



## EXAMINATION PAPER

<b>Examination Session:</b> May/June	<b>Year:</b> 2022	<b>Exam Code:</b> MATH2071-WE01
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<b>Title:</b> Mathematical Physics II
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Time:	3 hours	
Additional Material provided:		
Materials Permitted:		
Calculators Permitted:	No	Models Permitted: Use of electronic calculators is forbidden.

Instructions to Candidates:	<p>Answer all questions.</p> <p>Section A is worth 40% and Section B is worth 60%. Within each section, all questions carry equal marks.</p> <p>Students must use the mathematics specific answer book.</p>
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<b>Revision:</b>	
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## SECTION A

**Q1** Consider a Lagrangian of the form

$$L = \frac{1}{2}(\dot{q}_1^2 + \dot{q}_2^2) + \cos(q_1 + q_2).$$

- 1.1** Write down the equations of motion for the system. You do not need to solve them.
- 1.2** Show that  $q_1(t) = q_2(t) = 0$  is a solution of the equations of motion.
- 1.3** Find an approximate Lagrangian  $L_{\text{app}}$  describing small perturbations around  $q_1(t) = q_2(t) = 0$ .
- 1.4** Find the general solution of the equations of motion derived from  $L_{\text{app}}$ .

**Q2** **2.1** The expression for the energy of a system with Lagrangian  $L$  is given by

$$E = \left( \sum_{i=1}^n \dot{q}_i \frac{\partial L}{\partial \dot{q}_i} \right) - L.$$

Show that if  $L$  does not depend explicitly on time, then  $E$  is conserved along solutions of the equations of motion.

**2.2** Consider a theory with a Lagrangian of the form

$$L = \frac{1}{2}\dot{q}_1^2 + \frac{1}{4}q_1^3\dot{q}_2^4 - f(q_1, q_2),$$

with  $f$  an arbitrary function of two arguments. Construct the canonical momenta  $p_1$  and  $p_2$  associated to  $q_1$  and  $q_2$ .

**2.3** Find the Hamiltonian for the system.

**Q3** A quantum mechanical system is, at  $t = 0$ , prepared in a state described by the wave function

$$\psi(t = 0, x) = C \left( \frac{1}{\sqrt{2}} \psi_{E=1}(x) + e^{i\beta} \psi_{E=2}(x) \right),$$

where  $C$  and  $\beta$  are real constants. The functions  $\psi_{E=1}$  and  $\psi_{E=2}$  are normalised energy-eigenfunctions of the system, with eigenvalues as indicated.

- 3.1** Determine the constant  $C$ . Is it possible to observe the overall phase factor of this constant? Motivate your answer.
- 3.2** An energy measurement is made. What are the possible outcomes, and what are the probabilities of those outcomes?
- 3.3** Is it possible to observe the value of the phase  $\beta$ ? Motivate your answer.

- Q4** Consider a potential for a one-dimensional quantum particle, corresponding to a uniform constant force,

$$V(x) = -Fx,$$

where  $F$  is a real constant.

- 4.1** Write down the time-independent Schrödinger equation *in momentum space* for the wave function  $\tilde{\psi}(p)$ , including the potential given above.
- 4.2** Find the solution  $\tilde{\psi}(p)$  up to a normalisation constant.

## SECTION B

- Q5** Consider a system where the Lagrangian  $L(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}})$  depends on the positions  $\mathbf{q} = (q_1, \dots, q_n)$ , velocities  $\dot{\mathbf{q}} = (\dot{q}_1, \dots, \dot{q}_n)$  and accelerations  $\ddot{\mathbf{q}} = (\ddot{q}_1, \dots, \ddot{q}_n)$ . In this case the  $n$  Euler-Lagrange equations of motion are

$$\frac{\partial L}{\partial q_i} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_i} \right) + \frac{d^2}{dt^2} \left( \frac{\partial L}{\partial \ddot{q}_i} \right) = 0$$

for  $i \in \{1, \dots, n\}$ .

We assume that the transformation  $q_i \rightarrow q_i + \epsilon a_i(\mathbf{q})$  (with  $n$  independent generators  $a_1, \dots, a_n$ ) leaves the Lagrangian  $L(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}})$  invariant to first order in  $\epsilon$ .

- 5.1** Assume first that  $L$  does not depend on the accelerations  $\ddot{q}_i$ . Show that in this case the Noether charge

$$Q = \sum_{i=1}^n a_i \frac{\partial L}{\partial \dot{q}_i}$$

is conserved.

- 5.2** Coming back to the more general case where  $L(\mathbf{q}, \dot{\mathbf{q}}, \ddot{\mathbf{q}})$  does depend on the accelerations  $\ddot{q}_i$ , find a conserved charge  $Q$  associated to the transformation generated by the  $a_i$ . [*Hint*: Try to write the variation of  $L$  to first order in  $\epsilon$  as a total time derivative.]

- 5.3** As an example, find the explicit form for  $Q$  when the Lagrangian is

$$L_2 := -\frac{\alpha}{2}(\ddot{q}_1^2 + \ddot{q}_2^2) + \frac{\beta}{2}(\dot{q}_1^2 + \dot{q}_2^2) - \frac{\gamma}{2}(q_1^2 + q_2^2)$$

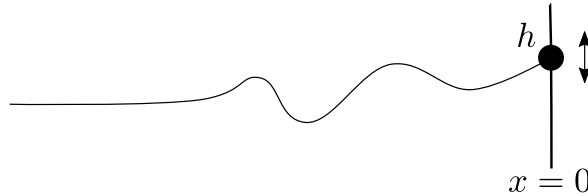
and the transformation is a rotation around the origin in the  $(q_1, q_2)$  plane.

