

On the gaps in spectrum of the Maxwell Operator: case of Photonic Crystal Fibres

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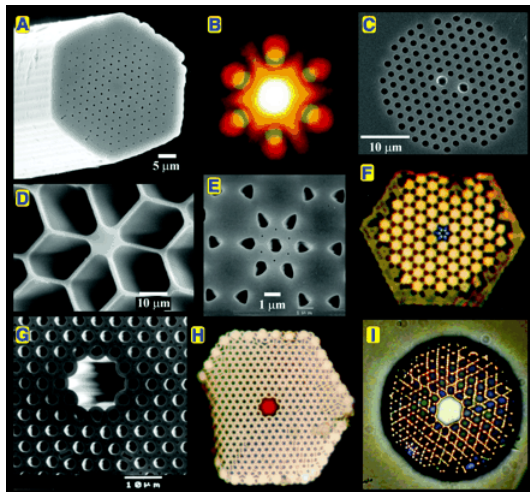
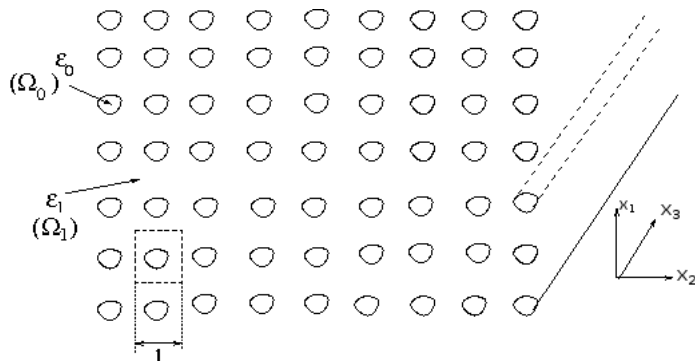


Figure: Taken from “Photonic Crystal Fibres” Phillip Russell, Science, 2003



$$\nabla \times E = -\mu \frac{\partial H}{\partial t}, \quad \nabla \times H = \epsilon \frac{\partial E}{\partial t},$$

$$\nabla \cdot (\epsilon E) = 0, \quad \nabla \cdot H = 0,$$

$$\epsilon = \epsilon_0 \chi_0(x) + \epsilon_1 \chi_1(x), \quad \epsilon_0 \neq \epsilon_1, \quad \mu \text{ constant} \quad (\mu = 1)$$

$$E = E(x_1, x_2) \exp(i(kx_3 + \omega t)), \quad H = H(x_1, x_2) \exp(i(kx_3 + \omega t))$$

If $k = 0$ then

$$\begin{aligned} -\Delta E_3(x) &= \epsilon(x)\omega^2 E_3(x), \quad x \in \mathbb{R}^2, \\ -\nabla \epsilon(x)^{-1} \nabla H_3(x) &= \omega^2 H_3(x), \quad x \in \mathbb{R}^2. \end{aligned}$$

We have spectral problem for

$$A[E_3, H_3] = \int_{\mathbb{R}^2} |\nabla E_3|^2 + \epsilon(x)^{-1} |\nabla H_3|^2 dx$$

and

$$B[E_3, H_3] = \int_{\mathbb{R}^2} \epsilon(x) |E_3|^2 + |H_3|^2 dx.$$

Gaps due to high-contrast

A.Figotin, P. Kuchment 1993, R.Hempel, K Lienau 2000, V. Zhikov, 2000 and many others.

Spectral problem for

$$a[u] = \int_{\Omega_0} |\nabla u|^2 dx + t^2 \int_{\Omega_1} |\nabla u|^2 dx,$$

and

$$b[u] = \int_{\mathbb{R}^2} |u|^2 dx.$$

Gaps appear as $t \rightarrow \infty$. High-contrast.

No high contrast in $\epsilon(x), \mu(x)$ for Photonic Crystal Fibers.

