

New Topological Constraints on Magnetic Relaxation



Anthony Yeates

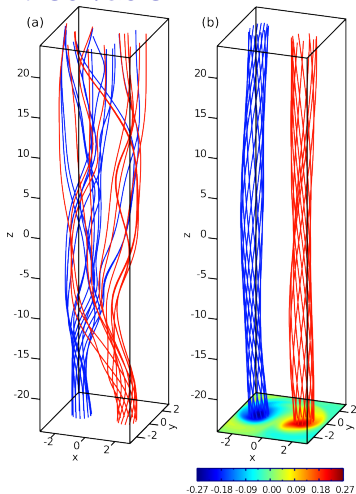
with

Gunnar Hornig (Dundee)

IUTAM Symposium, University College Dublin

July 2012

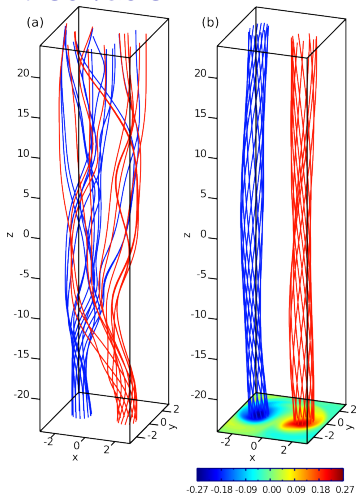
Motivation



- **Dundee simulation: relaxation of a braided magnetic field.**

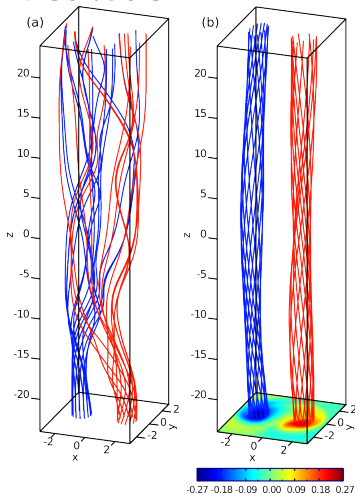
[Pontin et al. 2011; Wilmot-Smith et al. 2011]

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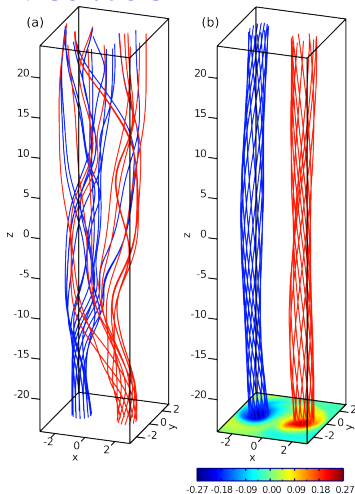
- ▶ **Dundee simulation: relaxation of a braided magnetic field.**
[Pontin et al. 2011; Wilmot-Smith et al. 2011]
- ▶ **Net effect of many “extreme” reconnection events.**

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- ▶ Dundee simulation: relaxation of a braided magnetic field.
[Pontin et al. 2011; Wilmot-Smith et al. 2011]
- ▶ Net effect of many “extreme” reconnection events.
- ▶ Taylor [1974]:
Turbulent reconnection destroys all ideal invariants except for total magnetic helicity
 \Rightarrow predict final state by minimising energy subject to constrained helicity.

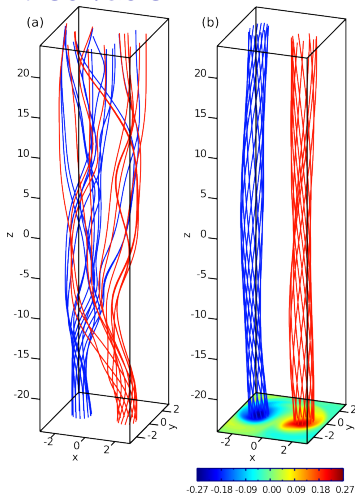
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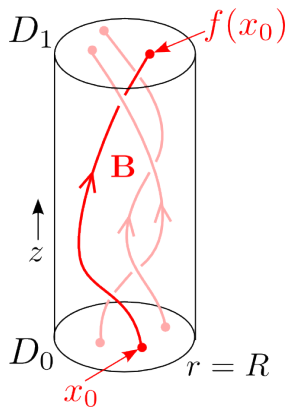
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Why is the Taylor state not reached?

