

# How good is the bipolar approximation of active regions for surface flux transport?

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#### **Motivation**

 The surface flux transport (SFT) model is remarkably effective at mimicking the large-scale decay of active region magnetic fields.
 [cf. Jiang et al., Space Sci Rev 2014;

Wang, Space Sci Rev 2017]

- Current interest:
  - Making solar cycle predictions before the end of the previous cycle.
  - Driving continuous simulations of the coronal magnetic field.





What is lost by the traditional assumption of symmetric bipolar magnetic sources?

#### Leading/following asymmetry

- Following polarities are often more diffuse.
- Iijima-Hotta-Imada 2019 -SFT simulations to investigate the effect (tilt angle proportional to latitude, all same asymmetry).

(a) Effect of tilt angle



(b) Effect of size asymmetry



 Accounting for asymmetry weakens dipole and gives better reversal time.

> area ratio 1.0 area ratio 0.4



#### Complexity

• Jiang et al. 2019 -

case studies of SFT evolution for two active regions.



<B> [G]

0.00

0.30

AR 12674

-0.30

50

#### My work

- Automated database of Bipolar Magnetic Regions from HMI/SHARPs.
- Compare SFT models with BMRs vs original SHARPs.

**A.R. Yeates**, How good is the bipolar approximation of active regions for surface flux transport?, *Solar Physics* **295**, 119 (2020)

Python code for extracting database: <u>https://github.com/antyeates1983/sharps-bmrs</u>

Ready-prepared file: <u>https://doi.org/10.7910/DVN/1Z7YMT</u> (Harvard Dataverse) for May 2010 to April 2020:

<pre>bmrsharps_evol.txt SHARPs from 2018-05-01 00:00:00 to 2028-04-06 00:00:00 Produced by anthany_weates.latidurban.ac.uk 1090 Grid resolution: 180 x 360, smoothing_param = 4 Selection criteria: [1] sep &gt; 1 deg, (11)  imbalance  &lt; 0.5 Last two columns use 10-year 10 SFT simulation with eta=350 km^2/s, v0=0.015 km/s, p=2.33, no decay term.</pre>															
l	SHARP	NO4A	CM Line	Latitude	Carr-Longitude	Unsond flux	Inbalance	Dicole	Bio-Separation	Bip-Till	Bip-Dipole	Pred-Dip-Real	Pred-Cip-Bip		
l	1	11067	2012-05-07	24.58757	172,95456	3.92405e+21	0.39477	4.60159e-04	5.37126e+80	3,20301e+00	4.527498-04	4.55467e-06	9.21374e-86		
l	2	11064	2012-05-03	12.12164	223,73779	1.35816c+21	-0.25154	-1.18855c-83	5.23233c+80	-2,35854c+01	-1.17753e-03	-1.10364c-03	-1.52803c-03		
l	10	11056	2012-05-04	-25.16836	286.65967	1.76575e+21	0.22983	2.84643e-04	2.82886e+80	1.73911e+02	2.05394e-04	2.34538e-05	1.92884e-86		
l	12	11058	2012-05-10	-19.89266	134.14394	3.52810e+21	0.31224	-1.20625e-23	4.11376e+20	-1.68141e+02	-1.19192e-03	-8.52827e-85	-1.33504e-04		
l	14	11070	2012-05-05	20.25744	194.78160	5.00410e+20	-0.03259	-1.69289e-84	1.22874e+30	-3.46432e+01	-1.86002e-04	-1.52582e-85	-1.68276e-05		
l	26	11072	2018-05-23	-15.21733	316.68772	6.99753e+21	-0.03534	1.53796e-03	3.65048e+20	1.71658e+02	1.53320e-03	6.97488e-04	8.00380e-84		
l	38	11073	2018-06-02	12.88378	194,19472	2.22336e+21	-0.89995	-1.87460c-03	3.53620c+80	-3,46612e+01	-1.86447c-03	-1.87190e-03	-1.91126c-03		
l	40	11075	2018-05-30	-19.66521	230.78519	1.16558e+21	0.81622	5.13498e-04	2.26383e+80	1.51270e+02	5.09108e-04	5.84839e-05	5.83880e-85		
l	43	11076	2012-06-01	-13.66200	196.18004	2.08237e+21	-0.11159	1.28108e-04	2.37732e+20	1.76333e+02	1.23387e-04	1.95448e-05	1.41248e-25		
l	44	0	2013-06-02	-33.12716	184.92852	5.97273e+20	-0.15399	3.26955e-05	1.46385e+80	1.74403e+02	2.84019e-05	1.53281e-08	1.42752e-38		
l	47	0	2012-06-10	15.19621	88.08352	7.74446c+20	0.20963	4.37134c-04	2.37736c120	3.48945c+01	4.35103e-04	2.21568e-04	2.23777e-84		
l	57	11082	2018-06-20	28.39895	383.46702	4.48532e+21	0.20622	4.19298e-03	3.60712e+20	4.32266e+01	4.15153e-03	1.52768e-05	1.67351e-85		
l	56	0	2012-06-28	15.71528	199.82528	2.13414e+21	-0.02659	-2.28356e-23	3.47132e+20	-4.81243e+01	-2.26704e-03	-1.00413e-23	-1.03158e-03		
l	67	11085	2013-06-28	-21.91311	230.43488	1.52761e+21	0.12114	-1.72414e-23	3.67840e+30	-1.29731e+02	-1.71825e-03	-9.65460e-85	-9.30580e-05		
l	36	11087	2012-07-15	19.49559	335.55594	1.83895e+22	-0.03850	1.92913e-02	8.97116e+30	1.67519e+01	1.98763e-02	4.51918e-03	3.39153e-83		
l	87	0	2012-07-09	19.52667	120.34961	3,89123c+20	0.31597	8.57507c-05	1.09127c+20	2.98810c+01	1.07181e-04	1.02349e-05	1.26450c-85		
l	89	11088	2018-07-15	-28.59370	337.44025	7.32340e+20	-0.40323	-4.75846e-85	2.18557e+80	-1.75765e+02	-5.06708e-05	-3.01310e-86	-4.10831e-06		
l	92	11089	2012-07-25	-22.90278	283.23528	2.05991e+22	-0.02122	9.70493e-03	5.68659e+20	1.67930e+02	9.61091e-03	3.75159e-04	3.86887e-24		
	97	0	2013-07-24	13.34482	220.46158	3.58470e+20	0.37365	-2.40031e-35	1.31483e+80	-6.98935e+00	-3.34932e-05	-1.63547e-85	-2.89587e-05		
	38	11090	2018-07-29	22.49662	149,96277	5.08242c+20	-0.32211	5.19955c-04	5.30587c+80	2.90281c+01	5.15122e-04	2.10756e-05	2.42236e-85		
	164	11032	2012-08-04	15.09601	75.61377	1.30375++22	-0.14607	3,233456-02	9.35438e+20	3.96531e+01	3.1976202	1.51923e-02	1.88848e - 22		

