

### 3) TRIGONOMETRIC INTERPOLATION

Many real-life functions (e.g. sounds) are periodic, meaning

$$f(x + 2\pi) = f(x) \quad \forall x \in \mathbb{R}. \quad \leftarrow \text{by suitably rescaling } x \text{ we can convert any other period to } 2\pi.$$

Such functions are well-approximated by trigonometric polynomials

$$P_n(x) = \sum_{k=0}^{n-1} (a_k \cos(kx) - b_k \sin(kx)),$$

called degree n if  $a_n$  or  $b_n$  are the highest non-zero coefficients.

Using Euler's identity  $e^{ikx} = \cos(kx) + i\sin(kx)$ , it is nicer to re-write  $p_n(x) = \operatorname{Re}\{q_n(x)\}$  for the complex polynomial

$$\begin{aligned} q_n(x) &= \sum_{k=0}^{n-1} (a_k + i b_k)(\cos(kx) + i \sin(kx)) \\ &= \sum_{k=0}^{n-1} (a_k + i b_k) e^{ikx} = \sum_{k=0}^{n-1} c_k e^{ikx}. \quad \leftarrow = (\cos x + i \sin x)^k \rightarrow \text{explains why it is a "polynomial"} \end{aligned}$$

#### 3.1) THE DISCRETE FOURIER TRANSFORM

Consider the problem of interpolating  $f$  at  $n$  equally-spaced points  $x_j = \frac{2\pi j}{n}$  for  $j=0, \dots, n-1$ . i.e. we want

$$q_n(x_j) = f_j \quad \text{for } j=0, \dots, n-1.$$

This gives

$$\sum_{k=0}^{n-1} c_k e^{ikx_j} = f_j \iff \sum_{k=0}^{n-1} c_k e^{i \frac{2\pi k j}{n}} = f_j \iff \sum_{k=0}^{n-1} c_k \omega^{jk} = f_j \quad \text{where } \omega = e^{i \frac{2\pi}{n}}.$$

This is a system of  $n$  linear equations

$$\underbrace{\begin{bmatrix} \omega^0 & \omega^0 & \omega^0 & \dots & \omega^0 \\ \omega^0 & \omega^1 & \omega^2 & & \omega^{n-1} \\ \omega^0 & \omega^2 & \omega^4 & & \omega^{2(n-1)} \\ \vdots & \vdots & \vdots & & \vdots \\ \omega^0 & \omega^{n-1} & \omega^{2(n-1)} & \dots & \omega^{(n-1)^2} \end{bmatrix}}_{F_n} \begin{bmatrix} c_0 \\ c_1 \\ c_2 \\ \vdots \\ c_{n-1} \end{bmatrix} = \begin{bmatrix} f_0 \\ f_1 \\ f_2 \\ \vdots \\ f_{n-1} \end{bmatrix}.$$

Inverting would give the interpolation coefficients  $\vec{c} = F_n^{-1} \vec{f}$ .

The matrix  $F_n$  is called the Fourier matrix.  $\leftarrow$  Note: books differ in whether they include the  $\frac{1}{\sqrt{n}}$ , or swap  $F_n$  and  $F_n^{-1}$ .

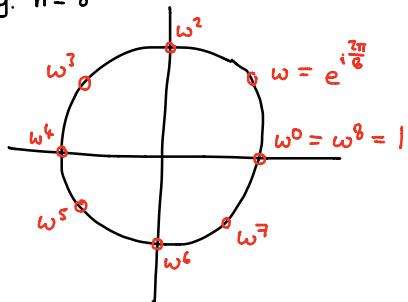
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In general, for any  $n$ -vector  $\vec{x}$ ,  $F_n^{-1} \vec{x}$  is called the discrete Fourier transform (DFT) of  $\vec{x}$ , and  $F_n \vec{x}$  is the inverse DFT of  $\vec{x}$ .  $\leftarrow$  the DFT matrix.

The trigonometric interpolation coefficients  $\vec{c}$  are given by the DFT of the data  $F$ .

Notice that  $\omega = e^{i\frac{2\pi}{n}}$  is a primitive  $n^{\text{th}}$  root of unity, so  $F_n$  contains roots of unity.  
 $\omega^n = 1$  and  $n$  is smallest integer of  $k=1, \dots, n$  for which  $\omega^k = 1$ .

e.g.  $n=8$



it is "almost" unitary,  $F_n^{-1} = F_n^*$  up to  $\frac{1}{n}$ .

Thm 3.1 — The Fourier matrix is symmetric and satisfies  $F_n^{-1} = \frac{1}{n} \bar{F}_n$ .

complex conjugate  
(same as  $F_n^+$  since symmetric)

Proof :-

We want to show that  $\frac{1}{n} \bar{F}_n F_n = I_n$ , i.e.

$$\frac{1}{n} \begin{bmatrix} \omega^0 & \omega^0 & \omega^0 & \cdots & \omega^0 \\ \omega^0 & \omega^{-1} & \omega^{-2} & \cdots & \omega^{-(n-1)} \\ \omega^0 & \omega^{-2} & \omega^{-4} & \cdots & \omega^{-2(n-1)} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \omega^0 & \omega^{-(n-1)} & \omega^{-2(n-1)} & \cdots & \omega^{-(n-1)^2} \end{bmatrix} \begin{bmatrix} \omega^0 & \omega^0 & \omega^0 & \cdots & \omega^0 \\ \omega^0 & \omega^1 & \omega^2 & \cdots & \omega^{(n-1)} \\ \omega^0 & \omega^2 & \omega^4 & \cdots & \omega^{2(n-1)} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \omega^0 & \omega^{(n-1)} & \omega^{2(n-1)} & \cdots & \omega^{(n-1)^2} \end{bmatrix} = \begin{bmatrix} 1 & 0 & \cdots & 0 \\ 0 & 1 & \cdots & \vdots \\ \vdots & \vdots & \ddots & \vdots \\ 0 & \cdots & 0 & 1 \end{bmatrix}.$$

Let

$$\vec{v}^{(k)} = (1, \omega^k, \omega^{2k}, \dots, \omega^{(n-1)k})$$

denote the  $k^{\text{th}}$  row of  $F_n$ .

Diagonal entries:

$$\frac{1}{n} (\bar{F}_n F_n)_{kk} = \frac{1}{n} \bar{v}^{(k)} \cdot \vec{v}^{(k)} = \frac{1}{n} \sum_{l=0}^{n-1} \overline{\omega^{kl}} \omega^{kl} = \frac{1}{n} \sum_{l=0}^{n-1} \omega^{-kl} \omega^{kl} = \frac{1}{n} \sum_{l=0}^{n-1} 1 = 1.$$

Off-diagonal entries:

$$\begin{aligned} \frac{1}{n} (\bar{F}_n F_n)_{jk} &= \frac{1}{n} \bar{v}^{(j)} \cdot \vec{v}^{(k)} = \frac{1}{n} \sum_{l=0}^{n-1} \overline{\omega^{jl}} \omega^{kl} \\ &= \frac{1}{n} \sum_{l=0}^{n-1} \omega^{(k-j)l} \\ &= \frac{1}{n} \sum_{l=0}^{n-1} (\omega^{k-j})^l \quad \text{geometric series.} \\ &= \frac{1}{n} \frac{(\omega^{k-j})^n - 1}{\omega^{k-j} - 1} \\ &= \frac{1}{n} \frac{e^{i(k-j)2\pi} - 1}{e^{i(k-j)2\pi} - 1} = 0. \end{aligned}$$

□

So there is a simple expression for the DFT matrix  $F_n^{-1}$ .