Previous ideas:

- ▶ Constrain multiple partial helicities over sub-volumes [Bhattacharjee & Dewar 1982; Dixon et al. 1989; Turnbull 2012]
- ▶ High-order linking of field lines [Ruzmaikin & Akhmetiev 1994; Hornig & Mayer 2002]

Magnetic Braids



- ▶ Magnetic field $\mathbf{B}(r, \phi, z, t)$ on a cylinder, satisfying

$$\mathbf{B} \neq \mathbf{0}, \quad (1)$$

$$\mathbf{B}_r|_{r=R} = \mathbf{0}, \quad (2)$$

$$\mathbf{B}_z|_{D_0} = \mathbf{B}_z|_{D_1} > \mathbf{0}. \quad (3)$$

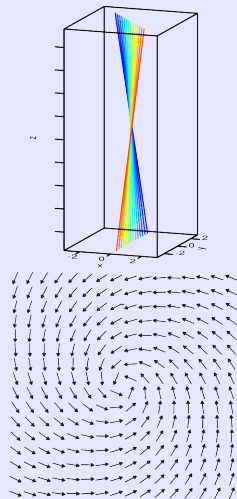
- ▶ Field line mapping

$$f : D_0 \rightarrow D_1.$$

Fixed Points

- Visualise with a “vector field” $\mathbf{v} = \mathbf{f} - \text{id}$.

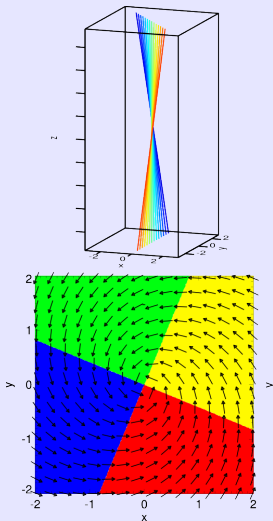
Example: uniform twist field



Fixed Points

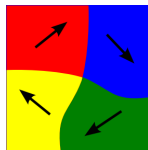
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- ▶ Colour map [Polymilis et al. 2003].

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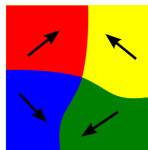
Fixed Points

- ▶ Visualise with a “vector field” $\mathbf{v} = \mathbf{f} - \text{id}$.
- ▶ Colour map [Polymilis et al. 2003].
- ▶ Each fixed point x_0 has a Poincaré index:



$$\text{ind}_{x_0} f = +1$$

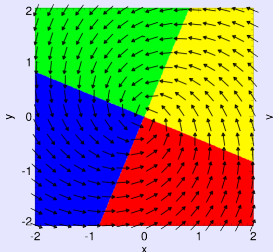
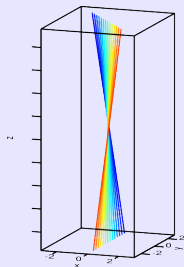
elliptic



$$\text{ind}_{x_0} f = -1$$

hyperbolic

Example: uniform twist field



Boundary Fixed Points

- ▶ To define index of fixed points on $f|_{\partial D_0}$, extend f outside D_0 .

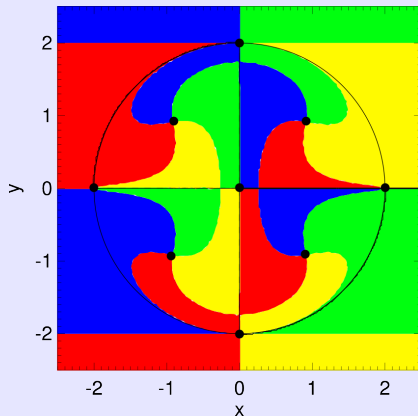
[Brown & Greene 1994; Ma & Wang 2001]

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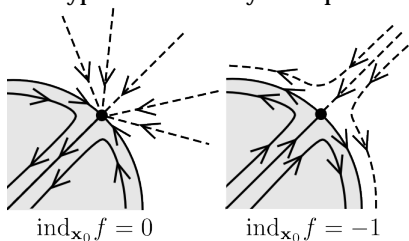


Boundary Fixed Points

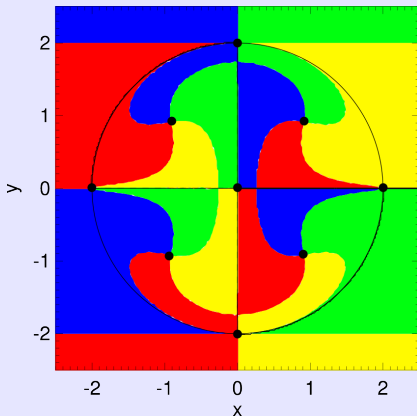
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[Brown & Greene 1994; Ma & Wang 2001]

- ▶ Two types of boundary fixed point:



Example:



Global Constraint

Define

- ▶ Total (net) interior index T_{int} (topological degree/Lefschetz number).
- ▶ Total (net) boundary index T_{∂} .