Gaps are not expected unless $\epsilon_1 \ll \epsilon_2$ or $\epsilon_2 \ll \epsilon_1$.

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No high contrast in $\epsilon(x), \mu(x)$ for Photonic Crystal Fibers.

Gaps are not expected unless $\epsilon_1 \ll \epsilon_2$ or $\epsilon_2 \ll \epsilon_1$ for $k = 0$.

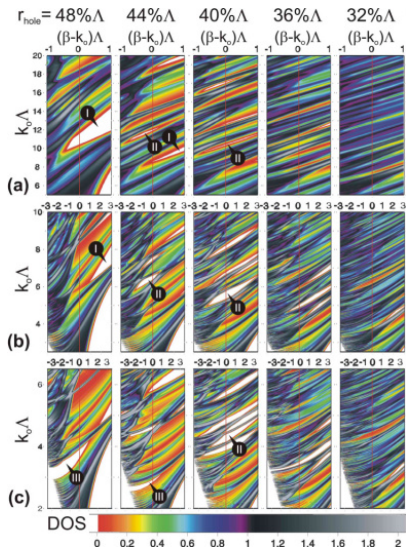
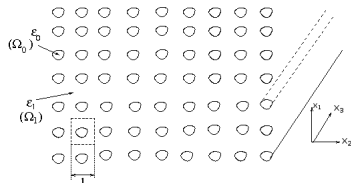


Figure: From J.M.Pottage, D.M.Bird, T.D.Hedley, T.A.Birks, J.C.Knight and P.St.J. Russell, Optics Express, 2003

Maxwell equations for plane waves in PCF, Oblique incidence (Case of PCF): $k \neq 0$

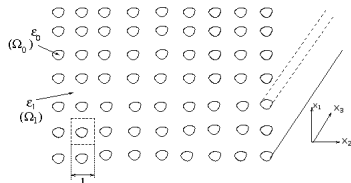


In each phase E_3 and H_3 satisfy the following equations

$$\Delta E_3 + (\omega^2 \epsilon_1 - k^2) E_3 = 0, \quad \Delta H_3 + (\omega^2 \epsilon_1 - k^2) H_3 = 0 \quad \text{in } \Omega_1$$

$$\Delta E_3 + (\omega^2 \epsilon_0 - k^2) E_3 = 0, \quad \Delta H_3 + (\omega^2 \epsilon_0 - k^2) H_3 = 0 \quad \text{in } \Omega_0$$

Maxwell equations for plane waves in PCF, Oblique incidence (Case of PCF): $k \neq 0$



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E_3 and H_3 **coupled across interface** $\Gamma = \partial\Omega_0$:

$$\omega \left[\frac{\epsilon}{a} \nabla E_3 \cdot n \right] = -k \left[\frac{1}{a} \nabla H_3 \cdot n^\perp \right], \quad k \left[\frac{1}{a} \nabla E_3 \cdot n^\perp \right] = \omega \left[\frac{1}{a} \nabla H_3 \cdot n \right]$$

where $a = \omega^2 \epsilon(x) - k^2$ **discontinuous** on Γ .

$$\begin{aligned} \partial_1 \left(\frac{\omega \epsilon}{a} \partial_1 E_3 \right) + \partial_2 \left(\frac{\omega \epsilon}{a} \partial_2 E_3 \right) + \partial_1 \left(\frac{k}{a} \partial_2 H_3 \right) - \partial_2 \left(\frac{k}{a} \partial_1 H_3 \right) &= -\omega \epsilon E_3, \\ -\partial_1 \left(\frac{k}{a} \partial_2 E_3 \right) + \partial_2 \left(\frac{k}{a} \partial_1 E_3 \right) + \partial_1 \left(\frac{\omega}{a} \partial_1 H_3 \right) + \partial_2 \left(\frac{\omega}{a} \partial_2 H_3 \right) &= -\omega H_3, \end{aligned}$$

where $a(x) = \omega^2 \epsilon(x) - k^2$.

