

# PROBLEMS CLASS 5

**Ex 4** Let  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ .

1. Show that the "vacuum Maxwell eqns" (i.e. with  $J^\mu = 0$ ) have plane wave solutions of the form

$$A_\mu = \text{Re} \left( p_\mu e^{ik_\nu x^\nu} \right)$$

wavenumber 4-vector  
 (momentum 4-vector)

polarization vector

for constant vectors  $p_\mu, k_\mu$ . Which conditions must  $p_\mu, k_\mu$  obey?

The EoM is

$$0 = \partial^\nu F_{\mu\nu} = \partial^\nu (\partial_\mu A_\nu - \partial_\nu A_\mu) = 0.$$

Letting  $A_\mu = \text{Re} (p_\mu e^{ik \cdot x})$  gives

$$0 = \text{Re} \left( \left( (ik^\nu)(ik_\mu) p_\nu - (ik^\nu)(ik_\nu) p_\mu \right) e^{ik \cdot x} \right)$$

NOTATION:  
 $a \cdot b \equiv a^\mu b_\mu$   
 $a^2 = a \cdot a$

$\forall x$

so the plane wave is a solution provided that

$$-k \cdot p k_\mu + k^2 p_\mu = 0 \quad \forall \mu = 0, 1, 2, 3.$$

2. Taking into account the freedom of gauge transformations, how many independent components of  $p_\mu$  appear in a general plane wave solution? [Hint: try  $\alpha = \text{Re}(c e^{ik \cdot x})$ ]

$\uparrow$  const.

Under a gauge transfo,

Note: this gauge transfo keeps the plane wave a plane wave

$$A_\mu \mapsto A_\mu + \partial_\mu \alpha$$

$$\Rightarrow \text{Re}(p_\mu e^{ikx}) \mapsto \text{Re}((p_\mu + ick_\mu) e^{ikx})$$

i.e. 
$$p_\mu \mapsto p_\mu + ick_\mu \quad (*)$$

Taking  $c \in i \cdot \mathbb{R}$  (purely imaginary), this allows us to eliminate the component of  $p_\mu$  parallel to  $k_\mu$ .  
(longitudinal component!)

$$k \cdot p \quad k_\mu = k^2 p_\mu$$

implies that

$$k \cdot p = k^2 = 0 \quad (**)$$

- $k^2 \equiv k^\mu k_\mu = 0$  is the dispersion relation of electromagnetism, which tells us that  $k_0 = \pm |\vec{k}|$ .  

$$\omega$$
- $k \cdot p \equiv k^\mu p_\mu = 0$  says that the polarization vector  $p_\mu$  is orthogonal (with the Minkowski metric) to the wave 4-vector  $k_\mu$  (or equivalently, momentum) of the electromagnetic wave.
- The gauge freedom (\*) eliminates the longitudinal component of  $p_\mu$ .
- The equation  $p \cdot k = 0$  eliminates another component.

$$\Rightarrow \left[ \begin{array}{l} \# \text{ of independent components ("degrees of freedom")} \\ \text{of } p_\mu \end{array} \right] = 4 - 1 - 1 = 2.$$

NOTE: using (\*\*) in (\*) we find

$$0 = k \cdot p \mapsto k \cdot p + ick^2 = k \cdot p = 0$$

$$p^2 \mapsto (p + ick)^2 = p^2 + 2ick \cdot p - c^2 k^2 = p^2$$

so  $k \cdot p$  and  $p^2$  are gauge invariant.

**Ex 6**

Consider a complex scalar field  $\phi$  with Lagrangian density

$$\mathcal{L} = -|\partial_\mu \phi|^2 - V(\phi, \bar{\phi}) \equiv -|\partial_\mu \phi|^2 - \lambda (|\phi|^2 - a^2)^2 \quad (\lambda, a > 0)$$

$$- \overbrace{\partial^\mu \phi \partial_\mu \phi}^{\equiv} = |\partial_0 \phi|^2 - |\vec{\nabla} \phi|^2$$

1. Show that the energy (or 'Hamiltonian') is

$$E = \int d^3x (|\partial_0 \phi|^2 + |\partial_i \phi|^2 + V(\phi, \bar{\phi}))$$

NOTATIONS:  $\overbrace{\partial_i \bar{\phi} \partial_i \phi}^{\equiv} = \sum_{i=1}^3 |\partial_i \phi|^2 \equiv |\vec{\nabla} \phi|^2$  (All these notations are fine)

