

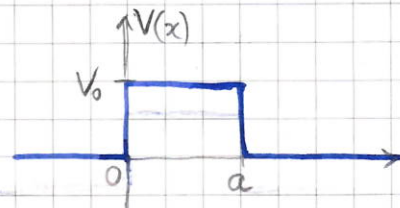
Ex 53

TISE

$$\left(-\frac{d^2}{dx^2} + V(x)\right)\psi(x) = k^2\psi(x)$$

SQUARE BARRIER/WELL POTENTIAL

$$V(x) = \begin{cases} 0, & x < 0 \\ V_0, & 0 < x < a \\ 0, & x > a \end{cases}$$



1. Show that the matching conditions at $x=0, a$, where $V(x)$ is discontinuous (but finite) are that $\psi(x), \psi'(x)$ are continuous.

This can be proven in generality for a potential that is finite in a neighbourhood of x_0 . Consider the infinitesimal neighbourhood $(x_0 - \epsilon, x_0 + \epsilon)$, with $\epsilon > 0$. Integrate the TISE once and twice:

ONCE:
$$-\left[\psi'(x)\right]_{x_0-\epsilon}^{x_0+\epsilon} + \int_{x_0-\epsilon}^{x_0+\epsilon} dx V(x)\psi(x) = k^2 \int_{x_0-\epsilon}^{x_0+\epsilon} dx \psi(x)$$

ψ, V bounded near $x_0 \Rightarrow$ integrals $\xrightarrow{\epsilon \rightarrow 0^+} 0$.

$\Rightarrow \psi'(x)$ continuous at x_0 .

TWICE:
$$-\left[\psi(x)\right]_{x_0-\epsilon}^{x_0+\epsilon} + 2\epsilon \psi'(x_0 - \epsilon) + \int_{x_0-\epsilon}^{x_0+\epsilon} dx \int_{x_0-\epsilon}^x dx' (V(x') - k^2)\psi(x') = 0$$

$\int_{x_0-\epsilon}^{x_0+\epsilon} dx \int_{x_0-\epsilon}^x dx'$
 $\int_{x_0-\epsilon}^{x_0+\epsilon} dx$
 $\int_{x_0-\epsilon}^{x_0+\epsilon} dx'$

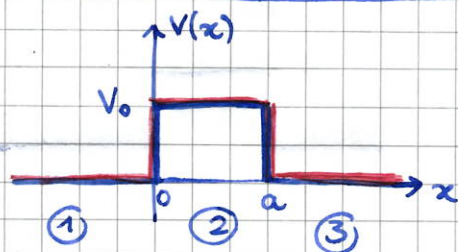
$\downarrow \epsilon \rightarrow 0$
 $\downarrow \epsilon \rightarrow 0$
 $\downarrow \epsilon \rightarrow 0$

0
 0
 0

$\Rightarrow \psi(x)$ continuous at x_0 .

(YOU CAN QUOTE THIS RESULT IN THE EXAM IF USEFUL.)

2. Solve the SE in the three regions and impose the matching conditions to find the scattering solutions associated to $k^2 > 0$ (continuous spectrum), and determine $R(k), T(k)$ in terms of a & $l = \sqrt{k^2 - V_0}$.



①, ③ : $-\frac{d^2}{dx^2} \psi = k^2 \psi \Rightarrow$ solutions $e^{\pm ikx}$

② : $-\frac{d^2}{dx^2} \psi = (k^2 - V_0) \psi = l^2 \psi \Rightarrow$ " $e^{\pm ilx}$ ($l^2 = k^2 - V_0$)

Scattering solution:

$$\psi(x) = \begin{cases} e^{ikx} + R(k)e^{-ikx} & , x < 0 & \text{①} \\ A^{(l)}e^{ilx} + B^{(l)}e^{-ilx} & , 0 < x < a & \text{②} \\ T(k)e^{ikx} & , x > a & \text{③} \end{cases}$$

4 unknowns. Impose 2+2 matching conditions at $x=0, a$ to find unique scattering solution with unit flux of incident waves of wavenumber $k > 0$.

ψ continuous at $x=0$: $1 + R = A + B$ ①

" $x=a$: $Ae^{ila} + Be^{-ila} = Te^{ika}$ ②

ψ' continuous at $x=0$: $ik(1-R) = il(A-B)$ ③

" $x=a$: $il(Ae^{ila} - Be^{-ila}) = ikTe^{ika}$ ④

First eliminate A & B . Since the RHS's of ② & ④ are simpler, I'll use them to eliminate A & B .

$$\textcircled{b} + \frac{\textcircled{d}}{il} : 2A e^{ila} = \left(1 + \frac{k}{l}\right) T e^{ika}$$

$$\textcircled{b} - \frac{\textcircled{d}}{il} : 2B e^{-ila} = \left(1 - \frac{k}{l}\right) T e^{ika}$$

$$\Rightarrow A = \frac{1}{2} \left(1 + \frac{k}{l}\right) e^{i(k-l)a} T, \quad B = \frac{1}{2} \left(1 - \frac{k}{l}\right) e^{i(k+l)a} T$$

