

§4.2 *Non-canonical kinetic terms*

Finally, we consider configurations with non-canonical kinetic terms of the form

$$L = \frac{1}{2} \sum_{i,j} \mathbf{B}_{ij}(q) \dot{q}_i \dot{q}_j - V(q). \quad (4.2.1)$$

We still obtain a linear differential operator if we restrict $\mathbf{B}(q) \rightarrow \mathbf{B}(0)$. Physically, this corresponds to considering oscillations with not too much kinetic energy, which makes sense if we want to stay at the minimum. The resulting equations of motion are

$$\mathbf{B}\ddot{\mathbf{q}} + \mathbf{A}\mathbf{q} = 0 \quad (4.2.2)$$

where we have defined $\mathbf{B} \equiv \mathbf{B}(0)$, a constant matrix. \mathbf{B} generally does not have zero eigenvalues (since this would correspond to generalised coordinates without a kinetic term), so we assume no zero eigenvalues. This implies that $\det(\mathbf{B}) \neq 0$, and so \mathbf{B}^{-1} exists. We then have an equivalent set of equations

$$\ddot{\mathbf{q}} + \mathbf{B}^{-1}\mathbf{A}\mathbf{q} = 0 \quad (4.2.3)$$

which reduces to the case we have already studied if you define $\mathbf{C} := \mathbf{B}^{-1}\mathbf{A}$.

There is one small subtlety that needs to be mentioned here: the fact that \mathbf{A} was symmetric was quite important in our discussion above, since it ensured that its eigenvalues were real, but in general $\mathbf{B}^{-1}\mathbf{A}$ will not be symmetric, even if both \mathbf{B}^{-1} and \mathbf{A} separately are. Let us assume that \mathbf{A} is positive semi-definite: that is, all its eigenvalues are either positive or zero. Then it exists a symmetric matrix $\mathbf{A}^{\frac{1}{2}}$ such that $(\mathbf{A}^{\frac{1}{2}})^2 = \mathbf{A}$.¹² We can use this matrix to rewrite

$$\mathbf{C} := \mathbf{B}^{-1}\mathbf{A} = \mathbf{A}^{-\frac{1}{2}} \left(\mathbf{A}^{\frac{1}{2}}\mathbf{B}^{-1}\mathbf{A}^{\frac{1}{2}} \right) \mathbf{A}^{\frac{1}{2}}$$

so we find that \mathbf{C} is similar (in the sense of similarity transformations of matrices) to $\mathbf{A}^{\frac{1}{2}}\mathbf{B}^{-1}\mathbf{A}^{\frac{1}{2}}$. This matrix is manifestly symmetric, so its eigenvalues are real. Since similar matrices have the same eigenvalues, the eigenvalues of \mathbf{C} will be real too. It is straightforward to check that if \mathbf{v} is an eigenvector of $\mathbf{A}^{\frac{1}{2}}\mathbf{B}^{-1}\mathbf{A}^{\frac{1}{2}}$ with eigenvalue λ , then $\mathbf{A}^{-\frac{1}{2}}\mathbf{v}$ will be an eigenvector of \mathbf{C} with the same eigenvalue. So, in practice, we can simply compute the eigenvalues and eigenvectors of \mathbf{C} , and proceed as we did above.

¹²The simplest way to prove this is to note that \mathbf{A} is real symmetric, and thus diagonalisable by an orthogonal transformation O as $\mathbf{A} = ODO^t$. We can then define $\mathbf{A}^{\frac{1}{2}} = OD^{\frac{1}{2}}O^t$.