The Hamiltonian is the integral over space  $\mathbb{R}^3$  of the Hamiltonian density

$$\mathcal{H} = \overbrace{\frac{\partial \mathcal{L}}{\partial \partial_0 \phi}}^{\pi_\phi} \partial_0 \phi + \overbrace{\frac{\partial \mathcal{L}}{\partial \partial_0 \bar{\phi}}}^{\pi_{\bar{\phi}}} \partial_0 \bar{\phi} - \mathcal{L}$$

(conjugate momenta to  $\phi, \bar{\phi}$ )

$$= (\partial_0 \bar{\phi}) (\partial_0 \phi) + (\partial_0 \phi) (\partial_0 \bar{\phi}) - (\partial_0 \bar{\phi}) (\partial_0 \phi) + \vec{\nabla} \bar{\phi} \cdot \vec{\nabla} \phi + V(\bar{\phi}, \phi)$$

$$= |\partial_0 \phi|^2 + |\vec{\nabla} \phi|^2 + V(\bar{\phi}, \phi)$$

$$\Rightarrow E = \int d^3x (|\partial_0 \phi|^2 + |\vec{\nabla} \phi|^2 + V(\bar{\phi}, \phi))$$

2. Find the configurations of least energy ("vacua"/"ground states") and show that they parametrize a circle in field space.

All 3 terms in  $\mathcal{H}$  are separately  $\geq 0$ , so the energy is minimized by setting them to zero separately:

$$\partial_0 \phi = 0, \quad \vec{\nabla} \phi = 0, \quad |\phi|^2 = a^2$$

Hence the ground states/vacua are constant field configurations that minimize the scalar potential. Here

$$\left\{ \phi = \underset{\substack{\uparrow \\ \text{constants}}}{a} e^{i\beta} \mid \beta \sim \beta + 2\pi \right\} \cong S^1 \quad \left( \begin{array}{l} \text{Circle of radius} \\ a \text{ in the complex} \\ \phi \text{ plane} \end{array} \right)$$

3. Show that any two vacua are related by a global  $U(1)$  transformation.

• Global (internal)  $U(1)$  transfo:

$$(*) \quad \phi \mapsto e^{i\alpha} \phi = \phi' \quad (\alpha \sim \alpha + 2\pi)$$

• Vacuum 1:  $\phi = a e^{i\beta_1} \equiv \phi_1$   
 Vacuum 2:  $\phi = a e^{i\beta_2} \equiv \phi_2$

Perform  $(*)$  with  $\alpha = \beta_2 - \beta_1$ :

$$\phi_1 = a e^{i\beta_1} \mapsto \phi'_1 = e^{i(\beta_2 - \beta_1)} a e^{i\beta_1} = a e^{i\beta_2} = \phi_2$$

**Ex 7** Show that the commutator of two covariant derivatives  $D_\mu = \partial_\mu - iA_\mu$  is given by

$$-i F_{\mu\nu} \equiv [D_\mu, D_\nu] = -i(\partial_\mu A_\nu - \partial_\nu A_\mu)$$

We need to use

$$[f, g] = 0 \quad \forall f, g \text{ functions of } x$$

$$[\partial_{\mu_1}, f] = \partial_{\mu_1} f - f \partial_{\mu_1} = (\partial_{\mu_1} f) \mathbb{1} + f \partial_{\mu_1} - f \partial_{\mu_1} = (\partial_{\mu_1} f) \mathbb{1} \quad \forall \text{ functions } f \text{ of } x$$

$$[\partial_\mu, \partial_\nu] = 0$$

where the above are viewed as differential operators (functions of  $x$  act by multiplication), and

$$[X, Y] = Z \quad \text{as operators}$$

iff  $[X, Y]\psi = Z\psi \quad \forall$  smooth functions  $\psi$  of  $x$ .

Indeed

$$[f, g]\psi = fg\psi - gf\psi = 0$$

$$[\partial_\mu, f]\psi = \partial_\mu(f\psi) - f(\partial_\mu\psi) = (\partial_\mu f)\psi + f(\cancel{\partial_\mu\psi}) - f(\partial_\mu\psi) = (\partial_\mu f)\psi$$

$$[\partial_\mu, \partial_\nu]\psi = \partial_\mu\partial_\nu\psi - \partial_\nu\partial_\mu\psi = 0.$$

Then

$$F_{\mu\nu}\mathbb{1} = i[D_\mu, D_\nu] = i[\partial_\mu - iA_\mu, \partial_\nu - iA_\nu]$$

$$= i([\partial_\mu, \partial_\nu] - i[\partial_\mu, A_\nu] - i[A_\mu, \partial_\nu] - [A_\mu, A_\nu])$$

$$= [\overset{\circ}{\partial}_\mu, \overset{\circ}{A}_\nu] - [\partial_\nu, A_\mu] = (\partial_\mu A_\nu - \partial_\nu A_\mu)\mathbb{1}.$$