S}_2^3 / \mathcal{S}_1^2$ - — $\bar{\mathcal{B}} \mathcal{S}_0^1 / \mathcal{S}_{-1}^0$ - — $\bar{\mathcal{B}} \mathcal{S}_1^2 / \mathcal{S}_{-1}^1$ - — $\bar{\mathcal{B}} \mathcal{S}_2^3 / \mathcal{S}_{-1}^2$



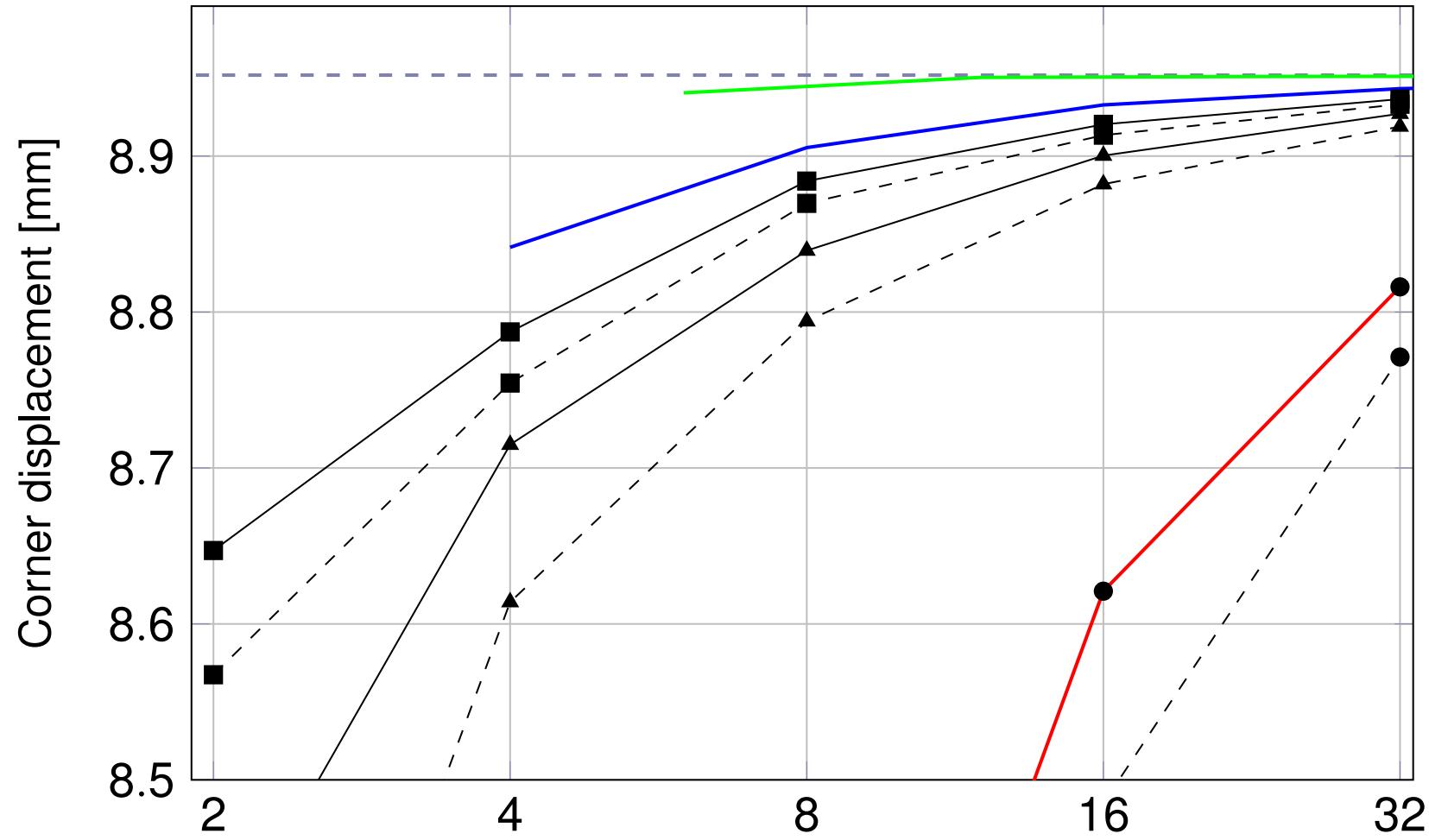
Cook Membrane $\nu = 0.4$

- ● - \mathcal{S}_0^1 - ▲ - \mathcal{S}_1^2 - ■ - \mathcal{S}_2^3 - ● - $\bar{\mathcal{B}} \mathcal{S}_0^1 / \mathcal{S}_{-1}^0$ - ▲ - $\bar{\mathcal{B}} \mathcal{S}_1^2 / \mathcal{S}_0^1$