Sub in \textcircled{a} & \textcircled{c} :

$$\textcircled{a} : 1 + R = \frac{1}{2} \left[\left(1 + \frac{k}{l}\right) e^{-ila} + \left(1 - \frac{k}{l}\right) e^{ila} \right] e^{ika} T$$

$$\frac{\textcircled{c}}{ik} : 1 - R = \frac{1}{2} \frac{l}{k} \left[\left(1 + \frac{k}{l}\right) e^{-ila} - \left(1 - \frac{k}{l}\right) e^{ila} \right] e^{ika} T$$

$$\begin{aligned} \Rightarrow \textcircled{a} + \frac{\textcircled{c}}{ik} : 2 &= \frac{1}{2} e^{ika} T \left[\left(1 + \frac{l}{k}\right) \left(1 + \frac{k}{l}\right) e^{-ila} + \left(1 - \frac{k}{l}\right) \left(1 - \frac{l}{k}\right) e^{ila} \right] \\ &= \frac{1}{2} e^{ika} T \left[\frac{(k+l)^2}{kl} e^{-ila} - \frac{(k-l)^2}{kl} e^{ila} \right] \end{aligned}$$

$$\Rightarrow T(k) = \frac{4kl e^{-ika}}{(k+l)^2 e^{-ila} - (k-l)^2 e^{ila}}$$

$(l = \sqrt{k^2 - V_0})$

$$= \frac{4kl e^{-ika}}{-2i(k^2 + l^2) \sin(la) + 4kl \cos(la)}$$

$$\begin{aligned} \textcircled{a} - \frac{\textcircled{c}}{ik} : 2R &= \frac{1}{2} e^{ika} T \left[\left(1 - \frac{l}{k}\right) \left(1 + \frac{k}{l}\right) e^{-ila} + \left(1 - \frac{k}{l}\right) \left(1 + \frac{l}{k}\right) e^{ila} \right] \\ &= \frac{e^{ika} T}{2lk} \left[(k^2 - l^2) e^{-ila} - (k^2 - l^2) e^{ila} \right] \\ &= \frac{e^{ika} T}{2lk} 2i \sin(la) (l^2 - k^2) = -\sin(la) \cdot i \frac{k^2 - l^2}{lk} e^{ika} T \end{aligned}$$

$$\Rightarrow R(k) = -i \sin(la) \frac{k^2 - l^2}{2lk} e^{ika} T(k) = \frac{-2i(k^2 - l^2) \sin(la)}{-2i(k^2 + l^2) \sin(la) + 4kl \cos(la)}$$

3. For which values of k is the square potential "transparent", i.e. $R(k) = 0$?

$$\underline{R(k) = 0} : \quad k^2 - l^2 = V_0 = 0 \quad (\text{no potential - trivial})$$

or

$$\sin(la) = 0 \Leftrightarrow la = n\pi, \quad n \in \mathbb{Z}$$

$$\Leftrightarrow k^2 - V_0 = l^2 = \left(\frac{n\pi}{a}\right)^2$$

$$\Leftrightarrow k^2 = V_0 + \left(\frac{n\pi}{a}\right)^2$$

For these values of k the incident wave is transmitted but not reflected:

$$R(k) = 0, \quad T(k) = \frac{e^{-ika}}{\cos(n\pi)} = (-1)^n e^{-ika}$$

4. Write down the bound state solutions corresponding to the discrete spectrum $k^2 = -\mu^2 < 0$. Find the equations that determine implicitly the allowed values of μ in terms of l (or V_0).

TO WRITE DOWN THE BOUND STATE SOLUTIONS (I won't do it explicitly):

1) Sub $k = i\mu$ in scattering solution, $\mu > 0$.

2) Divide through by $T(i\mu)$.

3) Choose $k = i\mu$ s.t. $\frac{1}{T(i\mu)} = 0$ to get rid of the unbounded $e^{\mu x}$ term for $x \rightarrow -\infty$.

So look for poles of $T(k)$ on the positive imaginary axis.

• Poles of $T(k)$: $(k+l)^2 e^{-ila} = (k-l)^2 e^{ila}$

$$\Rightarrow \frac{k+l}{k-l} = \varepsilon e^{ila}, \quad \varepsilon = \pm 1$$

$$\Rightarrow k(1 - \varepsilon e^{ila}) = -l(1 + \varepsilon e^{ila})$$

$$\Rightarrow k = -l \frac{1 + \varepsilon e^{ila}}{1 - \varepsilon e^{ila}} = l \frac{e^{ila/2} + \varepsilon}{e^{ila/2} - \varepsilon} = l \frac{e^{ila/2} + \varepsilon e^{-ila/2}}{e^{ila/2} - \varepsilon e^{-ila/2}}$$

$$\varepsilon = -1 : k = i l \tan\left(\frac{la}{2}\right), \quad \varepsilon = +1 : k = -i l \cot\left(\frac{la}{2}\right)$$

$$\uparrow \qquad \qquad \qquad \uparrow$$

$$\mu = l \tan\left(\frac{la}{2}\right) \qquad \qquad \qquad \mu = -l \cot\left(\frac{la}{2}\right)$$