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Theorem [Hopf, 1929]

For the disc, $T_{\text{int}} = 1 - T_{\partial}$.



Heinz Hopf

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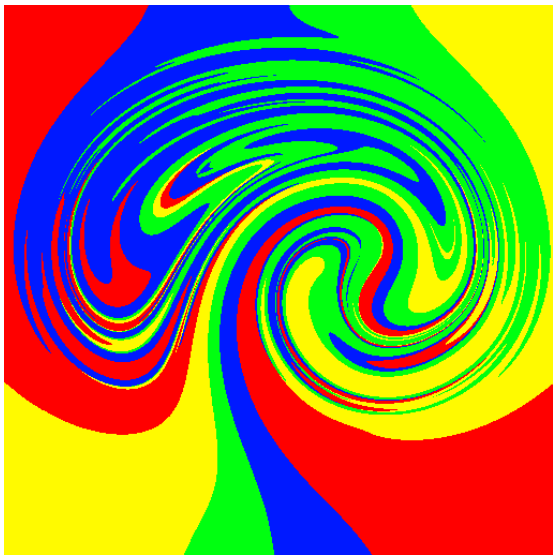
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Heinz Hopf

\Rightarrow If f remains fixed on the boundary,
then T_{int} must be conserved.

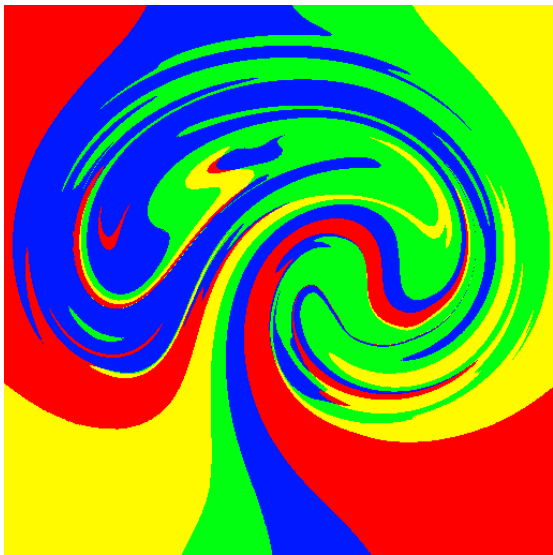
Dundee Simulation



$t = 0$

- ▶ Boundary shows that $T_{\text{int}} = 2$.