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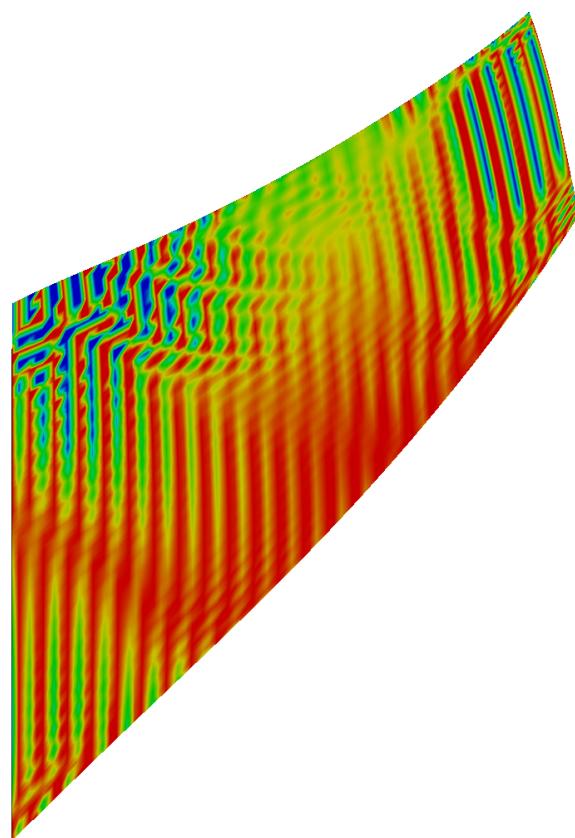
Cook Membrane $\nu = 0.4$

- ● - \mathcal{S}_0^1 - ▲ - \mathcal{S}_1^2 - ■ - \mathcal{S}_2^3 - ● - $\bar{\mathcal{B}} \mathcal{S}_0^1 / \mathcal{S}_{-1}^0$ - ▲ - $\bar{\mathcal{B}} \mathcal{S}_1^2 / \mathcal{S}_0^1$
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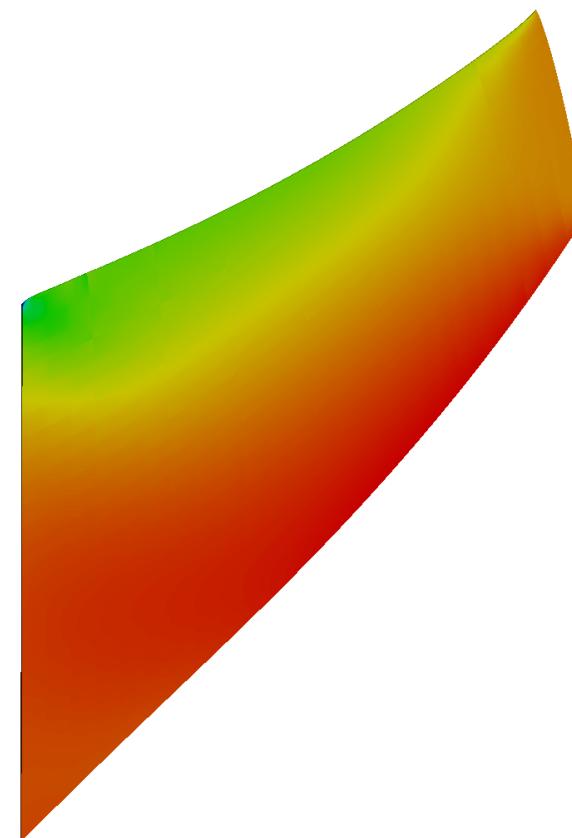


