## Assignment 4

## Due date: Friday, 2 December (8pm)

## Ex 25

This is yet another question about the KdV equation $u_{t}+6 u u_{x}+u_{x x x}=0$. This time we will focus on conservation laws and what they can teach us about the time evolution of the class of initial conditions that we experimented with in the first lecture.

1. Evaluate the first three KdV conserved charges

$$
\begin{equation*}
Q_{1}=\int_{-\infty}^{+\infty} d x u, \quad Q_{2}=\int_{-\infty}^{+\infty} d x u^{2}, \quad Q_{3}=\int_{-\infty}^{+\infty} d x\left(u^{3}-\frac{1}{2} u_{x}^{2}\right) \tag{4.1}
\end{equation*}
$$

for the initial state $u(x, 0)=A \operatorname{sech}^{2}(B x)$, where $A$ and $B$ are constants. [20 marks]

## SOLUTION:

I will use the definite integrals (where $n \in \mathbb{N}$ )

$$
\begin{gathered}
I_{n}:=\int_{-\infty}^{+\infty} d y \operatorname{sech}^{2 n}(y)=\frac{2^{2 n-1}((n-1)!)^{2}}{(2 n-1)!} \\
\Longrightarrow I_{1}=2, \quad I_{2}=\frac{4}{3}, \quad I_{3}=\frac{16}{15}, \quad I_{4}=\frac{32}{35}, \quad I_{5}=\frac{256}{315}, \quad \ldots
\end{gathered}
$$

which are tabulated at the end of the problem sheet, as well as the derivative formula

$$
\frac{d}{d y} \operatorname{sech} y=-\operatorname{sech} y \cdot \tanh y=-\operatorname{sech}^{2} y \cdot \sinh y
$$

which implies

$$
\begin{aligned}
\frac{d}{d y} \operatorname{sech}^{2} y & =-2 \operatorname{sech}^{3} y \cdot \sinh y \\
\left(\frac{d}{d y} \operatorname{sech}^{2} y\right)^{2} & =4 \operatorname{sech}^{6} y \cdot \sinh ^{2} y=4 \operatorname{sech}^{6} y \cdot\left(\cosh ^{2} y-1\right)=4\left(\operatorname{sech}^{4} y-\operatorname{sech}^{6} y\right)
\end{aligned}
$$

We can then evaluate $Q_{1}, Q_{2}, Q_{3}$ at $t=0$ by changing integration variable $y=B x$ :

$$
\begin{aligned}
Q_{1} & =\frac{A}{B} I_{1}=\frac{2 A}{B} \\
Q_{2} & =\frac{A^{2}}{B} I_{2}=\frac{4 A^{2}}{3 B} \\
Q_{3} & =\frac{A^{3}}{B} I_{3}-2 A^{2} B\left(I_{2}-I_{3}\right)=\frac{8 A^{2}\left(2 A-B^{2}\right)}{15 B} .
\end{aligned}
$$

2. The initial condition

$$
\begin{equation*}
u(x, 0)=N(N+1) \operatorname{sech}^{2}(x), \tag{4.2}
\end{equation*}
$$

where $N$ is an integer, is known to evolve at late times into $N$ well-separated solitons, with velocities $4 k^{2}, k=1 \ldots N$. So for $t \rightarrow+\infty$, this solution approaches the sum of $N$ single well-separated solitons

$$
\begin{equation*}
u(x, t) \approx \sum_{k=1}^{N} 2 k^{2} \operatorname{sech}^{2}\left[k\left(x-x_{k}-4 k^{2} t\right)\right] \tag{4.3}
\end{equation*}
$$

Since $Q_{1}, Q_{2}$ and $Q_{3}$ are conserved, their values at $t=0$ and $t \rightarrow+\infty$ must be equal. Use this fact to deduce formulae for the sums of the first $N$ integers, the first $N$ cubes, and the first $N$ fifth powers.

