

Ex 24

1. Consider a rep \underline{r} of a compact Lie group G , w/ generators $t_a^{(\underline{r})}$
 Show that the set of generators

$$t_a^{(\bar{r})} := -(t_a^{(\underline{r})})^T$$

defines another rep of G , which is called the complex conjugate rep \bar{r} .

We need to show that (drop underlines for ease of notation)

$$[t_a^{(\underline{r})}, t_b^{(\underline{r})}] = i f_{ab}^c t_c^{(\underline{r})} \Rightarrow [t_a^{(\bar{r})}, t_b^{(\bar{r})}] = i f_{ab}^c t_c^{(\bar{r})}$$

$$\begin{aligned} [t_a^{(\bar{r})}, t_b^{(\bar{r})}] &= [-(t_a^{(\underline{r})})^T, -(t_b^{(\underline{r})})^T] = [t_b^{(\underline{r})}, t_a^{(\underline{r})}]^T \\ &= -[t_a^{(\underline{r})}, t_b^{(\underline{r})}]^T = -i f_{ab}^c (t_c^{(\underline{r})})^T = i f_{ab}^c t_c^{(\bar{r})} \end{aligned}$$

2. Show that if the column vector ϕ transforms in the irrep \underline{r} , then its complex conjugate $\phi^* = \bar{\phi}$, which is also a column vector, transforms in the complex conjugate irrep \bar{r} .

$$\begin{aligned} \phi &\mapsto e^{i\alpha^a t_a^{(\underline{r})}} \phi \\ &\quad \uparrow \\ &\quad r(g), \quad g = e^{i\alpha^a t_a} \\ \Rightarrow \bar{\phi} &\mapsto e^{-i\alpha^a \overline{t_a^{(\underline{r})}}} \bar{\phi} = e^{i\alpha^a (-t_a^{(\underline{r})})^T} \bar{\phi} = e^{i\alpha^a t_a^{(\bar{r})}} \bar{\phi} \\ &\quad \uparrow \qquad \qquad \qquad \uparrow \\ &\quad \alpha^a \in \mathbb{R} \qquad \qquad \qquad t_a^{(\bar{r})} = \overline{(t_a^{(\underline{r})})^T} \end{aligned}$$

for reps of compact Lie groups

3. Denote $\bar{\phi}_j := (\phi^j)^*$, $\phi^T = \bar{\phi}^T = (\bar{\phi}_1, \dots, \bar{\phi}_r)$ ← row vector
 Show that the inner product $\phi^T \phi$ is invariant under the action of G .

$$\phi^T \mapsto \left(e^{i\alpha^a t_a^{(r)}} \phi \right)^T = \phi^T e^{-i\alpha^a t_a^{(r)}} = (t_a^{(r)})^T$$

$$\Rightarrow \phi^T \phi \mapsto \phi^T e^{-i\alpha^a t_a^{(r)}} e^{i\alpha^b t_b^{(r)}} \phi = \phi^T \phi$$

$= e^{-i\alpha^a t_a^{(r)}} e^{i\alpha^b t_b^{(r)}} = \mathbb{1}_{r \times r}$
↑ ↑
matrices

Ex 25

1. The adjoint action of the Lie algebra \mathfrak{g} on itself is

$$\begin{aligned} \text{ad}_x: \mathfrak{g} &\rightarrow \mathfrak{g} \\ y &\mapsto \text{ad}_x(y) := [x, y] \end{aligned}$$

$\forall x \in \mathfrak{g}$. Show that

$$\text{ad}_{t_a}(y^b t_b) = (t_a^{(\text{adj})})^b{}_c y^c t_b, \text{ where } (t_a^{(\text{adj})})^b{}_c = i f_{ac}{}^b$$

are the generators of the adjoint rep.

$$\begin{aligned} \text{ad}_{t_a}(y^b t_b) &= [t_a, y^b t_b] = y^b [t_a, t_b] = y^b \cdot i f_{ab}{}^c t_c \\ &= i f_{ac}{}^b y^c t_b = (t_a^{(\text{adj})})^b{}_c y^c t_b \end{aligned}$$

2. Express the Killing form

$$K_{ab} = \text{tr}_{\text{adj}}(t_a^{(\text{adj})} t_b^{(\text{adj})})$$

in terms of the structure constants.

