

Bounded solutions to T.I.S.E. $(-\frac{\partial^2}{\partial x^2} + v(x))\psi(x) = k^2 \psi(x)$ come in two flavours:

1. Positive energies $E = k^2 \in (0, \infty)$: "continuous spectrum" \rightarrow "scattering solution"

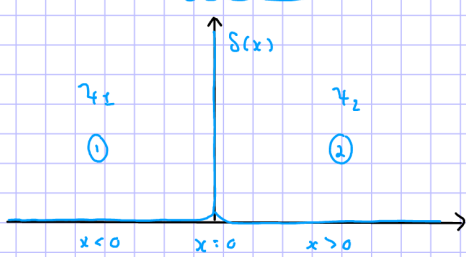
$$\psi(x) \approx \begin{cases} e^{ikx} + R(k)e^{-ikx}, & x \rightarrow -\infty \\ T(k)e^{ikx}, & x \rightarrow +\infty \end{cases} \quad (11.13)$$

2. Negative energies $E = -\mu^2 \in \mathbb{Z}^2 - \mu_1^2, -\mu_2^2, \dots, -\mu_N^2$ (possibly empty set): "discrete spectrum"

set $k = i\mu$ with $\mu > 0$

$$\psi(x) \approx \begin{cases} B e^{\mu x}, & x \rightarrow -\infty \\ C e^{-\mu x}, & x \rightarrow +\infty \end{cases} \rightarrow \text{"bound state solutions"}$$

Example 2: $V(x) = a \delta(x)$



$$\psi(x) = \begin{cases} \psi_1(x), & x < 0 \\ \psi_2(x), & x > 0 \end{cases}$$

$$\begin{aligned} \psi_1(x) &= A e^{ikx} + B e^{-ikx}, & x < 0 \\ \psi_2(x) &= C e^{ikx} + D e^{-ikx}, & x > 0 \end{aligned} \quad (11.23)$$

Matching condition # 1: $\psi_1(0) = \psi_2(0)$

Matching condition # 2: $\psi_2'(0) - \psi_1'(0) = a \psi(0)$

$$\Rightarrow \begin{cases} A = (1 - \frac{a}{2ik})C - \frac{a}{2ik}D \\ B = \frac{a}{2ik}C + (1 + \frac{a}{2ik})D \end{cases} \quad (11.28)$$

plug into (11.23) to give general solution for all x!

Two cases:

(a) Positive energies $E = k^2$ and $k > 0$. For scattering solution, set $D = 0$ and divide by $A = (1 - \frac{a}{2ik})C$:

$$\psi(x) = \begin{cases} e^{ikx} + \frac{a}{2ik - a} e^{-ikx}, & x < 0 \\ \frac{2ik}{2ik - a} e^{ikx}, & x > 0 \end{cases} \quad (11.29)$$

$$R(k) = \frac{a}{2ik - a}, \quad T(k) = \frac{2ik}{2ik - a}$$

(b) Bound states? $E = -\mu^2$ and wlog $\mu > 0$. Set $k = i\mu$ in (11.23) and (11.28)

$$\psi(x) = \begin{cases} A e^{-\mu x} + B e^{\mu x}, & x < 0 \\ C e^{-\mu x} + D e^{\mu x}, & x > 0 \end{cases}$$

ψ is bounded $\rightarrow A = D = 0$. Then we have to check consistency with (11.28):

$$0 = (1 + \frac{a}{2\mu})C$$

$$B = -\frac{a}{2\mu}C$$

two possibilities:

① $A = B = C = D = 0$ (trivial solution)