Find $u = (E_3, H_3)$ such that

$$\begin{aligned} \int_{\mathbb{R}^2} \frac{\omega}{a} (\epsilon \nabla u_1 \cdot \overline{\nabla \phi_1} + \nabla u_2 \cdot \overline{\nabla \phi_2}) + \frac{k}{a} (\{u_2, \overline{\phi_1}\} - \{u_1, \overline{\phi_2}\}) \, dx \\ = \omega \int_{\mathbb{R}^2} \epsilon u_1 \overline{\phi_1} + u_2 \overline{\phi_2} \, dx \quad \forall \phi \in C_0^\infty(\mathbb{R}^2) \end{aligned}$$

$\{f, g\} := f_{x_2} g_{x_1} - f_{x_1} g_{x_2}$.

The above form is symmetric, and positive if $k^2 < \omega^2 \min\{\epsilon_0, \epsilon_1\}$.

If $k = \omega\kappa$, $\kappa \geq 0$ then we have **usual spectral problem**: Find u such that

$$\int_{\mathbb{R}^2} \frac{1}{\epsilon(x) - \kappa^2} (\epsilon(x) \nabla u_1 \cdot \overline{\nabla \phi_1} + \nabla u_2 \cdot \overline{\nabla \phi_2}) +$$

$$+ \int_{\mathbb{R}^2} \frac{\kappa}{\epsilon(x) - \kappa^2} (\{u_2, \overline{\phi_1}\} - \{u_1, \overline{\phi_2}\}) \, dx = \omega^2 \int_{\mathbb{R}^2} \epsilon(x) u_1 \overline{\phi_1} + u_2 \overline{\phi_2} \, dx,$$

$$\forall \phi \in C_0^\infty(\mathbb{R}^2).$$

The above form is symmetric, and positive if $\kappa^2 < \min\{\epsilon_0, \epsilon_1\}$.

Anti-resonant reflecting optical waveguide (ARROW)

Assume $\epsilon_0 > \epsilon_1 = 1$.

Spectral problem

$$A_\kappa(u, \phi) = \lambda B(u, \phi), \quad \forall \phi \in C_0^\infty(\mathbb{R}^2),$$
$$\lambda = (\epsilon_0 - \kappa^2)\omega^2.$$

$$A_\kappa[u] := \int_{\Omega_1} \kappa \frac{\epsilon_0 - 1}{1 - \kappa^2} |\partial u|^2 + \frac{\epsilon_0 + \kappa}{1 + \kappa} |\nabla u|^2 dx + \int_{\Omega_0} \epsilon_0 |\nabla u_1|^2 + |\nabla u_2|^2 dx$$

where

$$|\partial u|^2 = |\partial_{x_1} u_1 + \partial_{x_2} u_2|^2 + |\partial_{x_2} u_1 - \partial_{x_1} u_2|^2$$

Scalar product is

$$B[u] := \int_{\Omega_1} |u|^2 dx + \int_{\Omega_0} \epsilon_0 |u_1|^2 + |u_2|^2 dx$$

If $\kappa < 1$: then A_κ is positive.

If $\kappa \rightarrow 1$ then there is high contrast.

Periodicity. Floquet transform

$$x \rightarrow (x, \theta), x \in \square = [-\pi, \pi]^2, \theta \in [-1/2, 1/2]^2.$$

Find $u \in H_{\theta}^1(\square)$ ($u(y) = e^{i\theta \cdot x} v(x)$, v \square -periodic) such that

$$A_{\kappa}(u, \phi) = \lambda B(u, \phi), \quad \forall \phi \in H_{\theta}^1(\square).$$

Spectrum:

$$0 \leq \lambda_1(\kappa, \theta) \leq \lambda_2(\kappa, \theta) \leq \dots \leq \lambda_n(\kappa, \theta) \leq \dots$$