- 5.4** Verify by taking the time derivative that the explicit  $Q$  that you just found is conserved along solutions of the equations of motion for the Lagrangian  $L_2$ .

**Q6** We describe a string oscillating in one dimension by a field  $u(x, t)$  with Lagrangian density

$$\mathcal{L} = \frac{1}{2}u_t^2 - \frac{1}{2}u_x^2 - \frac{1}{2}mu^2$$

with  $u_t := \frac{\partial u}{\partial t}$  and  $u_x := \frac{\partial u}{\partial x}$ . The string extends from  $x = -\infty$  to  $x = 0$ , where it ends on a bead of mass  $h$  that is constrained to move on the vertical line  $x = 0$ . The bead can slide without friction along  $x = 0$ , and we ignore the effect of gravity.



**6.1** Find the equation of motion for  $u$  valid in the region  $x < 0$ .

**6.2** Using that the energy-momentum tensor is given by

$$T_{ij} = \frac{\partial \mathcal{L}}{\partial u_j} \frac{\partial u}{\partial x_i} - \delta_{ij} \mathcal{L},$$

where  $u_j := \partial u / \partial x_j$ , compute the energy flux  $T_{tx}$  associated to  $\mathcal{L}$ .

**6.3** Consider an ansatz given by

$$u(x, t) = \text{Re} \left( e^{i\omega t} (e^{-ikx} + \rho e^{ikx}) \right),$$

where  $\text{Re}(z)$  indicates taking the real part of  $z$ , and  $\rho$  is a complex number, which for generic  $h$  we assume to be different from  $-1$ . Find the values of  $\omega$  and  $\rho$  (as functions of  $k$ ,  $m$  and  $h$ ) that make this ansatz a solution of the problem. [*Hint*: Impose energy conservation at the boundary.]

**6.4** Consider the  $(m, h) = (0, 0)$  case. Which standard boundary condition for the massless scalar does this correspond to? Similarly, which standard boundary condition do you obtain in the  $(m, h) \rightarrow (0, \infty)$  limit? Show in both cases that the form of  $u(x, t)$  is the expected one.

**Q7** Consider a system of two one-dimensional, distinguishable particles in a simple harmonic oscillator potential of frequency  $\omega$ . The coordinates of the two particles are  $x_1$  and  $x_2$ .

**7.1** Give the energy eigenfunctions for the two-particle system in terms of the energy eigenfunctions of the one-particle simple harmonic oscillator. Show that the energy eigenvalues for the two-particle wave functions take the form  $E = (n_1 + n_2 + 1)\hbar\omega$  (with  $n_1$  and  $n_2$  integers).

**7.2** Now assume that the particles are *indistinguishable*. This means that all probabilities have to remain unchanged under an exchange  $x_1 \leftrightarrow x_2$ , so

$$\left| \psi(x_1, x_2, t) \right|^2 = \left| \psi(x_2, x_1, t) \right|^2.$$

Show that there are now only *two* independent wave functions  $\psi(x_1, x_2, t)$  for a state with energy  $E = 2\hbar\omega$  (up to an irrelevant phase factor).

**7.3** The explicit form of the first two wave functions for the single particle system are given by

$$\psi_{n=0}(x) = C \exp\left(-\frac{m\omega x^2}{2\hbar}\right), \quad \psi_{n=1}(x) = \sqrt{\frac{2m\omega}{\hbar}} x \times \psi_{n=0}(x).$$

Rewrite the two wave functions found in **7.2** in terms of the centre of mass  $X$  and separation  $y$ ,

$$X = \frac{1}{2}(x_1 + x_2), \quad y = x_1 - x_2.$$

**7.4** Determine the probability density for the separation  $P(y)$  by averaging over  $X$ , separately for the two wave functions found in **7.3**. You do *not* have to work out the overall normalisation factors.

**7.5** Sketch the form of  $P(y)$  for the two cases. Which of the two describes a system in which particles repel each other?

**Q8** Consider the time-dependent Schrödinger equation for a single free particle. A Galilean transformation is a transformation on  $x$  and  $p$  which takes us to the coordinates of a moving observer, so in particular  $x \rightarrow x' = x + vt$  (and  $t = t'$ ).

- 8.1** Write down the Schrödinger equation (and include a generic potential  $V(x, t)$ ). Transform coordinates to  $x', t'$ , assuming that the wave function in primed coordinates is related to the one in unprimed coordinates by a phase factor,

$$\psi'(x', t') = \exp \left[ i f(x', t') \right] \psi(x, t),$$

where  $f(x', t')$  is real. Use that  $V'(x', t') = V(x, t)$ . [*Hint*: First show that

$$\frac{\partial}{\partial x} = \frac{\partial}{\partial x'}, \quad \frac{\partial}{\partial t} = \frac{\partial}{\partial t'} + v \frac{\partial}{\partial x'}.$$

and then apply this to the Schrödinger equation. ]

- 8.2** Show that the Schrödinger equation in primed coordinates has the same form as the one in unprimed coordinates when

$$f(x', t') = \frac{mvx' - \frac{1}{2}mv^2t'}{\hbar}.$$

- 8.3** Consider now the situation  $V(x) = 0$ . The momentum eigenfunction for a single particle reads

$$\psi(x, t) = e^{\frac{i}{\hbar}px - i\omega t}, \quad \omega = \frac{p^2}{2\hbar m}.$$

Transform it to  $\psi'(x', t')$  and show that the result is consistent with the transformation which you expect  $x' = x + vt$  to induce on the momentum.