But we also have

$$l^2 = k^2 - V_0 \quad \Leftrightarrow \quad l^2 + \mu^2 = -V_0$$

So we need

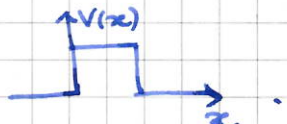
$$\begin{cases} \mu = l \tan\left(\frac{la}{2}\right) & \vee & \mu = -l \cot\left(\frac{la}{2}\right) \\ l^2 + \mu^2 = -V_0 \end{cases} \quad (*)$$

5. Do bound state solutions exist for $V_0 > 0$? And for $V_0 < 0$? In the latter case, use a graphical argument to show that a new bound state solution appears every time $\sqrt{-V_0}$ crosses a non-negative multiple of $\frac{\pi}{4a}$.

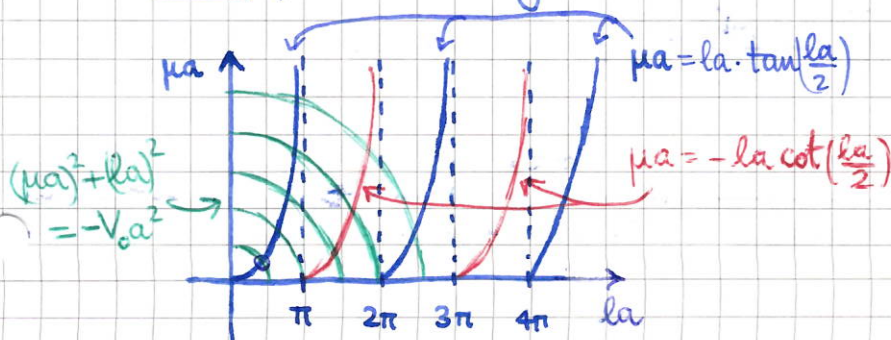
We seek solutions with $\mu > 0$.

• If $V_0 > 0$, need $l^2 < 0$. Set $l = iL$, ($L > 0$ wlog) and get

$$\mu = -L \tanh\left(\frac{aL}{2}\right) \leq 0 \quad \text{or} \quad \mu = -L \coth\left(\frac{aL}{2}\right) < 0$$

\Rightarrow No bound state solutions, as expected since 

• If $V_0 < 0$, we can again exclude $l^2 < 0$ as above. For $l, \mu > 0$



Bound state solutions exist in correspondence of the intersections of the green $\frac{1}{4}$ -circles with the blue or red curves.

• Since $la \cdot \tan\left(\frac{la}{2}\right) \approx \frac{(la)^2}{2}$ for $la \ll 1$, there is always at least a bound state solution if $V_0 < 0$.

• A new bound state solution appears every time

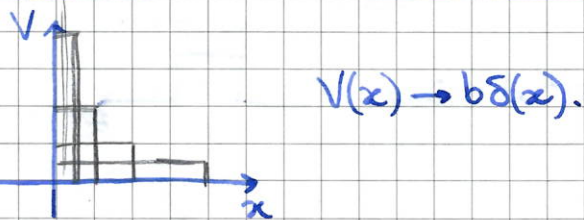
$$\underbrace{\sqrt{-V_0} \cdot a}_{\text{Radius of the circle}} = n\pi, \quad n=0, 1, 2, \dots$$

$$\Leftrightarrow \underline{\underline{\sqrt{-V_0} = \frac{n\pi}{a}}}, \quad n=0, 1, 2, \dots$$

The eigenvalues k^2 in the discrete spectrum move down as $-V_0$ increases, and whenever $\sqrt{-V_0}$ crosses an integer multiple of $\frac{\pi}{a}$ a new bound state solution appears with $k^2 = 0$.



6. Show that when $a \rightarrow 0$, $V_0 \rightarrow \infty$ with $aV_0 = b$ fixed, the reflection and transmission coeff's reduce to those for $V(x) = b\delta(x)$.



$$V_0 = \frac{b}{a} \rightarrow \infty, \quad \Rightarrow \quad l = \sqrt{k^2 - V_0} \approx \frac{\sqrt{-b}}{a} \rightarrow \infty \Rightarrow la \rightarrow 0.$$

$$T(k) = \frac{4kl e^{-ika}}{-2i(k^2 + l^2)\sin(la) + 4kl \cos(la)} \approx \frac{4kl}{-2il^2 a + 4kl} \xrightarrow{l^2 a \rightarrow -b} \frac{2k}{2k + ib} = \frac{2ik}{2ik - b}$$

$$R(k) = -i \sin(la) \frac{k^2 - l^2}{2lk} e^{ika} T(k) \approx \frac{-l^2 a}{2ik} \frac{2ik}{2ik - b} \rightarrow \frac{b}{2ik - b}$$