Theorem

$$\lim_{\kappa \nearrow 1} \lambda_n(\kappa, \theta) = \lambda_n(\theta).$$

Here $\lambda_n(\theta)$ are eigenvalues of the problem:

Find λ and $u \in V_\theta = \{u \in H_\theta^1(\square) : \partial u = 0 \text{ in } Q_1\}$ such that

$$B(u, \phi) = \lambda A(u, \phi), \quad \forall \phi \in V_\theta,$$

where

$$A[u] := \int_{\square} \epsilon_0 |\nabla u_1|^2 + |\nabla u_2|^2 dx,$$

and

$$B[u] := \int_{\square} |u|^2 dx + (\epsilon_0 - 1) \int_{Q_0} |u_1|^2 dx.$$

Lemma

There exists a constant $c > 0$ such that for any $u \in H_\theta^1(\square)$ there is $v \in V_\theta$ such that $\|u - v\|_{H^1(\square)} \leq c \|\partial u\|_{L_2(Q_1)}$.

Find $\lambda \in \mathbb{C}$ and $u \in V_\theta = \{u \in H_\theta^1(\square) : \partial u = 0 \text{ in } Q_1\}$ such that

$$B(u, \phi) = \lambda A(u, \phi), \quad \forall \phi \in V_\theta,$$

where

$$A[u] := \int_{\square} |\nabla u|^2 + (\epsilon_0 - 1) |\nabla u_1|^2 dx,$$

and

$$B[u] := \int_{\square} |u|^2 dx + (\epsilon_0 - 1) \int_{Q_0} |u_1|^2 dx.$$

What can we say about $\lambda_n(\theta)$?

Let

$$B_\delta := \{x : |x| < \delta\} \subset Q_0.$$

Theorem

$$\lambda_2(\theta) < 8\epsilon_0\delta^{-2} \left(1 + 4 \ln \frac{\pi}{\delta}\right)^{-1},$$

$$\lambda_3(\theta) > \epsilon_0^{-1}\Lambda_2(Q_0),$$

where $\Lambda_2(Q_0)$ is 2nd eigenvalue of Neumann Laplacian in Q_0 .

Corollary

If $8\epsilon_0^2 \leq \Lambda_2(Q_0)\delta^2 \left(1 + 4 \ln \frac{\pi}{\delta}\right)$, then there is a gap.

Find $\lambda \in \mathbb{C}$ and $u \in V_\theta = \{u \in H_\theta^1(\square) : \partial u = 0 \text{ in } Q_1\}$ such that

$$B(u, \phi) = \lambda A(u, \phi), \quad \forall \phi \in V_\theta,$$

where

$$A[u] := \int_{\square} |\nabla u|^2 + (\epsilon_0 - 1) |\nabla u_1|^2 dx,$$

and

$$B[u] := \int_{\square} |u|^2 dx + (\epsilon_0 - 1) \int_{Q_0} |u_1|^2 dx.$$

Find $\lambda \in \mathbb{C}$ and $u \in V_\theta = \{u \in H_\theta^1(\square) : \partial u = 0 \text{ in } Q_1\}$ such that

$$B(u, \phi) = \lambda A(u, \phi), \quad \forall \phi \in V_\theta,$$

where

$$A[u] := \int_{\square} |\nabla u|^2 dx,$$

and

$$B[u] := \int_{\square} |u|^2 dx.$$

Let $\theta \neq 0$. Then for any $u \in H_\theta^1(\square)$ there is $f \in L_2(\square)$ such that

$$u = \partial \Delta_\theta^{-1} f,$$

where

$$\partial = \begin{pmatrix} \partial_1 & \partial_2 \\ \partial_2 & -\partial_1 \end{pmatrix}$$

Then constrain $\partial u = 0$ in Q_1 is equivalent to $f = 0$ in Q_1 .