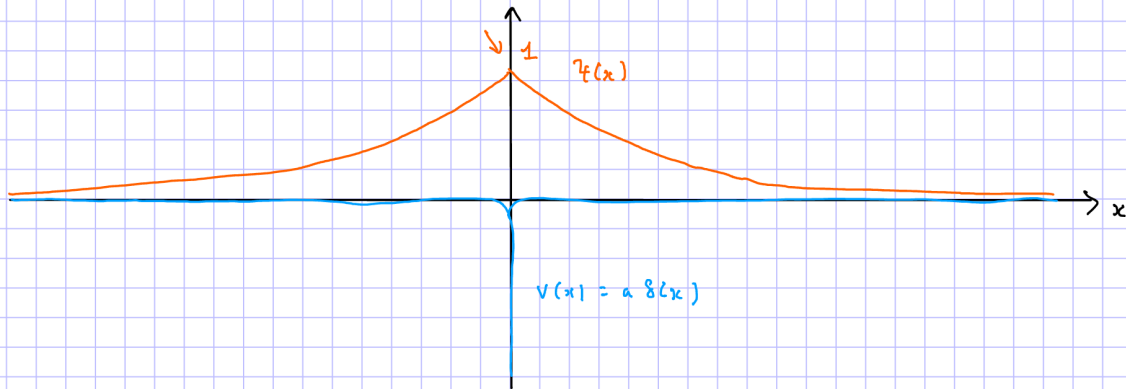
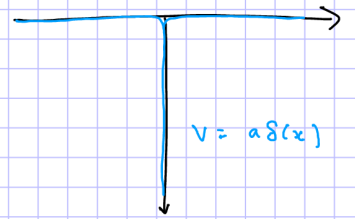
② $\mu = -a/2 \rightarrow B = C$ and $A = D = 0$ (since $\mu > 0$ this only exists for $a < 0$)

This is a simple bound state solution for $a < 0$ given by

(normalize $C = 1$):

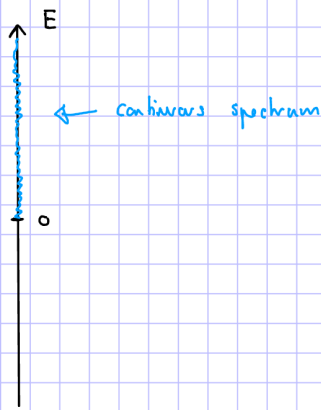
$$\psi(x) = \begin{cases} e^{-a/2 x} & , x < 0 \\ e^{a/2 x} & , x > 0 \end{cases}$$

with energy $E = -\frac{a^2}{4}$

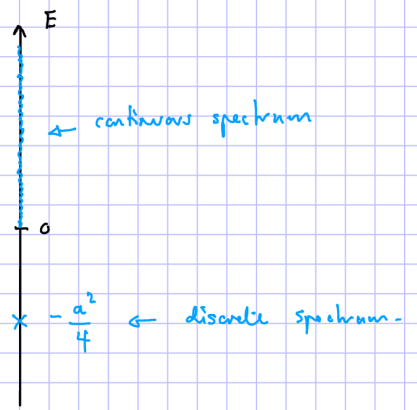


Summary of energy spectrum of bound solutions:

$a > 0$



$a < 0$



Question: Can I obtain bound state solutions from the knowledge of the scattering solution?

Answer: Yes!

The general story

(a) For all $E = k^2 > 0$ and wlog $k > 0$ we find the scattering solution:

$$\psi(x) \sim \begin{cases} e^{ikx} + R(k)e^{-ikx} & , x \rightarrow -\infty \\ T(k)e^{ikx} & , x \rightarrow +\infty \end{cases} \quad (11.37)$$

[N.B. R and T are meromorphic in the complex variable k and have simple poles on the imaginary axis.]