## SOLUTION:

To calculate the conserved charges $Q_{1}, Q_{2}, Q_{3}$ at the initial time $t=0$, we can just set $A=N(N+1)$ and $B=1$ in the results of part 1 :

$$
\begin{aligned}
& Q_{1}=2 N(N+1) \\
& Q_{2}=\frac{4}{3} N^{2}(N+1)^{2} \\
& Q_{3}=\frac{8}{15} N^{2}(N+1)^{2}\left(2 N^{2}+2 N-1\right)
\end{aligned}
$$

To calculate the conserved charges $Q_{1}, Q_{2}, Q_{3}$ at late times $t \rightarrow+\infty$, we use the fact that the $N$ solitons are well-separated, that is, they are separated by much larger distances than the widths of the solitons. So

$$
u \approx \sum_{k=1}^{N} u_{k, x_{k}}, \quad u^{2} \approx \sum_{k=1}^{N} u_{k, x_{k}}^{2}, \quad u^{3} \approx \sum_{k=1}^{N} u_{k, x_{k}}^{3}, \quad u_{x}^{2} \approx \sum_{k=1}^{N}\left(\partial_{x} u_{k, x_{k}}\right)^{2}
$$

where

$$
u_{k, x_{k}}(x, t)=2 k^{2} \operatorname{sech}^{2}\left[k\left(x-x_{k}-4 k^{2} t\right)\right] .
$$

The cross terms are negligible because the soliton solutions $u_{k, x_{k}}$ tend to zero exponentially fast away from their centres, and the solitons are assumed to be well-separated.

So the contributions of the $N$ well-separated solitons simply add up. Comparing with part 1, the $k$-th soliton $u_{k, x_{k}}$ has $A_{k}=2 k^{2}$ and $B_{k}=k$ (the shift of $x$ by $x_{k}+4 k^{2} t$ is inconsequential for the calculation of the charges as it can be absorbed by a shift of the integration variable). So we find the charges

$$
\begin{aligned}
Q_{1} & =2 \sum_{k=1}^{N} \frac{A_{k}}{B_{k}}=4 \sum_{k=1}^{N} k \\
Q_{2} & =\frac{4}{3} \sum_{k=1}^{N} \frac{A_{k}^{2}}{B_{k}}=\frac{16}{3} \sum_{k=1}^{N} k^{3} \\
Q_{3} & =\frac{8}{15} \sum_{k=1}^{N} \frac{A_{k}^{2}\left(2 A_{k}-B_{k}^{2}\right)}{B_{k}}=\frac{32}{5} \sum_{k=1}^{N} k^{5} .
\end{aligned}
$$

Equating the $t=0$ expressions and the $t \rightarrow+\infty$ expressions for the conserved charges we find

$$
\sum_{k=1}^{N} k=\frac{1}{2} N(N+1), \quad \sum_{k=1}^{N} k^{3}=\frac{1}{4} N^{2}(N+1)^{2}, \quad \sum_{k=1}^{N} k^{5}=\frac{1}{12} N^{2}(N+1)^{2}\left(2 N^{2}+2 N-1\right) .
$$

## Ex 26

1. Show that the pair of equations

$$
\begin{align*}
& (u-v)_{+}=\sqrt{2} e^{(u+v) / 2}  \tag{5.1}\\
& (u+v)_{-}=\sqrt{2} e^{(u-v) / 2}
\end{align*}
$$

provides a Bäcklund transformation linking solutions of $v_{+-}=0$ (the wave equation in light-cone coordinates) to those of $u_{+-}=e^{u}$ (the Liouville equation).
[15 marks]

## SOLUTION:

Cross-differentiate (5.1) to get

$$
\begin{aligned}
& (u-v)_{+-}=\frac{1}{\sqrt{2}} e^{(u+v) / 2}(u+v)_{-}=e^{(u+v) / 2} e^{(u-v) / 2}=e^{u} \\
& (u+v)_{-+}=\frac{1}{\sqrt{2}} e^{(u-v) / 2}(u-v)_{+}=e^{(u-v) / 2} e^{(u+v) / 2}=e^{u}
\end{aligned}
$$

