

1. Note that

$$df = \frac{\partial f}{\partial x_1} dx_1 + \frac{\partial f}{\partial x_2} dx_2 + \frac{\partial f}{\partial x_3} dx_3.$$

On the other hand, we have

$$\nabla f = \left(\frac{\partial f}{\partial x_1}, \frac{\partial f}{\partial x_2}, \frac{\partial f}{\partial x_3} \right),$$

i.e.,

$$\omega_{\nabla f} = \frac{\partial f}{\partial x_1} dx_1 + \frac{\partial f}{\partial x_2} dx_2 + \frac{\partial f}{\partial x_3} dx_3.$$

Next, note for a vector field $F = (f_1, f_2, f_3)$ we have

$$\text{curl } F = \left(\frac{\partial f_3}{\partial x_2} - \frac{\partial f_2}{\partial x_3}, \frac{\partial f_1}{\partial x_3} - \frac{\partial f_3}{\partial x_1}, \frac{\partial f_2}{\partial x_1} - \frac{\partial f_1}{\partial x_2} \right),$$

i.e.,

$$\eta_{\text{curl } F} = \left(\frac{\partial f_3}{\partial x_2} - \frac{\partial f_2}{\partial x_3} \right) dx_2 \wedge dx_3 + \left(\frac{\partial f_1}{\partial x_3} - \frac{\partial f_3}{\partial x_1} \right) dx_3 \wedge dx_1 + \left(\frac{\partial f_2}{\partial x_1} - \frac{\partial f_1}{\partial x_2} \right) dx_1 \wedge dx_2.$$

On the other hand, since $\omega_F = f_1 dx_1 + f_2 dx_2 + f_3 dx_3$,

$$d\omega_F = \left(\frac{\partial f_1}{\partial x_2} dx_2 + \frac{\partial f_1}{\partial x_3} dx_3 \right) \wedge dx_1 + \left(\frac{\partial f_2}{\partial x_1} dx_1 + \frac{\partial f_2}{\partial x_3} dx_3 \right) \wedge dx_2 + \left(\frac{\partial f_3}{\partial x_1} dx_1 + \frac{\partial f_3}{\partial x_2} dx_2 \right) \wedge dx_3,$$

which shows $\eta_{\text{curl } F} = d\omega_F$, after rearranging the latter expression and using the fact that $dx_i \wedge dx_j = -dx_j \wedge dx_i$. Finally, for $F = (f_1, f_2, f_3)$ and $\eta_F = f_1 dx_2 \wedge dx_3 + f_2 dx_3 \wedge dx_1 + f_3 dx_1 \wedge dx_2$, we obtain

$$d\eta_F = \frac{\partial f_1}{\partial x_1} dx_1 \wedge dx_2 \wedge dx_3 + \frac{\partial f_2}{\partial x_2} dx_2 \wedge dx_3 \wedge dx_1 + \frac{\partial f_3}{\partial x_3} dx_3 \wedge dx_1 \wedge dx_2,$$

i.e.,

$$d\eta_F = \left(\frac{\partial f_1}{\partial x_1} + \frac{\partial f_2}{\partial x_2} + \frac{\partial f_3}{\partial x_3} \right) dx_1 \wedge dx_2 \wedge dx_3 = \text{div } F dx_1 \wedge dx_2 \wedge dx_3.$$

The above calculations show that the following diagram commutes:

$$\begin{array}{ccccccc} C^\infty(\mathbb{R}^3) & \xrightarrow{\nabla} & \mathcal{X}(\mathbb{R}^3) & \xrightarrow{\text{curl}} & \mathcal{X}(\mathbb{R}^3) & \xrightarrow{\text{div}} & C^\infty(\mathbb{R}^3) \\ \downarrow \Phi_0 & & \downarrow \Phi_1 & & \downarrow \Phi_2 & & \downarrow \Phi_3 \\ \Omega^0(\mathbb{R}^3) & \xrightarrow{d} & \Omega^1(\mathbb{R}^3) & \xrightarrow{d} & \Omega^2(\mathbb{R}^3) & \xrightarrow{d} & \Omega^3(\mathbb{R}^3) \end{array}$$

Here, $\mathcal{X}(U)$ denotes the space of smooth vector fields on $U \subset \mathbb{R}^n$, and the vertical maps Φ_i are **bijective maps** between functions/vector fields and differential forms and defined as follows:

$$\Phi_0(f) = f, \quad \Phi_1(F) = \omega_F, \quad \Phi_2(F) = \eta_F, \quad \Phi_3(f) = f dx_1 \wedge dx_2 \wedge dx_3.$$

Since the vertical maps are bijective, we see that $d^2 = 0$ translates into $\text{curl} \circ \nabla = 0$ and $\text{div} \circ \text{curl} = 0$.