(b) $E = -\mu^2$ and wlog $\mu > 0$

Step #1: Set $k = i\mu$ in (11.37)

$$\psi(x) \approx \begin{cases} e^{-\mu x} + R(i\mu) e^{\mu x}, & x \rightarrow -\infty \\ T(i\mu) e^{\mu x}, & x \rightarrow +\infty \end{cases}$$

two problems:

1. This is not bounded ($e^{-\mu x} \rightarrow \infty$ as $x \rightarrow -\infty$)
2. R and T might contain poles at certain real values of μ

step #2: Divide by $T(i\mu)$:

$$\psi(x) \approx \begin{cases} \frac{1}{T(i\mu)} e^{-\mu x} + \frac{R(i\mu)}{T(i\mu)} e^{\mu x}, & x \rightarrow -\infty \\ e^{-\mu x}, & x \rightarrow +\infty \end{cases}$$

This is bounded at the values of μ where $T(i\mu)$ has a pole since $1/T(i\mu) = 0$

For such a value of μ :

$$\psi(x) \approx \begin{cases} \frac{R(i\mu)}{T(i\mu)} e^{\mu x}, & x \rightarrow -\infty \\ e^{-\mu x}, & x \rightarrow +\infty \end{cases}$$

This only exists for values of μ where $T(i\mu)$ has a pole for $\mu > 0$

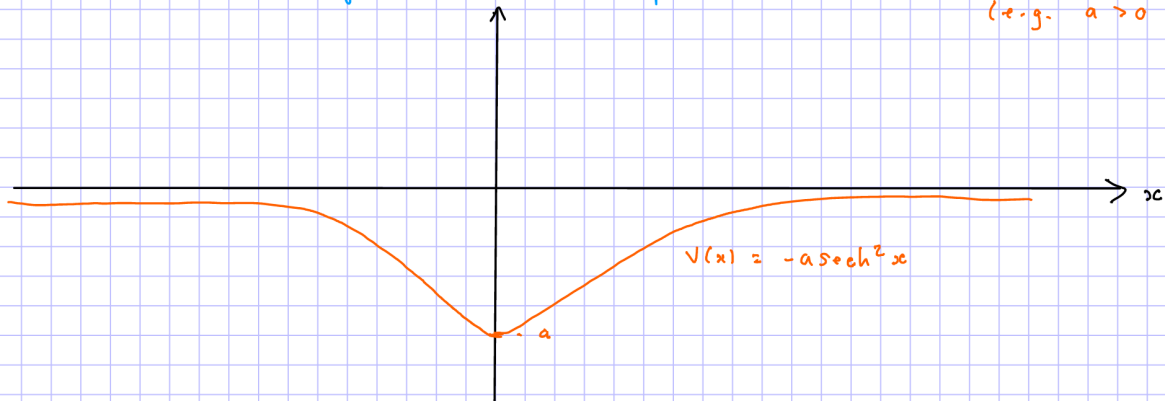
Poles in $T(i\mu)$ for $\mu > 0 \Leftrightarrow$ bound state solution.

11.3 Reflectionless potentials

Consider

$$V(x) = -a \operatorname{sech}^2 x \quad (11.42)$$

[recall: $u = -V = a \operatorname{sech}^2 x$ where for $a = n(n+1)$ and n is a positive integer this seemed to be the form of the initial data for pure n -soliton solutions to the KdV equation (e.g. $a > 0$)]



The T.I.S.E. takes the form:

$$-\psi''(x) - a \operatorname{sech}^2 x \psi = E \psi \quad E = k^2$$

goal: look for bounded solutions

make change of variables

$$y = \tanh x \in (-1, 1)$$

$$\begin{aligned} \frac{d}{dx} &= \frac{dy}{dx} \frac{d}{dy} = \operatorname{sech}^2 x \frac{d}{dy} \\ &= (1 - \tanh^2 x) \frac{d}{dy} \\ &= (1 - y^2) \frac{d}{dy} \end{aligned}$$



The T. I. S. E. becomes

$$\frac{d}{dy} \left[(1 - y^2) \frac{d\psi}{dy} \right] + \left(\frac{k^2}{1 - y^2} + a \right) \psi = 0$$

if we make the replacements:

$$k^2 = -m^2 \quad \text{and} \quad a = n(n+1)$$

This becomes

$$\frac{d}{dy} \left[(1 - y^2) \frac{d\psi}{dy} \right] + \left(n(n+1) - \frac{m^2}{1 - y^2} \right) \psi = 0$$

General (or associated) Legendre equation.