Taking sum and difference we obtain

$$
\begin{array}{ll}
u_{+-}=e^{u} & \text { (Liouville eqn) } \\
v_{+-}=0 & \text { (the wave eqn) }
\end{array}
$$

2. Starting from d'Alembert's general solution $v=f\left(x^{+}\right)+g\left(x^{-}\right)$of the wave equation, use the Bäcklund transform (5.1) to obtain the corresponding solutions of the Liouville equation for $u$.
[Hint: Set $u\left(x^{+}, x^{-}\right)=2 U\left(x^{+}, x^{-}\right)+f\left(x^{+}\right)-g\left(x^{-}\right)$. You might simplify the notation by setting $f\left(x^{+}\right)=\log \left(F^{\prime}\left(x^{+}\right)\right)$and $g\left(x^{-}\right)=-\log \left(G^{\prime}\left(x^{-}\right)\right.$, where prime means first derivative.]
[35 marks]

## SOLUTION:

Substituting the general solution $v=f\left(x^{+}\right)+g\left(x^{-}\right)$of the wave equation in the Bäcklund transform we obtain a system of two first order PDEs for a solution $u$ of the Liouville equation:

$$
\left\{\begin{array}{l}
u_{+}-f^{\prime}\left(x^{+}\right)=\sqrt{2} e^{\left[u+f\left(x^{+}\right)+g\left(x^{-}\right)\right] / 2} \\
u_{-}+g^{\prime}\left(x^{-}\right)=\sqrt{2} e^{\left[u-f\left(x^{+}\right)-g\left(x^{-}\right)\right] / 2}
\end{array}\right.
$$

The system simplifies if we make the substitution $u\left(x^{+}, x^{-}\right)=2 U\left(x^{+}, x^{-}\right)+f\left(x^{+}\right)-$ $g\left(x^{-}\right)$given in the hint:

$$
\left\{\begin{array}{l}
2 U_{+}=\sqrt{2} e^{U+f\left(x^{+}\right)} \\
2 U_{-}=\sqrt{2} e^{U-g\left(x^{-}\right)}
\end{array} \Longleftrightarrow \quad \Longleftrightarrow \begin{array}{l}
e^{-U} U_{+}=\frac{1}{\sqrt{2}} e^{f\left(x^{+}\right)} \equiv \frac{1}{\sqrt{2}} F^{\prime}\left(x^{+}\right) \\
e^{-U} U_{-}=\frac{1}{\sqrt{2}} e^{-g\left(x^{-}\right)} \equiv \frac{1}{\sqrt{2}} G^{\prime}\left(x^{-}\right)
\end{array}\right.
$$

where in the last expression we used $f\left(x^{+}\right)=\log \left(F^{\prime}\left(x^{+}\right)\right)$and $g\left(x^{-}\right)=-\log \left(G^{\prime}\left(x^{-}\right)\right)$ as suggested. This system can be integrated to get

$$
-e^{-U}=\frac{1}{\sqrt{2}}\left(F\left(x^{+}\right)+G\left(x^{-}\right)+c\right)
$$

where $c$ is an integration constant. Taking the logarithm of the previous equation,

$$
U=-\log \left[-\frac{1}{\sqrt{2}}\left(F\left(x^{+}\right)+G\left(x^{-}\right)+c\right)\right]
$$

so

$$
\begin{aligned}
u & =2 U+\log \left[F^{\prime}\left(x^{+}\right) G^{\prime}\left(x^{-}\right)\right] \\
& =-2 \log \left[-\frac{1}{\sqrt{2}}\left(F\left(x^{+}\right)+G\left(x^{-}\right)+c\right)\right]+\log \left[F^{\prime}\left(x^{+}\right) G^{\prime}\left(x^{-}\right)\right] \\
& =\log \frac{2 F^{\prime}\left(x^{+}\right) G^{\prime}\left(x^{-}\right)}{\left(F\left(x^{+}\right)+G\left(x^{-}\right)+c\right)^{2}} .
\end{aligned}
$$