Problem. $\lambda \in \mathbb{C}$ and $f \in L_2(\square)$, $\text{supp } f \subset \bar{Q}_0$ such that

$$a(f, \phi) = \lambda b(u, \phi), \quad \forall \phi \in L_2(\square), \text{supp } \phi \subset \bar{Q}_0$$

where

$$a(f, \phi) := \langle f, \phi \rangle := \int_{\square} f \bar{\phi} dx = \int_{Q_0} f \bar{\phi} dx,$$

and

$$b[f] := - \int_{Q_0} \bar{f} \Delta_\theta^{-1} f dx.$$

"Small inclusions"

Assume $\overline{Q_0} \subset B_\delta = \{x : |x| < \delta\}$, $\delta < \pi$.

Find $\lambda \in \mathbb{C}$ and $f \in L_2(Q_0)$ such that

$$a(f, \phi) = \lambda b(f, \phi), \quad \forall \phi \in L_2(Q_0),$$

where

$$a(f, f) := \langle f, f \rangle, \quad b(f, f) := \langle -\Delta_\theta^{-1} f, f \rangle.$$

Aim is to "replace" $-\Delta_\theta^{-1} f$ by $-\Delta^{-1} f$, where

$$(-\Delta^{-1} f)(x) = -\frac{1}{2\pi} \int_{Q_0} \ln|x-y| f(y) dy.$$

$$(-\Delta_\theta^{-1} f)(x) = (-\Delta^{-1} f)(x) + \int_{Q_0} g_\theta(x, y) f(y) dy.$$

Then

$$\langle -\Delta_\theta^{-1} f, f \rangle \approx \langle -\Delta^{-1} f, f \rangle + g_\theta(0, 0) \left| \int_{Q_0} f(y) dy \right|^2.$$

Lemma

Let $f \in L_2(\square)$, $\text{supp } f \subset B_\delta$ and $\delta < \pi$. Then

$$\left| \langle (-\Delta_\theta)^{-1} f, f \rangle - \langle (-\Delta)^{-1} f, f \rangle - g_\theta |\langle f, \mathbf{1} \rangle|^2 \right| < \frac{3\delta}{\pi} \left(\langle -\Delta^{-1} f, f \rangle + g_\theta |\langle f, \mathbf{1} \rangle|^2 \right),$$

where $g_\theta = g_\theta(0, 0) > (2\pi)^{-1} \ln \pi$.

Problem 1. Find $\lambda^{(1)} \in \mathbb{C}$ and $f \in L_2(Q_0)$ such that

$$a^{(1)}(f, \phi) = \lambda^{(1)} b^{(1)}(f, \phi), \quad \forall \phi \in L_2(Q_0),$$

where

$$a^{(1)}(f, f) := \langle f, f \rangle, \quad b^{(1)}(f, f) := \langle -\Delta^{-1} f, f \rangle + g_\theta |\langle f, \mathbf{1} \rangle|^2.$$

$$\lambda_n^{(1)} \left(1 + \frac{3\delta}{\pi} \right)^{-1} < \lambda_n < \lambda_n^{(1)} \left(1 - \frac{3\delta}{\pi} \right)^{-1}$$

Consider $Q_0 = \delta\Omega$, where $\Omega \subset B_1$. After rescaling $x = \delta y$ we obtain the following problem.

Problem 2. Find $\lambda^{(2)} \in \mathbb{C}$ and $f \in L_2(\Omega)$ such that

$$a^{(2)}(f, \phi) = \lambda^{(2)} b^{(2)}(f, \phi), \quad \forall \phi \in L_2(\Omega),$$

where

$$a^{(2)}(f, f) := \int_{\Omega} |f|^2 dy, \quad b^{(2)}(f, f) := - \int_{\Omega} \bar{f} \Delta^{-1} f dy + \left(-(2\pi)^{-1} \ln \delta + g_{\theta} \right) \left| \int_{\Omega} f dy \right|^2,$$

and

$$\lambda_n^{(1)} = \delta^{-2} \lambda_n^{(2)}.$$

Here

$$\nu := -(2\pi)^{-1} \ln \delta + g_{\theta}$$

is a big positive parameter.