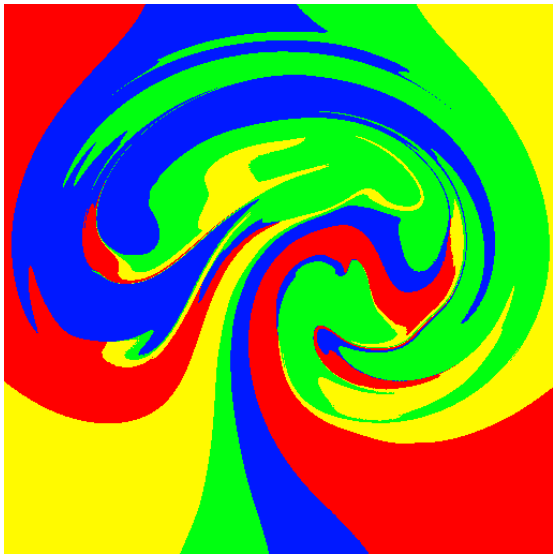
Dundee Simulation



$t = 35$

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Dundee Simulation



$t = 50$

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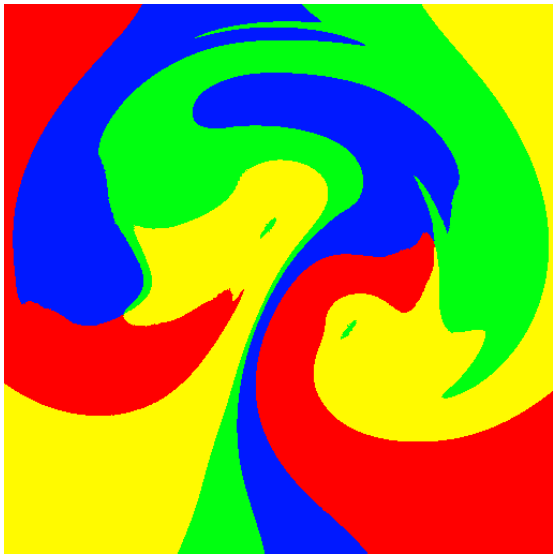
Dundee Simulation



$t = 80$

- ▶ Boundary shows that $T_{\text{int}} = 2$.

Dundee Simulation



$t = 110$

- ▶ Boundary shows that $T_{\text{int}} = 2$.

Dundee Simulation



$t = 290$

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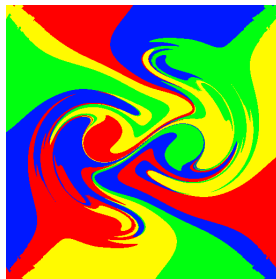
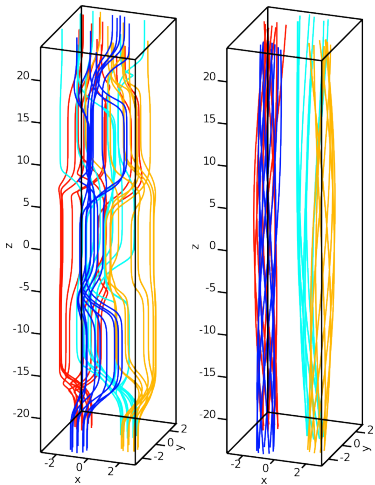


$t = 290$

- ▶ Boundary shows that $T_{\text{int}} = 2$.
- ▶ $T_{\text{int}} = 2$ is conserved, explaining failure to reach Taylor state.

The Silver Braid

- ▶ New initial state with $T_{\text{int}} = 3$ (inspired by “silver mixer” [Finn & Thiffeault 2011])



Initial



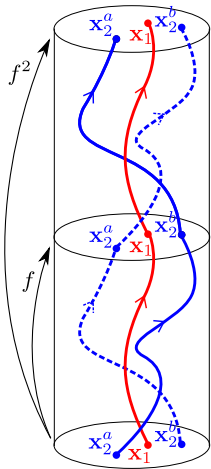
Final

Family of Invariants from Periodic Points

Denote the total index of $f^n = \underbrace{f \circ f \circ \dots \circ f}_{n \text{ times}}$ by T_{int}^n .

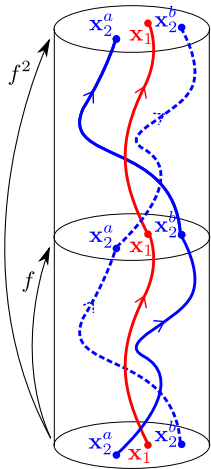
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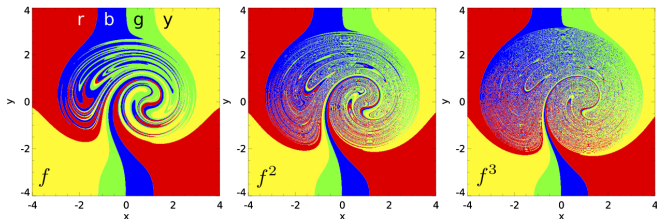


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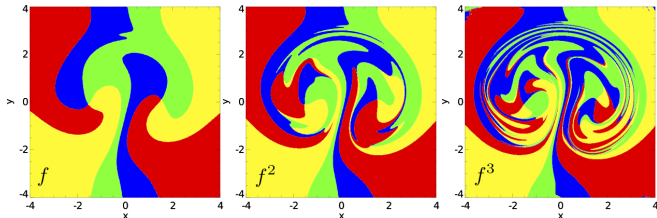
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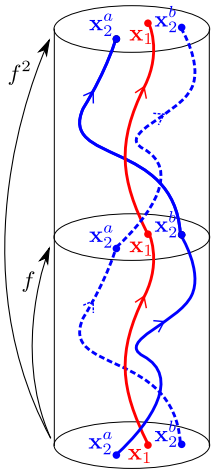