2. (a) We have

$$\begin{aligned} \left| \int_c \omega \right| &= \left| \int_a^b \omega_{c(t)}(c'(t)) dt \right| \leq \int_a^b |\langle F_\omega(c(t)), c'(t) \rangle| dt \leq \\ &\leq \int_a^b \|F_\omega(c(t))\| \cdot \|c'(t)\| dt \leq M \int_a^b \|c'(t)\| dt = ML(c). \end{aligned}$$

(b) According to Proposition 5.13, we only have to prove that we have $\int_c \omega = 0$ for all closed curves $c : [a, b] \rightarrow \mathbb{R}^n - 0$. Choose $M > 0$ and $r > 0$ such that $\|F_\omega(x)\| \leq M$ for all $\|x\| \leq r$. Let $\epsilon > 0$ be arbitrary. We consider the free homotopy $G : [a, b] \times [\epsilon, 1] \rightarrow \mathbb{R}^n - 0$, defined by $G(t, s) = s \cdot c(t)$. Since G is a free homotopy and ω is closed, we conclude that

$$\int_c \omega = \int_{c_\epsilon} \omega,$$

by Corollary 6.13. Note also that $L(c_\epsilon) = \epsilon \cdot L(c)$, since $c'_\epsilon(t) = \epsilon c'(t)$. This implies with (a) that

$$\left| \int_c \omega \right| = \left| \int_{c_\epsilon} \omega \right| \leq M \cdot L(c_\epsilon) = \epsilon \cdot M \cdot L(c).$$

Since $\epsilon > 0$ was arbitrary, we must have $\int_c \omega = 0$. This is what we wanted to show.

(c) Note that $F_\omega(x, y) = \left(-\frac{y}{x^2+y^2}, \frac{x}{x^2+y^2} \right)$ and

$$\|F_\omega(0, y)\| = \frac{1}{|y|}.$$

Note that $1/|y|$ is not bounded for any disk of centre 0, so we cannot apply (b) in this case.

3. (a) Since $c'_x(t) = x = (x_1, x_2)$, we have

$$\begin{aligned} f(x) &= \int_{c_x} \omega = \int_0^1 \omega_{c_x(t)}(c'_x(t)) dt = \int_0^1 f_1(c_x(t)) dx_1(c'_x(t)) + f_2(c_x(t)) dx_2(c'_x(t)) = \\ &= \int_0^1 f_1(tx_1, tx_2)x_1 + f_2(tx_1, tx_2)x_2 dt. \end{aligned}$$

(b) Since $\omega = f_1 dx_1 + f_2 dx_2$ is closed, we have $\frac{\partial f_1}{\partial x_2} = \frac{\partial f_2}{\partial x_1}$. Using this, we

obtain

$$\begin{aligned}
\frac{\partial f}{\partial x_1}(x) &= \int_0^1 \frac{\partial}{\partial x_1} (f_1(tx_1, tx_2)x_1 + f_2(tx_1, tx_2)x_2) dt \\
&= \int_0^1 \left(t \frac{f_1}{\partial x_1}(tx_1, tx_2)x_1 + f_1(tx_1, tx_2) + t \frac{f_2}{\partial x_1}(tx_1, tx_2)x_2 \right) dt \\
&= \int_0^1 \left(t \frac{f_1}{\partial x_1}(tx_1, tx_2)x_1 + t \frac{f_1}{\partial x_2}(tx_1, tx_2)x_2 \right) dt + \int_0^1 f_1(tx_1, tx_2) dt \\
&= \int_0^1 t \langle \nabla f(c_x(t)), c'_x(t) \rangle dt + \int_0^1 f_1(c_x(t)) dt \\
&= \int_0^1 t Df_1(c_x(t))(c'_x(t)) dt + \int_0^1 f_1(c_x(t)) dt \\
&= \int_0^1 t (f_1 \circ c_x)'(t) + f_1 \circ c_x(t) dt,
\end{aligned}$$

where in the last step we applied the chain rule. Partial integration yields

$$\frac{\partial f}{\partial x_1}(x) = [t f_1 \circ c_x(t)]_0^1 - \int_0^1 f_1 \circ c_x(t) dt + \int_0^1 f_1 \circ c_x(t) dt = 1 \cdot f_1(x) - 0 \cdot f_1(0) = f_1(x).$$

Similarly, one shows $\frac{\partial f}{\partial x_2}(x) = f_2(x)$, and we conclude that

$$df = \frac{\partial f}{\partial x_1} dx_1 + \frac{\partial f}{\partial x_2} dx_2 = f_1 dx_1 + f_2 dx_2 = \omega,$$

i.e., ω is exact.