$$K_{ab} = (t_a^{(\text{adj})})^d{}_c (t_b^{(\text{adj})})^c{}_d = if_{ac}{}^d \cdot if_{bd}{}^c$$

$$= -f_{ac}{}^d f_{bd}{}^c = f_{ac}{}^d f_{db}{}^c$$

3. Show that, in a basis where the Killing form is

$$K_{ab} = C(\text{adj}) \delta_{ab},$$

the quadratic invariant $C(\text{adj})$ of the adjoint rep is given by

$$C(\text{adj}) = \frac{\delta^{ab} f_{acd} f_b{}^{cd}}{\dim g}$$

$$C(\text{adj}) \delta_{ab} = f_{ac}{}^d f_{db}{}^c$$

$$\delta^{ba} \uparrow : C(\text{adj}) \delta_a^a = \delta^{ba} f_{ac}{}^d f_{db}{}^c = -\delta^{ab} f_{ac}{}^d f_{bd}{}^c$$

" $\dim g$ $f_{db}{}^c = -f_{bd}{}^c$

$$= -\delta^{ab} f_{acd} f_b{}^{dc} = \delta^{ab} f_{adc} f_b{}^{dc}$$

Lower & Raise d \uparrow $f_{acd} = -f_{adc}$

(either with κ or δ , both symmetric)

$$= \delta^{ab} f_{acd} f_b{}^{cd}$$

\uparrow relabel $c \leftrightarrow d$

$$\Rightarrow C(\text{adj}) = \frac{\delta^{ab} f_{acd} f_b{}^{cd}}{\dim g}$$

Ex 28

By considering infinitesimal gauge transformations ($|\alpha| \ll 1$)

$$g = e^{i\alpha^a t_a} \equiv e^{i\alpha} = \mathbb{1} + i\alpha + \mathcal{O}(\alpha^2)$$

and Taylor expanding finite gauge transformations to leading order in $\alpha \in \mathfrak{g} = \text{lie}(G)$, show that the infinitesimal gauge variations of the fields are

$$\delta_\alpha \phi = i\alpha \phi$$

$$\delta_\alpha A_\mu = i[\alpha, A_\mu] + \partial_\mu \alpha$$

$$\delta_\alpha F_{\mu\nu} = i[\alpha, F_{\mu\nu}]$$

where $\phi \mapsto \phi + \delta_\alpha \phi + \mathcal{O}(\alpha^2)$ and so on.

Finite gauge transformations:

$$\phi \mapsto \phi' = g \cdot \phi = e^{i\alpha} \phi$$

$$\begin{aligned} A_\mu \mapsto A'_\mu &= g A_\mu g^{-1} + i g (\partial_\mu g^{-1}) \\ &= e^{i\alpha} A_\mu e^{-i\alpha} + i e^{i\alpha} (\partial_\mu e^{-i\alpha}) \end{aligned}$$

$$F_{\mu\nu} \mapsto F'_{\mu\nu} = g F_{\mu\nu} g^{-1} = e^{i\alpha} F_{\mu\nu} e^{-i\alpha}$$

Taylor expand to first order:

$$\phi \mapsto (1 + i\alpha + \mathcal{O}(\alpha^2)) \phi = \phi + i\alpha \phi + \mathcal{O}(\alpha^2)$$

$$\Rightarrow \delta_\alpha \phi = i\alpha \phi$$

$$\begin{aligned} A_\mu \mapsto & (1 + i\alpha + \mathcal{O}(\alpha^2)) A_\mu (1 - i\alpha + \mathcal{O}(\alpha^2)) + i (1 + i\alpha + \mathcal{O}(\alpha^2)) (\partial_\mu (1 - i\alpha + \mathcal{O}(\alpha^2))) \\ &= A_\mu + i\alpha A_\mu - i A_\mu \alpha + i(-i) (\partial_\mu \alpha) + \mathcal{O}(\alpha^2) \\ &= A_\mu + i[\alpha, A_\mu] + (\partial_\mu \alpha) + \mathcal{O}(\alpha^2) \end{aligned}$$

$$\Rightarrow \delta_\alpha A_\mu = i[\alpha, A_\mu] + \partial_\mu \alpha. \quad \left(\begin{array}{l} \text{Similarly,} \\ \delta_\alpha F_{\mu\nu} = i[\alpha, F_{\mu\nu}] \end{array} \right)$$