Cook Membrane σ_{xx} for $\nu = 0.49999$ $p = 3$

Standard formulation

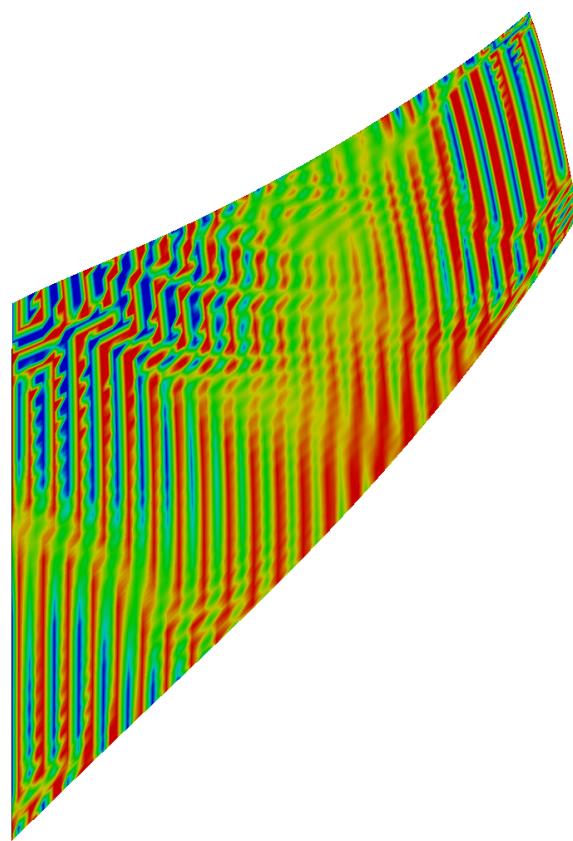


Discontinuous subgrid \bar{B}

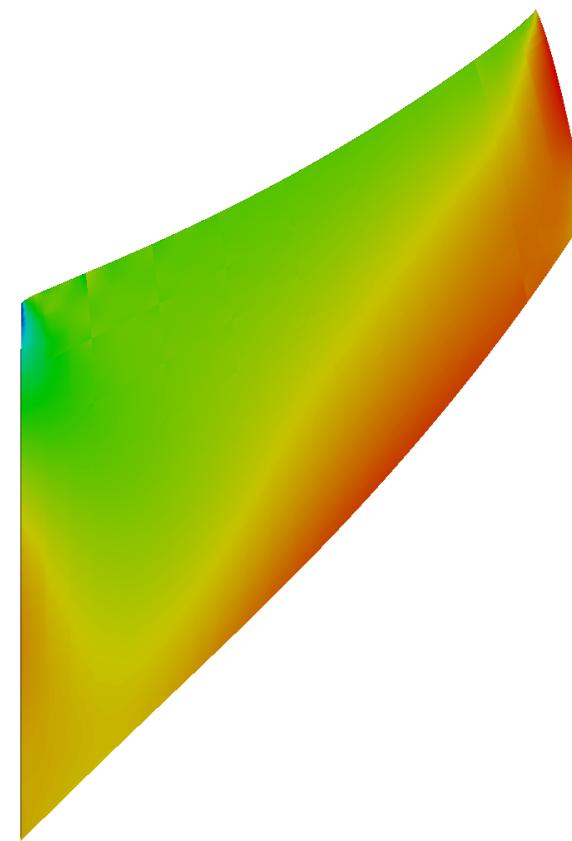


Cook Membrane σ_{yy} for $\nu = 0.49999$ $p = 3$

Standard formulation



Discontinuous subgrid \bar{B}



Conclusion

Based on the \bar{B} -method, we have proposed a new formulation with:

- symmetric and positive definite stiffness matrix
- sparse, with little increase of non-zeros w.r.t. plain Galerkin
- locking-free and optimally accurate

It uses smooth & discontinuous functions together :-)

Open ERC-funded positions:

<http://www-dimat.unipv.it/sangalli/higeom>

*** Thank you for your attention ***

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References

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Macroelement method [Bazilevs, Beirão da Veiga, Cottrell, Hughes, and Sangalli, 2006; Bressan and Sangalli, 2013]

Let $\mathcal{M} = \{M\}$ a covering of $\hat{\Omega}$ such that the (local) full-rank condition

$\forall \hat{q}_h \in \hat{Q}_{h,0}, \hat{q}_h$ non-constant on $M, \exists \hat{\mathbf{v}}_h \in \hat{V}_{h,0}$ with $\text{supp}(\hat{\mathbf{v}}_h) \subseteq M$

such that $\int_M \hat{q}_h \text{div} \hat{\mathbf{v}}_h \neq 0$

holds in each macroelement $M \in \mathcal{M}$ (+ some technicalities)