(b) Final state

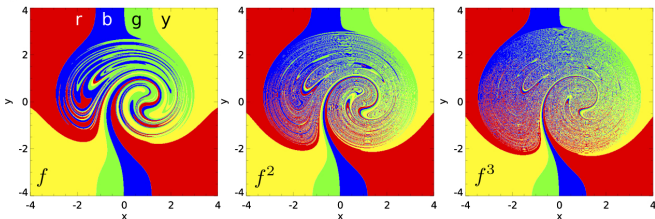


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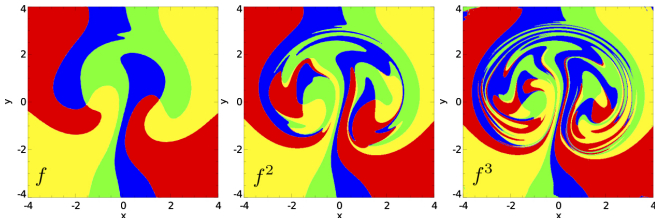
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Do these additional invariants constrain our predicted relaxed state?

Non-independence of Constraints

Usually “no”...

Theorem (Yeates & Hornig, 2011)

Suppose f has a finite number of fixed points at all iterations, and $\text{ind}_{x_0} f^q \in \{-1, 0\}$ for every periodic point x_0 on ∂D_0 . Then $T_{\text{int}}^q = T_{\text{int}}^1$ for all $q \in \mathbb{N}$ unless $T_{\text{int}}^1 = 1$ and $f|_{\partial D_0}$ has a rational rotation number greater than 1.

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$$T_{\text{int}}^1 > 1 \implies T_{\partial}^1 \neq 0 \quad (\text{all boundary fixed points already present at } f^1)$$

$$T_{\text{int}}^1 = 1 \implies T_{\partial}^1 = 0 \implies \text{minimal period could be higher (if rotation number } \in \mathbb{Q}\text{).}$$

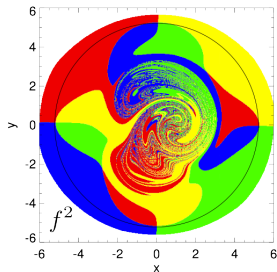
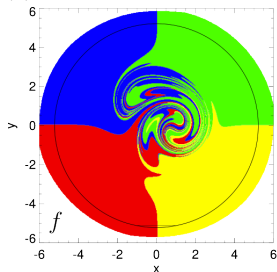
Exceptional Case

- ▶ Initial field with

$$T_{\text{int}}^{1,3,5,\dots} = 1,$$

$$T_{\text{int}}^{2,4,6,\dots} = 3.$$

(a) Initial state



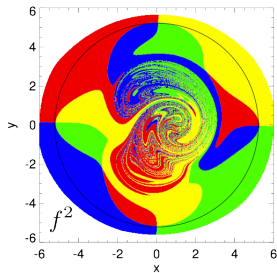
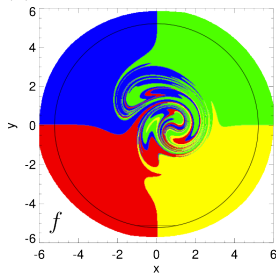
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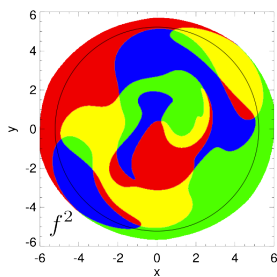
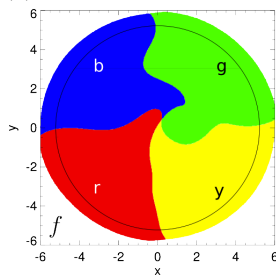
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(a) Initial state



(b) Final state



Conclusion

- ▶ Topological degree of the field line mapping can constrain the turbulent relaxation of a magnetised plasma.
- ▶ Applies to any continuous evolution of B providing that the field remains ideal on the boundary.
- ▶ If the degree of the initial state differs from that of the Taylor state, then the Taylor state will not be reached in the dynamical relaxation.
- ▶ There can be up to one further constraint from higher iterations of the field line mapping, but only for certain initial states with degree 1.

References

1. Yeates, Hornig & Wilmot-Smith, PRL 105, 085002 (2010).
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*Use periodic field lines to define a global reconnection rate.