## **Building the Database**

#### **SHARP data**

I use hmi.sharp\_cea\_720s — "definitive" data (after full-disk package).



#### **Magnetogram extraction**

Single observation for each SHARP — when flux-weighted centroid closest to CM.



#### **Bipolar approximation**

Fit based on polarity centroids:

$$B(s,\phi) = -B_0 rac{\phi}{
ho} \exp\left[-rac{\phi^2 + 2 \arcsin^2(s)}{\left(a
ho
ight)^2}
ight]$$

- Rotate to correct location and tilt angle.
- Scale unsigned flux to match.



Parameter a controls dipole moment for given flux. Set a = 0.56 so axial dipole moment of BMR matches that of original SHARP.



#### **Filtering**

Initially: 3671 regions, total flux 1.4e25 Mx

- 1. Remove SHARPs with too much flux imbalance. [2323 regions, total flux 0.18e25 Mx]
- 2. Remove SHARPs where separation too small to resolve on computational grid. [114 regions, total flux 0.004e25 Mx]
- 3. Remove repeat observations. [143 regions, total flux 0.17e25 Mx]

Remaining: 1090 regions, total flux 1.0e25 Mx.

#### **Automated removal of repeats**

- Every SHARP is compared with those passing CM between 20-34 days earlier.
- A "repeat" is where an earlier SHARP had more flux in its derotated footprint.



#### **Summary of emergence-time properties**



Database includes predicted final dipole moment from Surface Flux Transport of (a) the BMR and (b) the original SHARP.

## **Evolution**

#### Surface flux transport model

diffusion

To compute the axial dipole moment we only need to evolve the 1D longitudeaveraged field:

$$\frac{\partial \overline{B}}{\partial t} = \frac{\partial}{\partial s} \left[ \frac{D}{R_{\odot}^2} (1 - s^2) \frac{\partial \overline{B}}{\partial s} - \frac{v_s(s)}{R_{\odot}} \sqrt{1 - s^2} \overline{B} \right] \qquad s = \cos \theta$$
supergranular
meridional flow

 $v_s(s) = D_u s(1-s^2)^{p/2} \quad \left[ = D_u \cos \theta \sin^p \theta \right]$ 

Parameters set using full simulation [all SHARPs, initialised with smoothed HMI synoptic map].