Macroelement method

[Bazilevs, Beirão da Veiga, Cottrell, Hughes, and

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holds in each macroelement $M \in \mathcal{M}$ (+ some technicalities)

then the global inf-sup condition hold:

$$\forall \hat{q}_h \in \hat{Q}_{h,0}, \exists \hat{\mathbf{v}}_h \in \hat{V}_{h,0} \text{ such that } \frac{\int_{\hat{\Omega}} \hat{q}_h \operatorname{div} \hat{\mathbf{v}}_h}{\|\hat{q}_h\|_{L^2} \|\hat{\mathbf{v}}_h\|_{(H^1)^2}} \geq \hat{C}_{is} > 0$$

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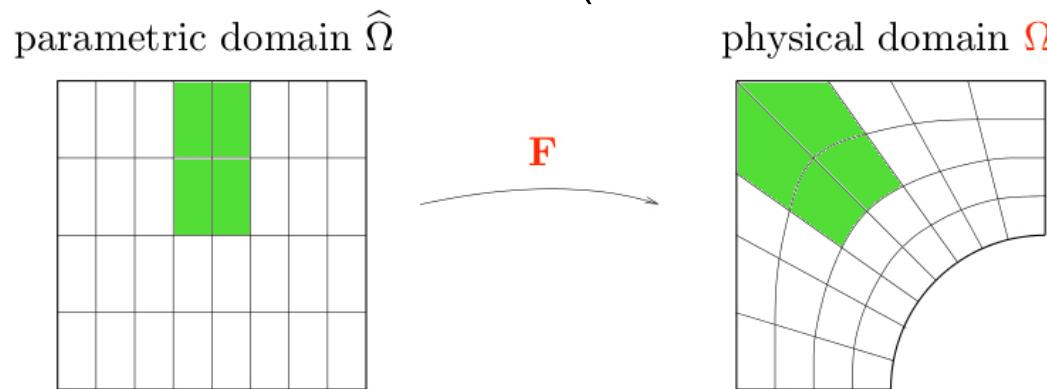
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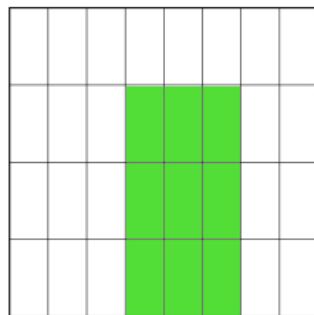
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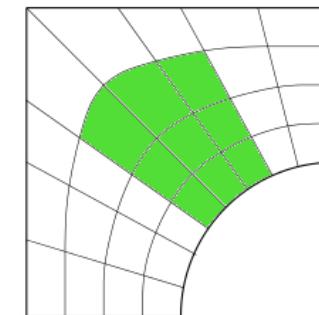
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holds in each macroelement $M \in \mathcal{M}$ (+ some technicalities)

parametric domain $\hat{\Omega}$



physical domain Ω



\mathbf{F}

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$$\begin{array}{c} \uparrow \\ \forall g \neq 0 \in S = S_1 \otimes S_2, \exists f \in T = T_1 \otimes T_2 : \int_M f(x_1, x_2) g(x_1, x_2) dx_1 dx_2 > 0 \\ \Downarrow \end{array}$$

$$\forall g \neq 0 \in S_i, \exists f \in T_i : \int f(x) g(x) dx > 0$$

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⇓

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Full-rank characterization of 1D scalar products

Let S a spline space on the partition ζ_0, \dots, ζ_m , let $J = [a, b[$ be a subinterval of $[\zeta_0, \zeta_m[$ and define:

$$S_{\subset J} = \text{span}\{B_i \in S : B_i(x) = 0, \forall x \in \setminus J\}, \quad (1)$$

$$S_{\cap J} = \text{span}\{B_i \in S : B_i|_J \neq 0\}, \quad (2)$$

Theorem [Bressan and Sangalli, 2013]

Let T be a spline space on the partition η_0, \dots, η_n and S be a spline space on ζ_0, \dots, ζ_m with $[\zeta_0, \zeta_m] = [\eta_0, \eta_n]$; the following two properties are equivalent:

for all $g \in S$, $g \neq 0$, there exists $f \in T$ such that $\int_{\mathbb{R}} f(x)g(x) dx \neq 0$;

for all $0 \leq i < j \leq n$, $\dim T_{\cap[\eta_i, \eta_j[} \geq \dim S_{\subset[\eta_i, \eta_j[}$.

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