Consider representation

$$f(y) = \alpha + h(y),$$

where $\alpha \in \mathbb{C}$ and $h \in L_2(\Omega)$, $\int_{\Omega} h dy = 0$.

Problem 3. Find $\lambda^{(3)} \in \mathbb{C}$ and $\alpha \in \mathbb{C}$, $h \in L_2(\Omega)$, $\int_{\Omega} h dy = 0$ such that

$$a^{(3)}(\alpha, h, \beta, \psi) = \lambda b^{(3)}(\alpha, h, \beta, \psi), \quad \forall \phi = \beta + \psi \in L_2(\Omega), \int_{\Omega} \psi dy = 0,$$

where

$$a^{(3)}[\alpha, h] := |\alpha|^2 |\Omega| + \int_{\Omega} |h|^2 dy, \quad b^{(3)}[\alpha, h] := - \int_{\Omega} \overline{(\alpha + h)} \Delta^{-1}(\alpha + h) dy + \nu |\alpha|^2 |\Omega|^2,$$

and

$$\lambda_n^{(2)} = \lambda_n^{(3)}.$$

Problem 4. Find $\lambda^{(4)} \in \mathbb{C}$ and $\alpha \in \mathbb{C}, h \in L_2(\Omega), \int_{\Omega} h dy = 0$ such that

$$a^{(4)}(\alpha, h, \beta, \psi) = \lambda b^{(4)}(\alpha, h, \beta, \psi), \quad \forall \phi \in L_2(\Omega),$$

where

$$a^{(4)}[\alpha, h] := |\alpha|^2 |\Omega| + \int_{\Omega} |h|^2 dy, \quad b^{(4)}[\alpha, h] := - \int_{\Omega} \bar{h} \Delta^{-1} h dy + \nu |\alpha|^2 |\Omega|^2,$$

and

$$\lambda_n^{(4)} \left(1 + c\nu^{-1/2}\right)^{-1} \leq \lambda_n^{(3)} \leq \lambda_n^{(4)} \left(1 - c\nu^{-1/2}\right)^{-1}.$$

Finally

$$\lambda_1^{(4)} = 2\pi |\Omega|^{-1} (2\pi g_{\theta} - \ln \delta)^{-1},$$

and $\lambda_2^{(4)}, \lambda_3^{(4)}, \dots$ are eigenvalues of the problem for

$$a^{(5)}[h] := \int_{\Omega} |h|^2 dy, \quad b^{(5)}[h] := - \int_{\Omega} \bar{h} \Delta^{-1} h dy,$$

with domain $h \in L_2(\Omega), \int_{\Omega} h dy = 0$.

We have similar results for PCF.

$$\lambda_{1,2}(\theta) = \delta^{-2} (2\pi g_\theta - \ln \delta)^{-1} \Lambda_{1,2} + O(\delta^{-2} (2\pi g_\theta - \ln \delta)^{-3/2}),$$

where $\Lambda_{1,2}$ are some positive numbers, and

$$\lambda_n(\theta) = \Lambda_n + O(\delta^{-2} (2\pi g_\theta - \ln \delta)^{-1/2}), \quad n = 3, 4, \dots,$$

where $\Lambda_3, \Lambda_4, \dots$ are the eigenvalues of operator generated by quadratic forms

$$a[f] := \int_{Q_0} |f|^2 dx + (\epsilon_0 - 1) \int_{Q_0} \bar{f} \operatorname{grad} \operatorname{div} \Delta^{-1} f dx,$$

and

$$b[f] := \int_{Q_0} \bar{f} \Delta^{-1} f dx + (\epsilon_0 - 1) \int_{Q_0} \overline{\operatorname{div} \Delta^{-1} f} \operatorname{div} \Delta_\theta^{-1} f dx.$$

with domain $f \in L_2(Q_0)$, $\int_{Q_0} f dx = 0$.

Thank you