 $D = 350 \,\mathrm{km}^2 \mathrm{s}^{-1}$   $D_u = 0.041 \,\mathrm{km} \,\mathrm{s}^{-1}$  p = 2.33

[same profile as Whitbread-Y-Muñoz-J 2018, but faster flow and slower diffusion]

No exponential decay term [didn't seem to be needed to match observed evolution].

#### **Complete simulation**



### **Complete simulation**

[1D results]

BMRs overestimate dipole moment by 24%.



#### **Individual evolutions**

 Solve with finite-differences for each region individually [10 years to steady state].



#### **Dipole amplification**

- Ratio of final to initial dipole moment [Jiang-Cameron-Schüssler 2014].
- As explained by Petrovay-Nagy-Yeates 2020 this follows a Gaussian in latitude for BMR sources.
- But for SHARPs not all regions fall on the curve:



- There are more SHARPs below the curve than above => weaker dipole.
- These are "non-dipolar" regions with enhanced cancellation.

#### **Regions with the largest discrepancy**

following polarity more diffuse [cf. lijima-Hotta-Imada 2019]



more complex initial shape [cf. Jiang et al 2019]

#### Reflections

Suggests that active region inflows are not the cause of the overestimate of dipole from BMRs.

[inflows may be accounted for already with my faster meridional flow]

- If SHARPs are used, dipole reverses at correct time even without decay term. [cf. Petrovay-Talafha 2019]
- Reducing the tilt angles has the same effect as slowing down the meridional flow:



# Driving Coronal Simulations [work in progress]

#### Magnetofrictional model

 Quasi-static evolution but preserving magnetic topology with induction equation [van Ballegooijen-Priest-Mackay 2000; Yeates 2014].

$$oldsymbol{B}=
abla imes oldsymbol{A}$$

$$\frac{\partial A}{\partial t} = \mathbf{v} \times \mathbf{B} - \eta \nabla \times \mathbf{B}$$
$$\mathbf{v} = \frac{(\nabla \times \mathbf{B}) \times \mathbf{B}}{\nu(B^2)} + \mathbf{v}_{\text{out}} \left(\frac{r}{R_{ss}}\right)^{11.5} \hat{\mathbf{r}}$$

- Boundary conditions:
  - zero gradient on  $r = R_{ss} = 2.5 R_{\odot}$

• on 
$$r = R_{\odot}$$
 set  $\frac{\partial A_h}{\partial t} = \mathbf{v}_h \times (B_r \hat{\mathbf{r}}) - D\nabla \times (B_r \hat{\mathbf{r}}) + \mathbf{E}_{em}$   
differential supergranular emergence rotation diffusion

 Previous simulations were driven by emergence of analytical bipoles [e.g. Yeates-Mackay-van Ballegooijen 2008].



#### How to emerge SHARPs

1. Compute local **A** on photosphere for the new region.



Local inductive" method:

minimize  $\sum_{D} |\mathbf{A}|^2$  subject to  $\hat{\mathbf{r}} \cdot \nabla \times \mathbf{A} = B_r$  with boundary condition  $\mathbf{A} \times \mathbf{n}|_{\partial D} = \mathbf{0}$ .

• Jumps in **A**.**n** are then removed by applying a "curl-free smoothing"  $\frac{\partial A}{\partial t} = \nabla (\nabla \cdot A)$ .

#### Why is the localization necessary?

Avoids spurious energization of the coronal field outside the emergence region.



cf. Fisher et al (ApJS) 2020 - electric field inversion techniques. cf. Yeates 2017 - sparse electric fields.

#### **How to emerge SHARPs**

2. Apply steady electric field **A**/dt for fixed emergence time - e.g. dt=12 hours.



• Still testing: addition of twist informed by HMI measurements of  $\alpha = \frac{J_r}{B_r}$ 

#### **Full simulation [preliminary]**

 $r = R_{\odot}$ 



BMRs [82% tilt]





▶ Longitude averages of *B<sub>r</sub>*.

 $r = 2.5 R_{\odot}$ 



#### BMRs [82% tilt] + MF



#### HMI+PFSS

r10

- 5

- i)

-10

5

- 0

-5

-10



#### Conclusion

- Automated database of Bipolar Magnetic Regions from HMI/SHARPs.
- Flux transport model with BMRs overestimates end-of-cycle dipole by 24% compared to model with original SHARP shapes.

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- Python code for extracting database: <u>https://github.com/antyeates1983/sharps-bmrs</u>
- Ready-prepared file: <u>https://doi.org/10.7910/DVN/1Z7YMT</u> (Harvard Dataverse) for May 2010 to April 2020.
- Ongoing work [with postdoc Prantika Bhowmik]: magnetofrictional simulations of Cycle 24 — how does coronal field depend on
  - bipolar approximation?
  - emergence electric field?
  - twisting of emerging regions?





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## The End!